FACULTY OF
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MASTER's THESIS

## Wilson loops in lattice gauge theories

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#### Abstract

This thesis is based on the article "Wilson loops in finite Abelian lattice gauge theories" by M. Forsström, J. Lenells and F. Viklund [1]. In the article, lattice gauge theories on $\mathbb{Z}^{4}$ for a finite Abelian structure group are considered and the expectation value for the Wilson loop observable at weak coupling is computed. The purpose of this thesis is to explain this article in more detail and to give the theory necessary to understand the article.

In this thesis, we consider the lattice $\mathbb{Z}^{4}$, the structure group $\mathbb{Z}_{n}$ and a faithful and one-dimensional representation. Basic theory for groups, representations and lattices is discussed. To state the main result, several definitions, e.g. the Wilson loop observable and Wilson action, are given. The main result is given as a theorem, where we have an inequality for the limit of the expectation value of the Wilson loop observable.

The theory necessary to prove the main result is given in this thesis. Theory for discrete exterior calculus is given in the third chapter. This includes theory for $k$-cells and $k$-forms as well as definitions and applications for both the exterior derivative and the co-derivative. Furthermore, two versions of the Poincaré lemma are given and applied to problems, e.g. for writing the given measure as a measure on plaquette configurations instead of spin configurations. The Hodge dual of the lattice $\mathbb{Z}^{r}$ is defined and both examples and lemmas, which are important for later proofs, are given.

In the fourth chapter, vortices and oriented surfaces are defined using the theory from the previous chapter. It is important to note that these definitions might differ from other sources. Various lemmas are stated and proved. The most important result in this chapter is a proposition, in which a probability is computed, that is applied several times in the proof of the main result.

Since the limit of the expectation of the Wilson loop observable is computed, both its existence and translation invariance must be proved. A more general theorem, which proves the existence and translation invariance for a real-valued function, is given and proved with Ginibre's inequality. This theorem is then applied to the Wilson loop observable.


The last part of this thesis is the proof of the main theorem. To prove this theorem, the theory and results given in the earlier chapters are applied. The proof is divided into two parts, which are then combined to achieve the desired result.

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## Chapter 1

## Introduction

Trying to mathematically describe the quantum field theories in the Standard Model is both a fascinating and difficult task. One method aiming to do this is lattice gauge theories. In these theories, space-time is discretised to a fourdimensional Euclidean lattice. Then, quantum field theories are approximated with methods from probability theory. Last, by taking a scaling limit, it might be possible to extend the model back to the continuous case and, thus, describe continuum quantum field theories.

Mathematical models for lattice gauge theories have been studied for a long time. They can be divided into two groups depending on whether its coupling constant is small or large. If the coupling constant is assumed to be small, we say that the gauge theory is in the weak coupling regime. The recently published articles by S. Chatterjee [2] and S. Cao [3] both compute the expectation values of a Wilson loop at weak coupling. This thesis will focus on the article by Forsström et al. [1, which extends Chatterjee's article. Since this model is in finite Abelian lattice gauge theories, it might not be relevant to the Standard Model, but it is still an interesting mathematical model to study.

The purpose of this thesis is to explain the third version of this article so that a master's student in mathematics understands it. The reader is expected to be familiar with probability theory, measure theory, complex analysis and algebra. Other necessary mathematical theory, e.g. group theory and discrete exterior calculus, will be discussed in this thesis and the reader is not required to have prior knowledge of them.

## Chapter 2

## The expectation value of the Wilson loop

In this chapter, the result for the expectation value of the Wilson loop observable will be given as a theorem. Before that, some theory necessary to understand this theorem is given. First, some basic theory for groups and representations is given. Second, the lattice used in the main theorem is described and terms necessary to the theorem, such as the Wilson action and the Wilson loop, are defined. Third, the main theorem is stated.

### 2.1 Groups

In this thesis, we will have a group $(G,+)$, which is a finite non-trivial Abelian group. Therefore, we have to define what that is, beginning with the definition of a group. A group is a set equipped with a binary operation that satisfies the group axioms in the definition below.

Definition 2.1. (Group) $(G,+)$ is a group if :
For all $a, b \in G: a+b \in G$ (Closure)
For all $a, b, c \in G:(a+b)+c=a+(b+c)$ (Associativity)
There exists an identity element $e \in G$ for all $a \in G: a+e=e+a=a$ (Identity)
There exists an inverse element $-a \in G$ for all $a \in G: a-a=-a+a=e$ (Inverse).

An Abelian group is a group, for which the elements commute.

Definition 2.2. (Abelian group) A group is Abelian if for all $a, b \in G: a+b=$ $b+a$ (Commutativity).

Example 2.3. Some examples of Abelian groups under ordinary addition are $(\mathbb{Z},+),(\mathbb{R},+)$ and $(\mathbb{Q},+)$. These groups are all infinite.

In this thesis, we assume that the group $G$, which is often called structure group, is $G=\left(\mathbb{Z}_{n},+\right)$, which is a finite group under addition modulo $n$. Here, the set $\mathbb{Z}_{n}$ is $\mathbb{Z}_{n}=\{0,1, \ldots, n-1\}$ and for $a, b \in \mathbb{Z}_{n}$ we have that $a+b:=a+b \bmod n$. The identity element for this group is $e=0$.

### 2.2 Representations

In this section, two different kinds of representations, which we will have use of later, will be defined. Last, the representation that will be used in this thesis is given. First, we must define what a representation is. Therefore, we define group homomorphisms and the general linear group before giving the definition for a representation.

Definition 2.4. (Group homomorphism) Let $(G, *)$ and $(H, \diamond)$ be groups. Then $f:(G, *) \mapsto(H, \diamond)$ is a homomorphism if for all $a, b \in G$ it holds that

$$
f(a * b)=f(a) \diamond f(b) .
$$

Definition 2.5. (General linear group of a vector space) The general linear group of a vector space $V$ is the group of linear invertible mappings from $V$ to $V$ and is denoted by $G L(V)$.

It is well known that this is a group with the operation matrix multiplication.
Definition 2.6. (Representation) A representation $\rho$ of $(G,+)$ on a vector space $V$ is a group homomorphism from $G$ to $G L(V)$.

From these definitions, it follows that the following property holds for a representation $\rho$ of $(G,+)$ :

$$
\begin{equation*}
\rho\left(g+g^{\prime}\right)=\rho(g) \rho\left(g^{\prime}\right) \quad \text { for } g, g^{\prime} \in G . \tag{2.1}
\end{equation*}
$$

We continue by discussing faithful representations.

Definition 2.7. (Faithful representation) A representation $\rho$ is faithful if it is injective.

A representation $\rho$ is injective if $\rho(g)=\rho\left(g^{\prime}\right)$ implies $g=g^{\prime}$. Let $\mathbb{I}_{V}$ be the identity matrix for the vector space $V$. Then, from equation (2.1) follows that

$$
\rho(e) \rho\left(g^{\prime}\right)=\rho\left(e+g^{\prime}\right)=\rho\left(g^{\prime}\right) \quad \text { for all } g^{\prime} \in G .
$$

Thus,

$$
\rho(e)=\mathbb{I}_{V} .
$$

In this thesis, $\rho$ is defined as a faithful and one-dimensional representation of the group $G=\mathbb{Z}_{n}$. That the representation is one-dimensional means that the dimension of the vector space $V$ is $\operatorname{dim}(V)=1$. We derive the representation and begin with calculating the value of $\rho(0)$. Since 0 is the identity element for this group, we have

$$
\begin{equation*}
\rho(0)=\rho(e)=1 \text {. } \tag{2.2}
\end{equation*}
$$

Since

$$
\rho(1)^{n}=\rho(n \cdot 1)=\rho(e)=1,
$$

it follows that $\rho(g)$ must be the $n$th roots of unity (Recall that a root of unity is a complex or real solution to $x^{n}=1$, where $n$ is a positive integer and the $n$th root of unity is given by $e^{2 k \pi i / n}$ ). Since $\rho$ is the roots of unity, we obtain

$$
\begin{equation*}
|\rho(g)|=1 \quad \text { for all } g \in G \tag{2.3}
\end{equation*}
$$

Last, the representation must take the same value for $g$ and $g+n$. Hence, it follows that the representation must be given by

$$
\begin{equation*}
\rho(g)=e^{g \cdot 2 \pi i m / n}, \quad g \in G=\mathbb{Z}_{n}, \tag{2.4}
\end{equation*}
$$

for some $m \in\{1, \ldots, n-1\}$ relatively prime to $n$. For faithfulness of the representation, it is necessary that $m$ is relatively prime to $n$.

We define a unitary representation and show that the representation $\rho$ is unitary.

Definition 2.8. (Unitary representation) A representation $\rho$ is unitary if $\rho(g)$ is a unitary operator. The operator $\rho(g)$ is unitary if $\rho(g)^{*}=\rho(g)^{-1}$ for every $g \in G$, where $\rho(g)^{*}$ is the conjugate transpose of $\rho(g)$.

To show that the representation $\rho$, defined in equation (2.4), is unitary, we calculate $\rho(g) \rho(g)^{*}$. The conjugate transpose of $\rho(g)$ is $\rho(g)^{*}=e^{-g \cdot 2 \pi i m / n}$ and

$$
\rho(g) \rho(g)^{*}=e^{g \cdot 2 \pi i m / n} e^{-g \cdot 2 \pi i m / n}=e^{0}=1 .
$$

Thus, $\rho(g)^{*}=\rho(g)^{-1}$ and the representation is unitary.
To summarise this section, the representation $\rho(g)=e^{g \cdot 2 \pi i m / n}$ is a faithful, unitary and one-dimensional representation of $G=\mathbb{Z}_{n}, \rho(0)=1$ and $\rho(g)$ is a root of unity. Last, the choice of $m$ will determine the representation $\rho$.

### 2.3 The lattice $\mathbb{Z}^{4}$

Consider the lattice $\mathbb{Z}^{4}$, which is four-dimensional with a vertex at every integer coordinate. An edge is called positively oriented if the coordinate increases when traversing it. From every vertex $x$, exactly four positively oriented edges emerge to its nearest neighbours. These edges are denoted by $d x_{i}$, where $i=1,2,3,4$. The opposite direction, the negative orientation, of an edge is given by $-d x_{i}$. An oriented plaquette $p$ is a pair of oriented edges that share a vertex and is defined as $p:=d x_{i} \wedge d x_{j}$, where $\wedge$ is the wedge product. The boundary $\partial p$ of a plaquette $p=d x_{j_{1}} \wedge d x_{j_{2}}$ is given by the edges $\partial p:=\left\{d x_{j_{1}}, d x_{j_{2}}+\left(d \mathbf{e}_{j_{1}}\right)_{j_{2}},-\left(d x_{j_{1}}+\right.\right.$ $\left.\left.\left(d \mathbf{e}_{j_{2}}\right)_{j_{1}}\right),-d x_{j_{2}}\right\}$, where $d \mathbf{e}_{j_{1}}$ and $d \mathbf{e}_{j_{1}}$ are unit vectors. An illustration of an oriented plaquette and its boundary is given in Figure 2.1 below. The terms oriented edges, oriented plaquettes and boundary are defined in more detail in Section 3.1 and Section 3.3.


Figure 2.1: An oriented plaquette $p=d x_{j_{1}} \wedge d x_{j_{2}}$ and the edges in its boundary $\partial p$.

We continue by defining some subsets of the lattice. For a given $N \geq 1$, let $B_{N}=\mathbb{Z}^{4} \cap[-N, N]^{4}$. Then the set $E_{N}$, whose elements are called $e$, is the set
containing all oriented edges, whose both endpoints are in $B_{N}$. An oriented edge with the opposite direction of $e$ is given by $-e$. An edge $e \in E_{N}$ is a boundary edge of $B_{N}$ if there is a plaquette $p$ with a corner that does not belong to $B_{N}$, such that $e \in \partial p$. The set of oriented plaquettes, which have all their boundary edges in $E_{N}$, are denoted by $P_{N}$. Elements in $P_{N}$ are called $p$. Let the set $\sum_{E_{N}}$ be the set of functions $\sigma: E_{N} \rightarrow G$ for which $\sigma_{e}=-\sigma_{-e}$ and $\sigma_{e} \neq 0$ for all $e \in E_{N}$. The elements $\sigma$ in $\Sigma_{E_{N}}$ are called spin configurations. For $\sigma \in \Sigma_{E_{N}}$, $(d \sigma)_{p}$ is defined as

$$
(d \sigma)_{p}:=\sigma_{e_{1}}+\sigma_{e_{2}}+\sigma_{e_{3}}+\sigma_{e_{4}}, \quad p \in P_{N}
$$

where $e_{i}$ are oriented edges in the boundary of $p$.
Last, we have some definitions for loops. Let $a_{0}, a_{1}, \ldots, a_{n}$ be vertices. A loop is then a sequence of oriented edges $e_{0}, e_{1}, \ldots, e_{n}$, where $e_{i}$ is a vertex between $a_{i-1}$ and $a_{i}$ for $i=1, \ldots, n$ such that $a_{n}=a_{0}$. Let $\gamma$ be a loop, then the length $|\gamma|$ is defined as the number of edges in it. A loop, which is closed and oriented, is called simple if all oriented edges in $\gamma$ are distinct from each other. Let $e$ be an edge in $\gamma$. Then the edge $e^{\prime} \in \gamma, e^{\prime} \neq e$, is a corner edge if $e^{\prime}$ or $-e^{\prime}$ shares a plaquette with the edge $e$. Edges which are not corner edges are called non-corner edges. This is illustrated in Figure 2.2, where the corner edges are given by the black edges and non-corner edges by the grey edges.


Figure 2.2: Corner and non-corner edges of a loop.

### 2.4 The Wilson loop observable

We begin with the definition of the Wilson action before defining the probability measure that will be used. Thereafter, the Wilson loop observable and the
expectation value of it are defined. In the next definition, we have the notation $\Re(x)$, which gives the real part of the complex number $x$.

Definition 2.9. The Wilson action corresponding to the structure group $G$, representation $\rho$ and fixed $N$ is given by

$$
S(\sigma):=-\sum_{p \in P_{N}} \Re \rho\left((d \sigma)_{p}\right), \quad \sigma \in \Sigma_{E_{N}}
$$

The probability measure $\mu_{H}$ is the uniform measure on the set $\sum_{E_{N}}$. From this measure, an associated probability measure $\mu_{\beta, N}$ on $\sum_{E_{N}}$ can be defined by weighting the uniform measure by $e^{-\beta S(\sigma)}$.

Definition 2.10. For each $\beta>0$, the probability measure $\mu_{\beta, N}$ on $\sum_{E_{N}}$ is

$$
\begin{equation*}
\mu_{\beta, N}(\sigma):=Z_{\beta, N}^{-1} e^{-\beta S(\sigma)} \mu_{H}(\sigma), \quad \sigma \in \Sigma_{E_{N}} \tag{2.5}
\end{equation*}
$$

where $Z_{\beta, N}$ is a normalising constant.
Definition 2.11. (The Wilson loop observable) Given a simple loop $\gamma \subseteq E_{N}$, the Wilson loop observable $W_{\gamma}$ is

$$
\begin{equation*}
W_{\gamma}:=W_{\gamma}(\sigma):=\rho\left(\sum_{e \in \gamma} \sigma_{e}\right), \quad \sigma \in \Sigma_{E_{N}} . \tag{2.6}
\end{equation*}
$$

Definition 2.12. (Expectation of the Wilson loop observable) The expectation of the Wilson loop observable with respect to the measure $\mu_{\beta, N}$, using free boundary conditions, is

$$
\mathbb{E}_{\beta, N}\left[W_{\gamma}\right]=\sum_{\sigma \in \sum_{E_{N}}} W_{\gamma}(\sigma) \mu_{\beta, N}(\sigma) .
$$

Free boundary conditions means that they are no modifications at the boundary of $B_{N}$. Last, the definition for the limit of this expectation.

Definition 2.13. (The limit of the expectation of the Wilson loop observable) The limit of the expectation of the Wilson loop observable $\left\langle W_{\gamma}\right\rangle_{\beta}$ is

$$
\left\langle W_{\gamma}\right\rangle_{\beta}:=\lim _{N \rightarrow \infty} \mathbb{E}_{\beta, N}\left[W_{\gamma}\right]
$$

when it exists.
In Chapter 5, we will see that the limit of the expected value of the Wilson loop observable both exists and is translation invariant.

### 2.5 The functions $\theta(\beta)$ and $\lambda(\beta)$

Before the main theorem is stated, two functions, $\theta(\beta)$ and $\lambda(\beta)$, must be defined. For $\beta \geq 0, \theta(\beta)$ is defined as

$$
\begin{equation*}
\theta(\beta):=\frac{\sum_{g \in G} \rho(g) e^{12 \beta \Re \rho(g)}}{\sum_{g \in G} e^{12 \beta \Re \rho(g)}} \tag{2.7}
\end{equation*}
$$

and $\lambda(\beta)$ as

$$
\begin{equation*}
\lambda(\beta):=\max _{g \in G \backslash\{0\}} \frac{e^{\beta \Re(\rho(g))}}{e^{\beta \Re(\rho(0))}} . \tag{2.8}
\end{equation*}
$$

We define a function $\phi_{\beta}(g)$ and write the functions $\lambda(\beta)$ and $\theta(\beta)$ with it. The function $\phi_{\beta}: G=\mathbb{Z}_{n} \rightarrow \mathbb{R}$ is defined by

$$
\begin{equation*}
\phi_{\beta}(g):=e^{\beta \Re \rho(g)}, \quad g \in G . \tag{2.9}
\end{equation*}
$$

Hence, the functions $\theta(\beta)$ and $\lambda(\beta)$ can be written as

$$
\begin{equation*}
\theta(\beta)=\frac{\sum_{g \in G} \rho(g) \phi_{\beta}(g)^{12}}{\sum_{g \in G} \phi_{\beta}(g)^{12}} \tag{2.10}
\end{equation*}
$$

and

$$
\begin{equation*}
\lambda(\beta)=\max _{g \in G \backslash\{0\}} \frac{\phi_{\beta}(g)}{\phi_{\beta}(0)} . \tag{2.11}
\end{equation*}
$$

### 2.6 The result for the expectation value

Theorem 2.14. Let $n \in \mathbb{Z}$ and $n \geq 2$. Consider lattice gauge theory with the structure group $G=\mathbb{Z}_{n}$ and a one-dimensional and faithful representation $\rho$ of G. Let $\gamma \in \mathbb{Z}^{4}$ be any simple oriented loop, $\ell=|\gamma|$ its length and $\ell_{c}$ the number of corner edges in $\gamma$. Then, the limit of the expectation value of the Wilson loop observable exists. Furthermore, for all $\beta_{0}>0$ chosen large enough, there exist constants $C^{\prime}=C^{\prime}\left(\beta_{0}\right)$ and $C^{\prime \prime}=C^{\prime \prime}\left(\beta_{0}\right)$ such that for any $\beta \geq \beta_{0}$, the following inequality holds:

$$
\begin{equation*}
\left|\left\langle W_{\gamma}\right\rangle_{\beta}-e^{-\ell(1-\theta(\beta))}\right| \leq C^{\prime}\left[\frac{\sqrt{\ell_{c}}}{\ell}+\lambda(\beta)^{2}\right]^{C^{\prime \prime}} \tag{2.12}
\end{equation*}
$$

The constants $C^{\prime}$ and $C^{\prime \prime}$ are defined by equation (6.66) in Section 6.3.
From this inequality, the value of $\left\langle W_{\gamma}\right\rangle_{\beta}$ (the limit of the expectation of the Wilson loop) can be estimated. The function $e^{-\ell(1-\theta(\beta))}$ takes values in $\left.] 0,1\right]$,
since the constant $\theta(\beta)$ only takes values in $] 0,1$ ] (This is proved in Section 6.3). The right-hand side is small, especially if we have a large $\beta$ and a long loop with few corners. If $\ell(1-\theta(\beta))$ is very large, the exponential function will be close to zero and hence $\left\langle W_{\gamma}\right\rangle_{\beta} \approx 0$. Likewise, if $\ell(1-\theta(\beta))$ is very small, the exponential function will be close to 1 and we have $\left\langle W_{\gamma}\right\rangle_{\beta} \approx 1$.

Another approach to this is that $\ell(1-\theta(\beta))$ is large when $\beta$ is chosen to be much smaller than $\ell$. Then it is likely that there are a lot of plaquettes $p$ near the loop $\gamma$ for which $d \sigma(p) \neq 0$. This implies that there are so many plaquettes that the model will have independence and therefore gives the result that the expectation is close to zero. Similarly, when $\ell(1-\theta(\beta))$ is small, $\beta$ is much larger than $\ell$. Hence, it is unlikely that there are any plaquettes $p$ near $\gamma$ for which $d \sigma(p) \neq 0$. Therefore, the expectation will be close to one.

Remark 2.15. The theorem holds when $\beta_{0}$ is chosen such that both $5(|G|-$ 1) $\lambda\left(\beta_{0}\right)^{2}<1$ and

$$
\begin{aligned}
& \max _{g_{1}, \ldots, g_{6} \in G}\left[\frac{\sum_{g \in G} e^{-2 \beta \sum_{k=1}^{6} \Re \rho\left(g+g_{k}\right)}}{\max _{g \in G} e^{-2 \beta \sum_{k=1}^{6} \Re \rho\left(g+g_{k}\right)}}-\left|\underset{g \in G}{\arg \max } e^{-2 \beta \sum_{k=1}^{6} \Re \rho\left(g+g_{k}\right)}\right|\right] \\
& \leq \frac{1-\cos (2 \pi / n)}{8}
\end{aligned}
$$

hold.

## Chapter 3

## Discrete exterior calculus

In this chapter, we focus on discrete exterior calculus for cell complexes of the lattice $\mathbb{Z}^{r}$ and its subset $B_{N}:=[-N, N]^{r} \cap \mathbb{Z}^{r}$ when $r \geq 1$. However, the theory will later only be used for cell complexes of $\mathbb{Z}^{4}$. We assume that the group $G$ is Abelian. Hence, the results can be applied to both $G=\mathbb{Z}$ and $G=\mathbb{Z}_{n}$. We begin with theory for $k$-cells and $k$-forms.

### 3.1 Oriented $k$-cells and $k$-forms

Earlier, both the oriented 1-cell and the oriented 2-cell were briefly discussed. Then they were called an oriented edge and an oriented plaquette. The 1-cell and the 2-cell will now be given new definitions and a general $k$-cell is defined. Consider the lattice $\mathbb{Z}^{r}$, which has a vertex at every point $x \in \mathbb{Z}^{r}$ with integer coordinates. Between every two neighbouring vertices, there exists an edge $\hat{e}$.

Definition 3.1. (Oriented 1-cell or edge) The edge ê between two neighbouring vertices can be divided into two oriented edges, e and $-e$, with the same endpoints. These oriented edges are called 1-cells.

To define the direction of an oriented edge, unit vectors must first be defined. Let $\mathbf{e}_{1}=(1,0,0, \ldots, 0), \mathbf{e}_{2}=(0,1,0, \ldots, 0), \ldots, \mathbf{e}_{r}=(0, \ldots, 0,1)$ and consider the edges from the origin to one of the vertices $\mathbf{e}_{1}, \mathbf{e}_{2}, \ldots, \mathbf{e}_{r}$. Edges, which have both the same length and direction as these edges, are unit vectors. These unit vectors are denoted by $d \mathbf{e}_{1}, \ldots, d \mathbf{e}_{r}$. Positive orientation is defined as the direction of the unit vectors and negative orientation as the opposite direction. Therefore, when $x \in \mathbb{Z}^{r}$ and $j \in\{1,2, \ldots, r\}$, a positively oriented edge $e$ can
be written as $e=d x_{j}:=x+d \mathbf{e}_{j}$. Then, the negatively oriented edge is given by $-e=-d x_{j}=x+\mathbf{e}_{j}-d \mathbf{e}_{j}$. Recall from Section 2.3 that in $\mathbb{Z}^{4}$ the set $E_{N}$ was denoted to be the set of oriented edges whose both endpoints are in the set $\mathbb{Z}^{4} \cap[-N, N]^{4}$. We expand this definition to the lattice $\mathbb{Z}^{r}$ and define $E_{N}$ to be the set of oriented edges, whose both endpoints are in the set $B_{N}=\mathbb{Z}^{r} \cap[-N, N]^{r}$. Furthermore, the set of only positively oriented edges in $B_{N}$ is given by the set $E_{N}^{+}$and the set of only negatively oriented edges is given by $E_{N}^{-}$.

In the definition for oriented $k$-cells, the notation $e_{1} \wedge e_{2}$ will be used. This is the wedge product between two oriented edges $e_{1}$ and $e_{2}$. The wedge product is defined to have the following properties: For two edges $e_{1}, e_{2} \in E_{N}$, we have

$$
e_{1} \wedge e_{2}=-\left(e_{2} \wedge e_{1}\right)=\left(-e_{2}\right) \wedge e_{1}=e_{2} \wedge\left(-e_{1}\right)
$$

and

$$
e_{1} \wedge e_{1}=0
$$

For two or more oriented edges $e_{1}, e_{2}, \ldots, e_{k}$ in $E_{N}$, the wedge product $e_{1} \wedge \cdots \wedge e_{k}$ is zero if the edges $e_{1}, \ldots, e_{k}$ do not have a common endpoint.

Definition 3.2. (Oriented $k$-cell) Let $e_{1}, \ldots, e_{k}$ be oriented edges with $e_{1} \wedge \cdots \wedge e_{k} \neq 0$. Then $e_{1} \wedge \cdots \wedge e_{k}$ is an oriented $k$-cell.

A $k$-cell $e_{1} \wedge \cdots \wedge e_{k}$ is positively oriented if there is an $x \in \mathbb{Z}^{r}$ and $j_{1}<$ $j_{2}<\cdots<j_{k}$ such that $e_{i}=d x_{j_{i}}$. Correspondingly, the k-cell $-\left(e_{1} \wedge \cdots \wedge e_{k}\right)$ is negatively oriented. Some examples of $k$-cells are illustrated in the figure below. Notice that a 0 -cell is a vertex.


Figure 3.1: Examples of $k$-cells.

From now on, a $k$-cell will refer to an oriented $k$-cell.

Definition 3.3. (Oriented plaquette) An oriented plaquette $p$ is an oriented 2cell.

Since the oriented plaquette is a $k$-cell, the oriented plaquette $p=e_{1} \wedge e_{2}$ is positively oriented if there is an $x \in \mathbb{Z}^{r}$ and $j_{1}<j_{2}$ such that $e_{1}=d x_{j_{1}}$ and $e_{2}=d x_{j_{2}}$. Therefore, a positively oriented plaquette can be written as $p:=d x_{j_{1}} \wedge d x_{j_{2}}$. The set of oriented plaquettes whose all boundary edges are in $E_{N}$ is given by $P_{N}$.

We continue with theory about $k$-forms, which are discrete differential forms. Definition 3.4. ( $k$-form) A function $f$, which is defined on a subset of $k$-cells in $\mathbb{Z}^{r}$, for which $f(c)=-f(-c)$ is a $k$-form.

Consider a $k$-form $f$, which has the value $f_{j_{1}, \ldots j_{k}}(x)$ on the $k$-cell $c=d x_{j_{1}} \wedge$ $\cdots \wedge d x_{j_{k}}$. Then the $k$-form $f$ is represented by

$$
\begin{equation*}
f(x)=\sum_{1 \leq j_{1}<\cdots<j_{k} \leq r} f_{j_{1}, \ldots, j_{k}}(x) d x_{j_{1}} \wedge \cdots \wedge d x_{j_{k}} . \tag{3.1}
\end{equation*}
$$

The function $f_{j_{1}, \ldots j_{k}}(x)$ can also be written shortly as $f_{c}$.
Example 3.5. Two examples of 2 -forms in $\mathbb{Z}^{4}$ are

$$
f_{1}(x)=f_{12}(x) d x_{1} \wedge d x_{2}+f_{34}(x) d x_{3} \wedge d x_{4}
$$

and

$$
f_{2}(x)=f_{12}(x) d x_{1} \wedge d x_{2}+f_{13}(x) d x_{1} \wedge d x_{3} .
$$

Next, we define both non-trivial $k$-forms and the set support.
Definition 3.6. (Non-trivial $k$-form) A $k$-form $f$ is non-trivial if there exists one or more $k$-cells c for which $f_{c} \neq 0$.

Definition 3.7. (Support) The support of a $k$-form $f$ is the set of oriented $k$ cells $c$, for which the $k$-form $f$ satisfies $f(c) \neq 0$. The support of $f$ is denoted by supp $f$.

Example 3.8. Let $A$ be a subset of $E_{N}$ and define for each $e \in E_{N}$ the function $f$ by

$$
f(e):= \begin{cases}1 & \text { if } e \in A \\ -1 & \text { if }-e \in A \\ 0 & \text { otherwise }\end{cases}
$$

Since $f(e)=-f(-e)$ and $e$ is a 1 -cell, we have that $f$ is a 1 -form. The function only takes the value zero outside $A$, hence supp $f=A$.

The set $\Sigma_{E_{N}}$ in $\mathbb{Z}^{r}$ is defined as the set of $G$-valued 1-forms, whose support is in $E_{N}$. Its elements are denoted by $\sigma$ and called spin configurations. Recall the definition for $\Sigma_{E_{N}}$ in $\mathbb{Z}^{4}$ : the set of functions $\sigma: E_{N} \rightarrow G$ for which $\sigma_{e}=-\sigma_{-e}$ and $\sigma_{e} \neq 0$ for all $e \in E_{N}$. We see that this agrees with the definition for $\mathbb{Z}^{r}$.

Last, the restriction $\left(\left.\sigma\right|_{E}\right)_{e}$ for a spin configuration $\sigma \in \Sigma_{E_{N}}$, a set $E \subseteq E_{N}$ and an edge $e \in E_{N}$ is defined by

$$
\left(\left.\sigma\right|_{E}\right)_{e}:= \begin{cases}\sigma_{e} & \text { if } e \in E  \tag{3.2}\\ 0 & \text { otherwise }\end{cases}
$$

### 3.2 The exterior derivative

In this section, the exterior derivative is defined and some examples are given. When the exterior derivative operates on a $k$-form, the result is a $k+1$-form.

Definition 3.9. (Exterior derivative) Given $h: \mathbb{Z}^{r} \rightarrow G, x \in \mathbb{Z}^{r}$, and $i \in$ $\{1,2, \ldots, r\}$, let

$$
\begin{equation*}
\partial_{i} h(x):=h\left(x+e_{i}\right)-h(x) . \tag{3.3}
\end{equation*}
$$

If $k \in\{0,1, \ldots, r-1\}$ and $f$ is a $G$-valued $k$-form, the exterior derivative $d$ of $f$ is defined as

$$
\begin{equation*}
d f(x)=\sum_{1 \leq j_{1}<\cdots<j_{k} \leq r} \sum_{i=1}^{r} \partial_{i} f_{j_{1}, \ldots, j_{k}}(x) d x_{i} \wedge\left(d x_{j_{1}} \wedge \ldots \wedge d x_{j_{k}}\right), \quad x \in \mathbb{Z}^{r} \tag{3.4}
\end{equation*}
$$

Next is an example for how the exterior derivative is calculated for 2 -forms.
Example 3.10. We calculate the exterior derivative of the 2 -forms
$f_{1}=f_{12}(x) d x_{1} \wedge d x_{2}+f_{34}(x) d x_{3} \wedge d x_{4}$ and $f_{2}=f_{12}(x) d x_{1} \wedge d x_{2}+f_{13}(x) d x_{1} \wedge d x_{3}$. We begin with calculating the exterior derivative for the 2 -form $f_{1}$. By the definition of the exterior derivative, we have

$$
\begin{aligned}
d f_{1}= & d\left(f_{12}(x) d x_{1} \wedge d x_{2}+f_{34}(x) d x_{3} \wedge d x_{4}\right) \\
= & \partial_{3} f_{12}(x) d x_{1} \wedge d x_{2} \wedge d x_{3}+\partial_{4} f_{12}(x) d x_{1} \wedge d x_{2} \wedge d x_{4} \\
& +\partial_{1} f_{34}(x) d x_{3} \wedge d x_{4} \wedge d x_{1}+\partial_{2} f_{34}(x) d x_{3} \wedge d x_{4} \wedge d x_{2}
\end{aligned}
$$

We calculate the exterior derivative for $f_{2}$ :

$$
\begin{aligned}
d f_{2}= & d\left(f_{12}(x) d x_{1} \wedge d x_{2}+f_{13}(x) d x_{1} \wedge d x_{3}\right) \\
= & \partial_{3} f_{12}(x) d x_{1} \wedge d x_{2} \wedge d x_{3}+\partial_{4} f_{12}(x) d x_{1} \wedge d x_{2} \wedge d x_{4} \\
& +\partial_{2} f_{13}(x) d x_{1} \wedge d x_{3} \wedge d x_{2}+\partial_{4} f_{13}(x) d x_{1} \wedge d x_{3} \wedge d x_{4} \\
= & \left(\partial_{3} f_{12}(x)-\partial_{2} f_{13}(x)\right) d x_{1} \wedge d x_{2} \wedge d x_{3}+\partial_{4} f_{12}(x) d x_{1} \wedge d x_{2} \wedge d x_{4} \\
& +\partial_{4} f_{13}(x) d x_{1} \wedge d x_{3} \wedge d x_{4}
\end{aligned}
$$

A closed $k$-form can now be defined.
Definition 3.11. (Closed $k$-form) $A k$-form $f$ is closed if $d f=0$.
Last, we see that we earlier had an definition for the exterior derivative on a spin configuration on the lattice $\mathbb{Z}^{4}$.

Example 3.12. Recall the definitions in Section 2.3. For $\sigma \in \Sigma_{E_{N}}$, we defined

$$
(d \sigma)_{p}:=\sigma_{e_{1}}+\sigma_{e_{2}}+\sigma_{e_{3}}+\sigma_{e_{4}}, \quad p \in P_{N}
$$

where $e_{i}$ are edges in the boundary of $p$. This is the exterior derivative of $\sigma \in$ $\Sigma_{E_{N}}$.

### 3.3 Boundary and co-boundary

Recall the definition for the boundary of a plaquette on $\mathbb{Z}^{4}$ in Section 2.3. We now expand this definition to $\mathbb{Z}^{r}$. When $x \in \mathbb{Z}^{r}$ and $j_{1}<j_{2}$, the boundary of a plaquette $p=d x_{j_{1}} \wedge d x_{j_{2}}$ is defined as

$$
\partial p:=\left\{d x_{j_{1}}, d x_{j_{2}}+\left(d \mathbf{e}_{j_{1}}\right)_{j_{2}},-\left(d x_{j_{1}}+\left(d \mathbf{e}_{j_{2}}\right)_{j_{1}}\right),-d x_{j_{2}}\right\} .
$$

This is illustrated in Figure 3.2 on page 17. The boundary for a plaquette can also be defined as the set of edges $d x_{j}$ for which

$$
\left(d\left(\mathbb{I}_{x=\hat{x}} d x_{j}\right)\right)_{p}=1
$$

where $\mathbb{I}_{x=\hat{x}}$ is the Kronecker delta function of $x$ with mass at $\hat{x} \in \mathbb{Z}^{r}$. Next, we prove that these two definitions are equivalent.


Figure 3.2: The boundary of the plaquette $p=d x_{j_{1}} \wedge d x_{j_{2}}$.
We prove this by letting the exterior derivative operate on the 1 -form $\mathbb{I}_{x=\hat{x}} d x_{j}$, when $\hat{x} \in \mathbb{Z}^{r}$ and $j \in\{1,2, \ldots, r\}$. By equations (3.4) and (3.3) for the function $h=\mathbb{I}_{x=\hat{x}}$, we have

$$
\begin{aligned}
d\left(\mathbb{I}_{x=\hat{x}} d x_{j}\right) & =\sum_{i=1}^{r}\left(\partial_{i} \mathbb{I}_{x=\hat{x}}\right) d x_{i} \wedge d x_{j}=\sum_{i=1}^{r}\left(\mathbb{I}_{x+\mathbf{e}_{i}=\hat{x}}-\mathbb{I}_{x=\hat{x}}\right) d x_{i} \wedge d x_{j} \\
& =\sum_{i=1}^{r}\left(\mathbb{I}_{x+\mathbf{e}_{i}=\hat{x}}-\mathbb{I}_{x=\hat{x}}\right) d x_{i} \wedge d x_{j}-\sum_{i=j+1}^{r}\left(\mathbb{I}_{x+\mathbf{e}_{i}=\hat{x}}-\mathbb{I}_{x=\hat{x}}\right) d x_{j} \wedge d x_{i} .
\end{aligned}
$$

For edges that are not in the boundary for neither the positive nor the negative direction, both $\mathbb{I}_{x+\mathbf{e}_{i}=\hat{x}}=0$ and $\mathbb{I}_{x=\hat{x}}=0$. Hence,

$$
\left(d\left(\mathbb{I}_{x=\hat{x}} d x_{j}\right)\right)_{p}= \begin{cases}\left(\mathbb{I}_{x+\mathbf{e}_{i}=\hat{x}}-\mathbb{I}_{x=\hat{x}}\right) & \text { if } e \in \partial p \\ -\left(\mathbb{I}_{x+\mathbf{e}_{i}=\hat{x}}-\mathbb{I}_{x=\hat{x}}\right) & \text { if }-e \in \partial p \\ 0 & \text { otherwise }\end{cases}
$$

where $p$ is chosen as an oriented plaquette and $e=\hat{x}+d \mathbf{e}_{j}=d \hat{x}_{j}$ as an oriented edge. Since $\mathbb{I}_{x+\mathbf{e}_{i}=\hat{x}}-\mathbb{I}_{x=\hat{x}}=1$ for edges in the boundary, we have

$$
\left(d\left(\mathbb{I}_{x=\hat{x}} d x_{j}\right)\right)_{p}= \begin{cases}1 & \text { if } e \in \partial p  \tag{3.5}\\ -1 & \text { if }-e \in \partial p \\ 0 & \text { otherwise }\end{cases}
$$

Thus, the definitions are equivalent.
From this follows that if $f$ is a 1 -form, $1 \leq j_{1}<j_{2} \leq r$ and the plaquette $p$ is given by $p=d x_{j_{1}} \wedge d x_{j_{2}}$, then we have that

$$
\begin{equation*}
(d f)_{p}=(d f)_{j_{1}, j_{2}}(x)=\sum_{e \in \partial p} f_{e} . \tag{3.6}
\end{equation*}
$$

Next, the definition for the boundary for general $k$-cells and a general version of formula (3.6) for $k$-forms.

Definition 3.13. (Boundary) Let $k \in\{1,2, \ldots, r\}$ and $c$ be a $k$-cell. The boundary $\partial \hat{c}$ of $c$ is the set of $(k-1)$-cells $\hat{c}=d x_{j_{1}} \wedge \cdots \wedge d \hat{x}_{j_{k_{1}}}$ for which

$$
\left(d\left(\mathbb{I}_{x=\hat{x}} d \hat{x}_{j_{1}} \wedge \cdots \wedge d \hat{x}_{j_{k-1}}\right)\right)_{c}=1
$$

Let $c_{0}$ be a $(k+1)$-cell, then

$$
\begin{equation*}
(d f)_{c_{0}}=\sum_{c \in \partial c_{0}} f_{c} . \tag{3.7}
\end{equation*}
$$

The boundary is connected with the co-boundary, which is defined next.
Definition 3.14. (Co-boundary) The co-boundary $\hat{\partial} c$ of $a k$-cell $c$ is defined as the set of every $(k+1)$-cells $\hat{c}$ for which $c \in \partial \hat{c}$.

Example 3.15. The co-boundary $\hat{\partial} e$ of an edge $e$ is the set of all oriented plaquettes $p$ that contain $e$. An illustration is given in Figure 3.3. Furthermore, $e \in \partial p$ is equivalent to $p \in \hat{\partial} e$.


Figure 3.3: An edge and two of the plaquettes in its co-boundary.

Example 3.16. For an edge $e \in E_{N}$, a plaquette $p$ cannot contain both $e$ and $-e$. Therefore, the intersection $\hat{\partial} e \cap \hat{\partial}(-e)$ is empty. Moreover, if $p \in \hat{\partial}(-e)$, then $-p \in \hat{\partial} e$.

Last, boundary cells are defined and one example is given. Notice that a $k$-cell is in $B_{N}$ if all corners of the $k$-cell is in $B_{N}$.

Definition 3.17. (Boundary cell) For $k \in\{0,1, \ldots, r-1\}$, a $k$-cell c in $B_{N}$ is a boundary cell of $B_{N}$ if there exists a $(k+1)$-cell $\hat{c} \in \hat{\partial} c$ which is not in $B_{N}$.

Example 3.18. An edge $e \in E_{N}$ is a boundary edge of $B_{N}$ if there exists one plaquette $p \in \hat{\partial} e$ which is outside $B_{N}$. A plaquette $p \in P_{N}$ is a boundary plaquette of $B_{N}$ if there exists a 3 -cell $c \in \hat{\partial} p$ which is outside $B_{N}$. This is illustrated in Figure 3.4, where one of the boundary edges is given by the blue edge and one of the boundary plaquettes is given by the blue plaquette.


Figure 3.4: Example of a boundary edge and a boundary plaquette.

### 3.4 The Poincaré Lemma

In this section, the Poincaré Lemma will be given without proof and then applied to both the set $\Sigma_{P_{N}}$ (which will be defined later in this section) and for writing the earlier given measure as a measure on elements in the set $\Sigma_{P_{N}}$ instead of spin configurations. Another version of this lemma, the Poincaré Lemma for the co-derivative, will be given in the next section. Before the Poincaré lemma is stated, the definition for a box and a cube is given.

Definition 3.19. (Box and cube) $A$ set of the form $\left(\left[a_{1}, b_{1}\right] \times \cdots \times\left[a_{r}, b_{r}\right]\right) \cap \mathbb{Z}^{r}$ is a box if for each $j \in\{1,2, \ldots, r\},\left\{a_{j}, b_{j}\right\} \subset \mathbb{Z}$ satisfies $a_{j}<b_{j}$. A box is a cube if all intervals $\left[a_{j}, b_{j}\right], 1 \leq j \leq r$, are of the same length.

An example of a box is the set $B_{N}$.
Lemma 3.20. (The Poincaré lemma). Let $k \in\{0,1, \ldots, r-1\}$ and $B$ be a box in $\mathbb{Z}^{r}$. Then the exterior derivative $d$ is a surjective mapping from the set of $G$-valued $k$-forms with support contained in $B$ onto the set of $G$-valued closed $(k+1)$-forms with support contained in B. If the group $G$ is finite and the number of closed $G$-valued $k$-forms with support contained in $B$ is $m$, then the map is an
$m$-to- 1 correspondence. If $k \in\{0,1,2, \ldots, r-1\}$ and $f$ is a closed $(k+1)$-form that vanishes on the boundary of $B$, then there is a $k$-form $h$ that also vanish on the boundary of $B$ and satisfies $d h=f$.

For the proof, see [2, Lemma 2.2].

### 3.4.1 The set $\Sigma_{P_{N}}$

The set $\sum_{P_{N}}$ is defined as the set of closed $G$-valued 2-forms $\omega$ with support in $P_{N}$. The 2-forms in $\sum_{P_{N}}$ are called plaquette configurations. Since the set $\Sigma_{P_{N}}$ only contains closed forms, we obtain

$$
d \omega=0 \quad \text { if } \omega \in \Sigma_{P_{N}}
$$

Recall that the 1-forms in the set $\Sigma_{E_{N}}$ do not have to be closed, while the set $\Sigma_{P_{N}}$ is defined to only contain closed 2-forms. The restriction for a plaquette configuration $\omega \in \Sigma_{P_{N}}$ is defined similar to how a restriction was defined for a spin configuration $\sigma \in \Sigma_{E_{N}}$ in equation (3.2). For $\omega \in \Sigma_{P_{N}}$, a set $P \subseteq P_{N}$ and a plaquette $p \in P_{N}$, the restriction $\left(\left.\omega\right|_{P}\right)_{p}$ is defined as

$$
\left(\left.\omega\right|_{P}\right)_{p}:= \begin{cases}\omega_{p} & \text { if } p \in P  \tag{3.8}\\ 0 & \text { otherwise }\end{cases}
$$

Next is an example of a restriction for a 2-form in $P_{N}$.
Example 3.21. For $\omega \in \Sigma_{P_{N}}$, consider the restriction $\left.\omega\right|_{\text {supp } \omega}$. By definition, each $\omega$ has support in $P_{N}$. Hence, $\left.\omega\right|_{\text {supp }} \omega=\omega$.

We apply the Poincaré lemma to the $k$-form $\omega \in \Sigma_{P_{N}}$ and the box $B_{N}$. Recall, that the set of $G$-valued 1-forms with support contained in $P_{N}$ is the set $\Sigma_{E_{N}}$ and that the 1 -forms in $\Sigma_{E_{N}}$ are called $\sigma$. Furthermore, we have that both $\sigma$ and $\omega$ vanishes on the boundary of $P_{N}$. By the Poincaré Lemma, we have that the exterior derivative from the set $\Sigma_{E_{N}}$ onto the set $\Sigma_{P_{N}}$ is a surjective map. Since $G$ is finite, there exists an $m$ for which the map is a $m$-to- 1 correspondence. Furthermore, since $\omega$ is closed, there is a $\sigma$ such that $d \sigma=\omega$. As a result of these three conclusions, we have that $\omega \in \Sigma_{P_{N}}$ if and only if there exists a $\sigma \in \Sigma_{E_{N}}$ for which $d \sigma=\omega$.

### 3.4.2 The measure

Earlier, the measure $\mu_{\beta, N}$ was defined in Definition 2.5 as a measure on spin configurations:

$$
\mu_{\beta, N}(\sigma)=Z_{\beta, N}^{-1} e^{\beta \sum_{p \in P_{N}} \Re \rho\left((d \sigma)_{p}\right)} \mu_{H}(\sigma), \quad \sigma \in \Sigma_{E_{N}} .
$$

We will later use this measure on plaquette configurations instead of spin configurations. Therefore, we have to map spin configurations to plaquette configurations. A problem is that every plaquette configuration does not arise from a spin configuration. Nonetheless, the Poincaré Lemma says that $\omega \in \Sigma_{P_{N}}$ if and only if there exists a $\sigma \in \Sigma_{E_{N}}$ such that $d \sigma=\omega$. Furthermore, the map from spin configurations $\sigma \in \Sigma_{E_{N}}$ to plaquette configurations $\omega \in \Sigma_{P_{N}}$ is a many-to-1 correspondence. Since $\left|\left\{\sigma \in \Sigma_{E_{N}}: d \sigma=\omega\right\}\right|$ does not depend on the choice of $\omega$, we have

$$
\sum_{\sigma \in \Sigma_{E_{N}}: d \sigma=\omega} e^{\beta \sum_{p \in P_{N}} \Re \rho\left((d \sigma)_{p}\right)}=\left|\left\{\sigma \in \Sigma_{E_{N}}: d \sigma=\omega\right\}\right| e^{\beta \sum_{p \in P_{N}} \Re \rho\left(\omega_{p}\right)} .
$$

Thus, a measure for plaquette configurations can be created:

$$
\begin{align*}
\mu_{\beta, N}\left(\left\{\sigma \in \Sigma_{E_{N}}: d \sigma=\omega\right\}\right) & =\frac{\sum_{\sigma \in \Sigma_{E_{N}}: d \sigma=\omega} e^{\beta \sum_{p \in P_{N}} \Re \rho\left((d \sigma)_{p}\right)}}{\sum_{\sigma \in \Sigma_{E_{N}}} e^{\beta \sum_{p \in P_{N}} \Re \rho\left((d \sigma)_{p}\right)}} \\
& =\frac{e^{\beta \sum_{p \in P_{N}} \Re \rho\left(\omega_{p}\right)}}{\sum_{\omega^{\prime} \in \Sigma_{P_{N}}} e^{\beta \sum_{p \in P_{N}} \Re \rho\left(\omega_{p}^{\prime}\right)}}  \tag{3.9}\\
& =\frac{\prod_{p \in P_{N}} \phi_{\beta}\left(\omega_{p}\right)}{\sum_{\omega^{\prime} \in \Sigma_{P_{N}}} \prod_{p \in P_{N}} \phi_{\beta}\left(\omega_{p}^{\prime}\right)}=\mu_{\beta, N}(\{\omega\}) .
\end{align*}
$$

### 3.5 The co-derivative and its Poincaré Lemma

The definition of the co-derivative is similar to the definition of the exterior derivative. They both operate on a $k$-form, but when the co-derivative operates the result is a $k-1$-form instead of a $k+1$-form.

Definition 3.22. (Co-derivative) Given $h: \mathbb{Z}^{r} \rightarrow G, x \in \mathbb{Z}^{r}$, and an $i \in$ $\{1,2, \ldots, r\}$, let

$$
\begin{equation*}
\bar{\partial}_{i} h(x):=h(x)-h\left(x-\mathbf{e}_{i}\right) . \tag{3.10}
\end{equation*}
$$

If $k \in\{1,2, \ldots, r-1\}$ and $f$ is a $G$-valued $k$-form, the co-derivative $\delta$ of $f$ is defined by
$\delta f(x):=\sum_{1 \leq j_{1}<\cdots<j_{k} \leq r} \sum_{i=1}^{k}(-1)^{i} \bar{\partial}_{j_{i}} f_{j_{1}, \ldots, j_{k}}(x) d x_{j_{1}} \wedge \cdots \wedge d x_{j_{i-1}} \wedge d x_{j_{i+1}} \wedge \ldots \wedge d x_{j_{k}}$
for $x \in \mathbb{Z}^{r}$.
Example 3.23. We calculate the co-derivative of the 2 -forms $f_{1}=f_{12}(x) d x_{1} \wedge$ $d x_{2}+f_{34}(x) d x_{3} \wedge d x_{4}$ and $f_{2}=f_{12}(x) d x_{1} \wedge d x_{2}+f_{13}(x) d x_{1} \wedge d x_{3}$.

By definition, we have that the co-derivative of $f_{1}$ is

$$
\delta f_{1}(x)=-\bar{\partial}_{1} f_{12} d x_{2}+\bar{\partial}_{2} f_{12} d x_{1}-\bar{\partial}_{3} f_{34} d x_{4}+\bar{\partial}_{4} f_{34} d x_{3}
$$

and the co-derivative of $f_{2}$ is

$$
\delta f_{2}(x)=-\bar{\partial}_{1} f_{12} d x_{2}+\bar{\partial}_{2} f_{12} d x_{1}-\bar{\partial}_{1} f_{13} d x_{3}-\bar{\partial}_{3} f_{13} d x_{1}
$$

Similarly to the calculations of $\left(d\left(\mathbb{I}_{x=\hat{x}} d x_{j}\right)\right)_{p}$, we calculate the co-derivative of the 2 -form $f:=\mathbb{I}_{x=\hat{x}} d x_{i_{1}} \wedge d x_{i_{2}}$ where $x \in \mathbb{Z}^{r}$. Let $\hat{x} \in \mathbb{Z}^{r}, 1 \leq i_{1}<i_{2} \leq r$ and $p_{0}=d \hat{x}_{i_{1}} \wedge \hat{x}_{i_{2}} \in P_{N}$. For $j, j_{1}, j_{2} \in\{1,2, \ldots, r\}$, we have by equation (3.10) that

$$
\bar{\partial}_{j} f_{j_{1}, j_{2}}(x)=\bar{\partial}_{j} \mathbb{I}_{x=\hat{x}, j_{1}=i_{1}, j_{2}=i_{2}}=\mathbb{I}_{x=\hat{x}, j_{1}=i_{1}, j_{2}=i_{2}}-\mathbb{I}_{x=\hat{x}-\mathbf{e}_{j}, j_{1}=i_{1}, j_{2}=i_{2}} .
$$

Letting the co-derivative operate on $f$, we obtain

$$
\begin{aligned}
\delta f(x) & =\sum_{1 \leq j_{1}<\cdots<j_{2} \leq r} \sum_{i=1}^{2}(-1)^{i} \bar{\partial}_{i} f_{j_{1}, j_{2}}(x) d x_{j_{1}} \wedge \cdots \wedge d x_{j_{i-1}} \wedge d x_{j_{i+1}} \wedge \ldots \wedge d x_{j_{2}} \\
& =(-1)^{1} \bar{\partial}_{i} \mathbb{I}_{x=\hat{x}} d x_{i_{2}}+(-1)^{2} \bar{\partial}_{i} \mathbb{I}_{x=\hat{x}} d x_{i_{1}}=-\bar{\partial}_{i} \mathbb{I}_{x=\hat{x}} d x_{i_{2}}+\bar{\partial}_{i} \mathbb{I}_{x=\hat{x}} d x_{i_{1}} \\
& =-\mathbb{I}_{x=\hat{x}} d x_{i_{2}}+\mathbb{I}_{x-\mathbf{e}_{i_{1}}=\hat{x}} d x_{i_{2}}+\mathbb{I}_{x=\hat{x}} d x_{i_{1}}-\mathbb{I}_{x-\mathbf{e}_{i_{2}}=\hat{x}} d x_{i_{1}} .
\end{aligned}
$$

Thus,

$$
(\delta f)_{e}= \begin{cases}1 & \text { if } e \in \partial p_{0} \\ -1 & \text { if }-e \in \partial p_{0} \\ 0 & \text { otherwise }\end{cases}
$$

As a result, when $1 \leq j \leq n$ and $f$ is a 2-form, we obtain for an edge $e=d x_{j}$ that

$$
(\delta f)_{e}=(\delta f)_{j}(x)=\sum_{p \in \hat{\partial} c} f_{p}
$$

This result can be extended to general $k$-forms. For a $k$-form $f$ and a ( $k-1$ )-cell $c_{0}$, we have

$$
\begin{equation*}
(\delta f)_{c_{0}}=\sum_{c \in \hat{\partial} c_{0}} f_{c} . \tag{3.12}
\end{equation*}
$$

Last, the version of the Poincare lemma for the co-derivative is stated.
Lemma 3.24. (The Poincaré lemma for the co-derivative) Let $k \in\{1,2, \ldots, r-$ $1\}$ and $f$ be a $G$-valued $k$-form on $\mathbb{Z}^{r}$ which is zero outside a finite region and satisfies $\delta f=0$. Then there is a $(k+1)$-form $h$ such that $f=\delta h$. Moreover, if $f$ is equal to zero outside a box $B$, then there is a choice of $h$ that is equal to zero outside $B$.

For the proof, see [2, Lemma 2.7]. Notice that if a $k$-form is zero outside a finite region, the support of the $k$-form is contained in this finite region. This lemma will be applied when theory for oriented surfaces is discussed.

### 3.6 The Hodge dual

We call the lattice $\mathbb{Z}^{r}$ for the primal lattice and create a copy of it, which does not have its vertices at the same points as the primal lattice $\mathbb{Z}^{r}$. This copy, whose vertices are at the centres of the $r$-cells of the primal lattice, is called the dual lattice and denoted $* \mathbb{Z}^{r}$. Therefore, there is a bijection between the $r$-cells in the primal lattice and the 0-cells, i.e. vertices, in the dual lattice. Furthermore, we have a bijection between the set of $k$-cells of the primal lattice $\mathbb{Z}^{r}$ and the set of $(r-k)$-cells of the dual lattice $* \mathbb{Z}^{r}$, which we will define. For a cell $c$, the Hodge dual of the cell is denoted by $*$. The operator $*$ is called the Hodge star operator and is additive.

We aim to define the bijection between a $k$-cell and its Hodge dual. First, some theory for permutations of sets is given. The sign of a permutation $p$ is given by $\operatorname{sgn}(p)$ and takes the value 1 if the permutation is even and -1 if the permutation is odd. The sign can be calculated with the formula $\operatorname{sgn}(p)=(-1)^{m}$, where $m$ is the number of transpositions used to obtain the rearrangement from the given set.

We begin with a vertex $x$ and a $r$-cell $d x_{1} \wedge \cdots \wedge d x_{r}$ in the primal lattice $\mathbb{Z}^{r}$. The point at the centre of this $r$-cell is in the dual lattice given by $y:=*\left(d x_{1} \wedge\right.$ $\left.\cdots \wedge d x_{r}\right)$. The negatively oriented edges emerging from $y$ are denoted by $d y_{1}=$
$y-d \mathbf{e}_{1}, \ldots, d y_{r}=y-d \mathbf{e}_{r}$. Let $k \in\{0,1, \ldots, r\}$ and $1 \leq i_{1}<\ldots<i_{k} \leq r$ be given. Since $x \in \mathbb{Z}^{r}$, the $k$-cell $c=d x_{i_{1}} \wedge \ldots \wedge d x_{i_{k}}$ is also in the primal lattice $\mathbb{Z}^{r}$. Let $j_{1}, \ldots, j_{r-k}$ be an enumeration of the set $\{1,2, \ldots, r\} \backslash\left\{i_{1}, \ldots, i_{k}\right\}$. Consider the permutation that maps $(1,2, \ldots, r)$ to $\left(i_{1}, \ldots, i_{k}, j_{1}, \ldots, j_{r-k}\right)$. Then the Hodge dual of the $k$-cell $c$ is defined as

$$
\begin{equation*}
*\left(d x_{i_{1}} \wedge \cdots \wedge d x_{i_{k}}\right):=\operatorname{sgn}\left(i_{1}, \ldots, i_{k}, j_{1}, \ldots, j_{r-k}\right) d y_{j_{1}} \wedge \cdots \wedge d y_{j_{r-k}} \tag{3.13}
\end{equation*}
$$

Likewise, we can define the Hodge dual of a $r-k$-cell in the dual lattice by

$$
\begin{align*}
*\left(d y_{j_{1}} \wedge \cdots \wedge d y_{j_{r-k}}\right) & :=\operatorname{sgn}\left(j_{1}, \ldots, j_{r-k}, i_{1}, \ldots, i_{k}\right) d x_{i_{1}} \wedge \cdots \wedge d x_{i_{k}} \\
& =(-1)^{k(r-k)} \operatorname{sgn}\left(i_{1}, \ldots, i_{k}, j_{1}, \ldots, j_{r-k}\right) d x_{i_{1}} \wedge \cdots \wedge d x_{i_{k}} \tag{3.14}
\end{align*}
$$

We continue with some examples.
Example 3.25. Let the primal lattice be $\mathbb{Z}^{3}$. For a 3 -cell in the primal lattice, we have that $r-k=3-3=0$. Hence, the Hodge dual of a 3 -cell is a 0 -cell at the centre of the 3-cell. For a 2-cell in the primal lattice, the Hodge dual is a 1 -cell in the dual lattice. This is illustrated in the figure below. Furthermore, the Hodge dual of a 1-cell is a 2-cell and the Hodge dual of a 0 -cell is a 3-cell.


Figure 3.5: A 3-cell and a 2-cell in the primal lattice (the black cells) and their Hodge duals in the dual lattice (the blue cells).

Example 3.26. We calculate the Hodge dual of the 2 -cells $d x_{1} \wedge d x_{2}$ and $d x_{2} \wedge d x_{4}$ in $\mathbb{Z}^{4}$. In the first case, the permutation of $\{1,2,3,4\}$ is even. The permutation
of $\{2,4,1,3\}$ is odd, since the set $\{1,2,3,4\}$ is obtained from $\{2,4,1,3\}$ with three transpositions. Thus,

$$
\begin{aligned}
& *\left(d x_{1} \wedge d x_{2}\right)=\operatorname{sgn}(1,2,3,4)\left(d y_{3} \wedge d y_{4}\right)=d y_{3} \wedge d y_{4} \\
& *\left(d x_{2} \wedge d x_{4}\right)=\operatorname{sgn}(2,4,1,3)\left(d y_{1} \wedge d y_{3}\right)=-d y_{1} \wedge d y_{3}
\end{aligned}
$$

Example 3.27. We calculate the Hodge dual of the 3 -cell $d x_{1} \wedge d x_{3} \wedge d x_{5}$ in $\mathbb{Z}^{5}$. The permutation of $\{1,3,5,2,4\}$ is odd, since the set $\{1,2,3,4,5\}$ is obtained from $\{1,3,5,2,4\}$ with three transpositions. Thus,

$$
*\left(d x_{1} \wedge d x_{3} \wedge d x_{5}\right)=\operatorname{sgn}(1,3,5,2,4)\left(d y_{2} \wedge d y_{4}\right)=-d y_{2} \wedge d y_{4}
$$

Next is the definition for the Hodge dual of a box.
Definition 3.28. (Hodge dual $* B$ of a box B) For a box $B \in \mathbb{Z}^{r}$, the corresponding box in the dual lattice $* \mathbb{Z}^{r}$ is

$$
* B:=\left\{y \in * \mathbb{Z}^{r}: \exists x \in B \text { such that } y \text { is a corner in } * x\right\}
$$

Example 3.29. $A$ cube $B \in \mathbb{Z}^{2}$ is represented by the black cube in Figure 3.6, where the vertices in the primal lattice are given by the grey points and the vertices in the dual lattice by the blue crosses. Then $* B \in * \mathbb{Z}^{2}$ is given by the blue cube and $* * B \in \mathbb{Z}^{2}$ is given by the red cube. Assume that $B$ is of width $b$. Then we see from the figure that $* B$ is of width $b+1$ and $* * B$ of width $b+2$.


Figure 3.6: The Hodge dual of a cube in $\mathbb{Z}^{2}$.

Example 3.30. The box $B$ is defined as $B:=\left(\left[a_{1}, b_{1}\right] \times \ldots \times\left[a_{r}, b_{r}\right]\right) \cap \mathbb{Z}^{r}$. Then, the Hodge dual is

$$
* B=\left(\left[a_{1}-\frac{1}{2}, b_{1}+\frac{1}{2}\right] \times \ldots \times\left[a_{r}-\frac{1}{2}, b_{r}+\frac{1}{2}\right]\right) \cap * \mathbb{Z}^{r} .
$$

Taking the Hodge dual again, we obtain

$$
* * B=\left(\left[a_{1}-1, b_{1}+1\right] \times \ldots \times\left[a_{r}-1, b_{r}+1\right]\right) \cap \mathbb{Z}^{r} .
$$

Thus, the width of a side in $* B$ is 1 larger than its corresponding side in $B$ and the width of a side in $* * B$ is 2 larger than its corresponding side in $B$. Consequently, $B \subsetneq * * B$. For example, if $B$ is a cube of width $b$, then $* * B$ is a cube of width $b+2$ containing $B$.

Next, one lemma for the Hodge dual of a box is given without proof.
Lemma 3.31. Let $B$ be a box in $\mathbb{Z}^{r}$. Then a $k$-cell $c$ is outside $B$ if and only if $* c$ is either outside $* B$ or in the boundary of $* B$. If $c$ is a $k$-cell outside $B$ that contains a $(k-1)$-cell of $B$, then $* c$ belongs to the boundary of $* B$.

For the proof, see [2, Lemma 2.4]. We define the Hodge dual for $k$-forms.
Definition 3.32. (Hodge dual of a $G$-valued $k$-form) Given a $G$-valued $k$-form $f$ on $\mathbb{Z}^{r}$, the Hodge dual $* f$ of $f$ is an $(r-k)$-form defined as

$$
\begin{equation*}
* f(y):=\sum_{1 \leq i_{1}<\cdots<i_{k} \leq r} f_{i_{1}, \ldots, i_{k}}(x) \operatorname{sgn}\left(i_{1}, \ldots, i_{k}, j_{1}, \ldots, j_{r-k}\right) d y_{j_{1}} \wedge \cdots \wedge d_{y_{j_{r-k}}}, \tag{3.15}
\end{equation*}
$$

where $y=*\left(d x_{1} \wedge \cdots \wedge d x_{r}\right)$ and the sequence $j_{1}, \ldots, j_{r-k}$ depends on $i_{1}, \ldots, i_{k}$.
From this definition and (3.13) follows that

$$
* f(* c)=f(c)
$$

for a $k$-cell $c=d x_{1} \wedge \cdots \wedge d x_{r}$.
Example 3.33. We calculate the Hodge dual of the 2 -form $f=f_{12}(x) d x_{1} \wedge d x_{2}+$ $f_{34}(x) d x_{3} \wedge d x_{4}$ in $\mathbb{Z}^{4}$. By Definition 3.32, we have

$$
\begin{aligned}
* f(y) & =f_{12}(x) \operatorname{sgn}(1,2,3,4) d y_{3} \wedge d y_{4}+\operatorname{sgn}(3,4,1,2) f_{34}(x) d y_{1} \wedge d y_{2} \\
& =f_{12}(x) d y_{3} \wedge d y_{4}+f_{13}(x) d y_{2} \wedge d y_{4} .
\end{aligned}
$$

For the Hodge dual of a $k$-form holds the following formula:

$$
\begin{equation*}
*(* f)=(-1)^{k(r-k)} f . \tag{3.16}
\end{equation*}
$$

We prove this by calculating $*(* f)$ :

$$
\begin{aligned}
& *(* f) \stackrel{\sqrt{3.155}=}{=}\left(\sum_{1 \leq i_{1}<\cdots<i_{k} \leq r} f_{i_{1}, \ldots, i_{k}}(x) \operatorname{sgn}\left(i_{1}, \ldots, i_{k}, j_{1}, \ldots, j_{r-k}\right) d y_{j_{1}} \wedge \cdots \wedge d_{y_{j_{r-k}}}\right) \\
&=\sum_{1 \leq i_{1}<\cdots<i_{k} \leq r} f_{i_{1}, \ldots, i_{k}}(x) \operatorname{sgn}\left(i_{1}, \ldots, i_{k}, j_{1}, \ldots, j_{r-k}\right) *\left(d y_{j_{1}} \wedge \cdots \wedge d_{y_{j_{r-k}}}\right) \\
& \stackrel{3.14}{=} \sum_{1 \leq i_{1}<\cdots<i_{k} \leq r} f_{i_{1}, \ldots, i_{k}}(x) \operatorname{sgn}\left(i_{1}, \ldots, i_{k}, j_{1}, \ldots, j_{r-k}\right)(-1)^{k(r-k)} \\
& \cdot \operatorname{sgn}\left(i_{1}, \ldots, i_{k}, j_{1}, \ldots, j_{r-k}\right) d x_{i_{1}} \wedge \cdots \wedge d x_{i_{k}} \\
&=\sum_{1 \leq i_{1}<\cdots<i_{k} \leq r} f_{i_{1}, \ldots, i_{k}}(x)(-1)^{k(r-k)} d x_{i_{1}} \wedge \cdots \wedge d x_{i_{k}} \\
&=(-1)^{k(r-k)} \sum_{1 \leq i_{1}<\cdots<i_{k} \leq r} f_{i_{1}, \ldots, i_{k}}(x) d x_{i_{1}} \wedge \cdots \wedge d x_{i_{k}} \\
& \stackrel{\text { 3.1] }}{=}(-1)^{k(r-k)} f .
\end{aligned}
$$

The exterior derivative has different definitions on the primal and on the dual lattice.

Definition 3.34. (Exterior derivative on the dual cell lattice) The exterior derivative d on the dual cell lattice is defined by

$$
d f(y)=\sum_{1 \leq j_{1}<\cdots<j_{k} \leq r} \sum_{i=1}^{r} \bar{\partial}_{i} f_{j_{1}, \ldots, j_{k}}(y) d y_{i} \wedge\left(d y_{j_{1}} \wedge \ldots \wedge d y_{j_{k}}\right), \quad y \in * \mathbb{Z}^{r}
$$

Note that $\bar{\partial}_{i}$ is used instead of $\partial_{i}$, although the exterior derivative is denoted by the same symbol $d$ in both cases. Last, a lemma is given.

Lemma 3.35. For any $G$-valued $k$-form $f$ on $\mathbb{Z}^{r}$ and any $x \in \mathbb{Z}^{r}$,

$$
\delta f(x)=(-1)^{r(k+1)+1} *(d(* f(y)))
$$

where $y=*\left(d x_{1} \wedge \cdots \wedge d x_{r}\right)$ is the centre of the $r$-cell $\left(d x_{1} \wedge \cdots \wedge d x_{r}\right)$.
For the proof, see [2, Lemma 2.3].

## Chapter 4

## Vortices and oriented surfaces

In this chapter, vortices and oriented surfaces are discussed. Lemmas crucial for the proof of the main theorem are stated and proved. From now on, the lattice is always $\mathbb{Z}^{4}$, the group $G$ is a finite Abelian group and the representation $\rho$ of $G$ is fixed.

### 4.1 Vortices

We begin with defining vortices before discussing both decompositions and distributions of vortices. Some theory for minimal vortices will also be given. First, the definition for an irreducible plaquette configuration is given. This definition is necessary for the definition of a vortex.

Definition 4.1. (Irreducible plaquette configuration) A plaquette configuration $\omega \in \Sigma_{P_{N}}$ is irreducible if there does not exist a non-empty set $P \subsetneq \operatorname{supp} \omega$ such that $\left.\omega\right|_{P} \in \Sigma_{P_{N}}$.

Recall from the definition of $\Sigma_{P_{N}}$ that $d \omega=0$ for $\omega \in \Sigma_{P_{N}}$. Therefore, in the above definition $\left.\omega\right|_{P} \in \Sigma_{P_{N}}$ is equivalent to $d\left(\left.\omega\right|_{P}\right)=0$.

Definition 4.2. (Vortex) Let $\sigma \in \Sigma_{E_{N}}$. A non-trivial and irreducible plaquette configuration $\nu \in \Sigma_{P_{N}}$ is a vortex in $\sigma$ if $\left.(d \sigma)\right|_{\text {supp } \nu}=\nu$.

### 4.1.1 Vortex decompositions

The goal of this section is to prove that if a spin configuration $\sigma \in \Sigma_{E_{N}}$, then $d \sigma$ can be written as a sum of vortices with disjoint supports in $\sigma$. To prove this
claim, we need the results from the following two lemmas.
Lemma 4.3. Let $\omega \in \Sigma_{P_{N}}$ and let $\nu \in \Sigma_{P_{N}}$ be such that $\left.\omega\right|_{\text {supp } \nu}=\nu$. Then $\left.\omega\right|_{P_{N} \backslash \text { supp } \nu} \in \Sigma_{P_{N}}$.

Proof. We prove this by showing that $d\left(\left.\omega\right|_{P_{N} \backslash \text { supp } \nu}\right)=0$. Since $\omega \in \Sigma_{P_{N}}$ and $\nu \in \Sigma_{P_{N}}$, we have that $d \omega=0$ and $d \nu=0$. Since $\omega \in \Sigma_{P_{N}} \subseteq P_{N}$ and supp $\nu$, it follows that $\left.\omega\right|_{P_{N} \backslash \text { supp } \nu}=\left.\omega\right|_{P_{N}}-\left.\omega\right|_{\text {supp } \nu}=\omega-\nu$. Thus,

$$
d\left(\left.\omega\right|_{P_{N} \backslash \operatorname{supp} \nu}\right)=d(\omega-\nu)=d \omega-d \nu=0-0=0
$$

and $\left.\omega\right|_{P_{N} \backslash \text { supp } \nu} \in \Sigma_{P_{N}}$.
Lemma 4.4. Let $\omega \in \Sigma_{P_{N}}$. Then either $\omega$ is irreducible or there exist nontrivial $\nu^{(1)}, \nu^{(2)} \in \Sigma_{P_{N}}$ with disjoint supports contained in $\operatorname{supp} \omega$, such that $\omega=\nu^{(1)}+\nu^{(2)}$.

Proof. We assume that $\omega$ is reducible and then prove that there exists nontrivial $\nu^{(1)}, \nu^{(2)} \in \Sigma_{P_{N}}$ with disjoint supports contained in supp $\omega$ such that $\omega=\nu^{(1)}+\nu^{(2)}$. By Definition 4.1, there exists at least one non-empty set $P \subsetneq \operatorname{supp} \omega$ such that $\nu^{(1)}:=\left.\omega\right|_{P} \in \Sigma_{P_{N}}$. Consequently, $P=\operatorname{supp} \nu^{(1)}$ and

$$
\begin{equation*}
\left.\omega\right|_{\text {supp } \nu^{(1)}}=\nu^{(1)} . \tag{4.1}
\end{equation*}
$$

Therefore, $\nu^{(1)}$ fulfils the assumptions for Lemma 4.3 and

$$
\begin{equation*}
\nu^{(2)}:=\left.\omega\right|_{P_{N} \backslash \operatorname{supp} \nu^{(1)}} \in \Sigma_{P_{N}} . \tag{4.2}
\end{equation*}
$$

Since supp $\nu^{(1)}$ is non-empty, $\nu^{(1)}$ contains at least one 2-cell for which $\nu^{(1)} \neq 0$. Similarly, $\nu^{(2)}$ contains at least one 2 -cell for which $\nu^{(2)} \neq 0$. Thus, they are both non-trivial. Since supp $\nu^{(2)} \subseteq P_{N} \backslash \operatorname{supp} \nu^{(1)}$, the intersection of supp $\nu^{(1)}$ and supp $\nu^{(2)}$ must be empty. Hence, they have disjoint supports. Last, we conclude that

$$
\nu^{(1)}+\left.\nu^{(2)} \stackrel{\sqrt[4.11]{=}, \sqrt[4.22]{ }}{\omega}\right|_{\text {supp } \nu^{(1)}}+\left.\omega\right|_{P_{N} \backslash \operatorname{supp} \nu^{(1)}}=\left.\omega\right|_{P_{N}}=\omega .
$$

We now have everything necessary to prove and state the main result of this section.

Lemma 4.5. Let $\sigma \in \Sigma_{E_{N}}$. Then $d \sigma$ can be written as a sum of vortices in $\sigma$ with disjoint supports.

Proof. Since $\sigma \in \Sigma_{E_{N}}$ and $\left.d \sigma\right|_{\text {supp } d \sigma}=d \sigma$, it follows from Definition 4.2 that there are two cases to consider. First, if $d \omega$ is irreducible, then $d \sigma$ is a vortex. Second, we consider the case when $d \omega$ is not irreducible. Since $d \sigma \in \Sigma_{P_{N}}$, by Lemma 4.4 there exists non-trivial $\nu^{(1)}, \nu^{(2)} \in \Sigma_{P_{N}}$ with disjoint supports in supp $d \sigma$ for which $d \sigma=\nu^{(1)}+\nu^{(2)}$.

If $\nu^{(1)}$ and $\nu^{(2)}$ are both irreducible, we have a sum of two non-trivial vortices with disjoint supports. Otherwise, by Lemma 4.4, $\nu^{(1)}$ and/or $\nu^{(2)}$ can be decomposed further. This decomposition of 2-forms is repeated until we only have irreducible and non-trivial plaquette configurations $\nu^{(j)} \in \Sigma_{P_{N}}$. Since this decomposition cannot consist of more than $\left|P_{N}\right|<\infty$ parts, there must exist an $m \leq\left|P_{N}\right|<\infty$ for which $d \sigma$ can be written as a finite sum of $m$ vortices with disjoint supports, i.e. $d \sigma=\sum_{j=1}^{m} \nu^{(j)}$.

### 4.1.2 Minimal vortices

In this section, we define minimal vortices and prove that a minimal vortex $\nu$, whose support does not contain any boundary plaquettes of $P_{N}$, can be written using an edge in $E_{N}$ and an element $g \in G \backslash\{0\}$. To achieve this result, two lemmas are necessary. We begin with the definition for a minimal vortex.

Definition 4.6. (Minimal vortex) Let $\sigma \in \Sigma_{E_{N}}$. A vortex $\nu$ in $\sigma$ is minimal if $|\operatorname{supp} \nu|=12$.

The following lemma says that the support of a minimal vortex, which contains no boundary plaquettes of $P_{N}$, can be written as a union of the co-boundary of a positively oriented edge $e_{0}$ and the co-boundary of the negatively oriented edge $-e_{0}$.

Lemma 4.7. Let $\sigma \in \Sigma_{E_{N}}$ and let $\nu$ in $\sigma$ be a vortex for which supp $\nu$ contains no boundary plaquettes of $P_{N}$. Then $|\operatorname{supp} \nu| \geq 12$. If $|\operatorname{supp} \nu|=12$, then there exists an edge $e_{0} \in E_{N}$ such that $\operatorname{supp} \nu=\hat{\partial} e_{0} \cup \hat{\partial}\left(-e_{0}\right)$.

For a proof of this lemma, see [3, Lemma 3.4.6]. Before stating the next lemma about minimal vortices, the Bianchi Lemma is given.

Lemma 4.8. (Bianchi Lemma) If $\omega \in \sum_{P_{N}}$ and $c$ is an oriented 3-cell in $B_{N}$, then

$$
\begin{equation*}
\sum_{p \in \partial c} \omega_{p}=0 . \tag{4.3}
\end{equation*}
$$

Proof. Let $c$ be an oriented 3 -cell in $B_{N}$. Since $\omega \in \Sigma_{P_{N}}$, we have that

$$
\sum_{p \in \partial c} \omega_{p} \stackrel{3.7}{=}(d \omega)_{c}=0
$$

The following lemma is the main result of this section.
Lemma 4.9. Let $\sigma \in \Sigma_{E_{N}}$ and let $\nu$ be a minimal vortex in $\sigma$. If the support of $\nu$ contains no boundary plaquettes of $P_{N}$, then there exists an edge $d x_{j} \in E_{N}$ and $g \in G \backslash\{0\}$ such that for all $p \in P_{N}$, we have

$$
\nu_{p}=\left(d\left(g d x_{j}\right)\right)_{p} .
$$

Proof. By definition, $|\operatorname{supp} \nu|=12$. Since the support of $\nu$ contains no boundary plaquettes of $P_{N}$, it follows from Lemma 4.7 that there exists an edge $e_{0}=d x_{j} \in$ $E_{N}$ such that supp $\nu=\hat{\partial} e_{0} \cup \hat{\partial}\left(-e_{0}\right)$.

A 3-cell $c$ can not contain more than two plaquettes in $\hat{\partial} e_{0} \cup \hat{\partial}\left(-e_{0}\right)$. Hence, we have two plaquettes $p_{1}$ and $p_{2}$ such that $p_{1} \in \hat{\partial} e_{0}$ and $p_{2} \in \hat{\partial}\left(-e_{0}\right)$. By Lemma 4.8,

$$
\sum_{p \in \partial c} \nu_{p}=\nu_{p_{1}}+\nu_{p_{2}}=0 .
$$

Since the plaquettes are in the support of $\nu$, we have $\nu_{p_{1}} \neq 0 \neq \nu_{p_{2}}$. Hence, $\nu_{p_{1}}=g$ and $\nu_{p_{2}}=-g$ for a $g \in G \backslash\{0\}$. To conclude,

$$
\nu_{p}= \begin{cases}g & \text { if } p \in \hat{\partial} e_{0}, \\ -g & \text { if } p \in \hat{\partial}\left(-e_{0}\right)\end{cases}
$$

Since $p \in \hat{\partial}\left(-e_{0}\right)$ is equivalent to $-e_{0} \in \partial p$, we see from equation (3.5) that $\left(d\left(g d x_{j}\right)\right)_{p}$ can be written as

$$
\left(d\left(g d x_{j}\right)\right)_{p}= \begin{cases}g & \text { if } p \in \hat{\partial} e_{0} \\ -g & \text { if } p \in \hat{\partial}\left(-e_{0}\right)\end{cases}
$$

Thus, $\nu_{p}=\left(d\left(g d x_{j}\right)\right)_{p}$.

This lemma will be applied in the proof of Proposition 6.1.

### 4.1.3 Distribution of vortices

In this section, the focus is on distributions of vortices. With the help of two lemmas, we will state and prove Proposition 4.12 that is of great importance in Chapter 6. We begin with the first lemma, which gives us an upper bound for a useful probability.

Lemma 4.10. Let $\nu \in \Sigma_{P_{N}}$. Then

$$
\mu_{\beta, N}\left(\left\{\sigma \in \Sigma_{E_{N}}:\left.(d \sigma)\right|_{\operatorname{supp} \nu}=\nu\right\}\right) \leq \lambda(\beta)^{\mid \text {supp } \nu \mid}
$$

Proof. Let $P:=\operatorname{supp} \nu$ and define the sets $\mathcal{E}_{P}^{\nu}:=\left\{\omega \in \Sigma_{P_{N}}:\left.\omega\right|_{P}=\nu\right\}$ and $\mathcal{E}_{P}^{0}:=\left\{\omega \in \Sigma_{P_{N}}:\left.\omega\right|_{P}=0\right\}$. By the Poincare Lemma, there exists a $\sigma \in \Sigma_{E_{N}}$ such that $d \sigma=\omega$ if and only if $\omega \in \Sigma_{P_{N}}$. Therefore, by equation (3.9) we have that $\mu_{\beta, N}\left(\mathcal{E}_{P}^{\nu}\right)=\mu_{\beta, N}\left(\left\{\sigma \in \Sigma_{E_{N}}:\left.(d \sigma)\right|_{\text {supp } \nu}=\nu\right\}\right)$. We calculate an upper bound for the probability. Since $\mathcal{E}_{P}^{0} \subseteq \Sigma_{P_{N}}$ and $\phi_{\beta}$ is positive, we have

$$
\begin{align*}
\mu_{\beta, N}\left(\left\{\sigma \in \Sigma_{E_{N}}:\left.(d \sigma)\right|_{\text {supp } \nu}=\nu\right\}\right) & =\mu_{\beta, N}\left(\mathcal{E}_{P}^{\nu}\right)=\frac{\sum_{\omega \in \mathcal{E}_{P}^{\nu}} \prod_{p \in P_{N}} \phi_{\beta}\left(\omega_{p}\right)}{\sum_{\omega \in \Sigma_{P_{N}}} \prod_{p \in P_{N}} \phi_{\beta}\left(\omega_{p}\right)}  \tag{4.4}\\
& \leq \frac{\sum_{\omega \in \mathcal{E}_{P}^{\nu}} \prod_{p \in P_{N}} \phi_{\beta}\left(\omega_{p}\right)}{\sum_{\omega \in \mathcal{E}_{P}^{0}} \prod_{p \in P_{N}} \phi_{\beta}\left(\omega_{p}\right)} .
\end{align*}
$$

We aim to compare the denominator to a sum over $\omega \in \mathcal{E}_{P}^{\nu}$. From the definitions of the sets $\mathcal{E}_{P}^{\nu}$ and $\mathcal{E}_{P}^{0}$ follows that

$$
\sum_{\omega \in \mathcal{E}_{P}^{0}} \prod_{p \in P_{N}} \phi_{\beta}\left(\omega_{p}\right)=\sum_{\omega \in \mathcal{E}_{P}^{\nu}} \prod_{p \in P_{N}} \phi_{\beta}\left((\omega-\nu)_{p}\right) .
$$

Since the mapping $\omega \rightarrow \omega-\nu$ from $\mathcal{E}_{P}^{\nu}$ to $\mathcal{E}_{P}^{0}$ is bijective, this gives

$$
\begin{align*}
& \frac{\sum_{\omega \in \mathcal{E}_{P}^{\nu}} \prod_{p \in P_{N}} \phi_{\beta}\left(\omega_{p}\right)}{\sum_{\omega \in \mathcal{E}_{P}^{0}} \prod_{p \in P_{N}} \phi_{\beta}\left(\omega_{p}\right)}=\frac{\sum_{\omega \in \mathcal{E}_{P}^{\nu}} \prod_{p \in P_{N}} \phi_{\beta}\left(\omega_{p}\right)}{\sum_{\omega \in \mathcal{E}_{P}^{\nu}} \prod_{p \in P_{N}} \phi_{\beta}\left((\omega-\nu)_{p}\right)}  \tag{4.5}\\
& =\frac{\sum_{\omega \in \mathcal{E}_{P}^{\nu}} \prod_{p \in P} \phi_{\beta}\left(\omega_{p}\right) \prod_{p \in P_{N} \backslash P} \phi_{\beta}\left(\omega_{p}\right)}{\sum_{\omega \in \mathcal{E}_{P}^{\nu}} \prod_{p \in P} \phi_{\beta}\left((\omega-\nu)_{p}\right) \prod_{p \in P_{N} \backslash P} \phi_{\beta}\left((\omega-\nu)_{p}\right)} .
\end{align*}
$$

To simplify it further, notice that $\nu_{p}=0$ when $p \in P_{N} \backslash P$. Therefore,

$$
\begin{equation*}
\phi_{\beta}\left((\omega-\nu)_{p}\right)=\phi_{\beta}\left(\omega_{p}\right) \quad \text { for } p \in P_{N} \backslash P . \tag{4.6}
\end{equation*}
$$

For a plaquette $p \in P$ and a plaquette configuration $\omega \in \mathcal{E}_{P}^{\nu}$, we have by definition that $\omega_{p}=\nu_{p}$. Thus,

$$
\begin{equation*}
(\omega-\nu)_{p}=\omega_{p}-\nu_{p}=0 \quad \text { for } p \in P \text { and } \omega \in \mathcal{E}_{P}^{\nu} . \tag{4.7}
\end{equation*}
$$

Combining the previous conclusions, we obtain

$$
\begin{aligned}
& \mu_{\beta, N}\left(\mathcal{E}_{P}^{\sqrt{4.44}}\right)^{4.5]} \frac{\sum_{\omega \in \mathcal{E}_{P}^{\nu}} \prod_{p \in P} \phi_{\beta}\left(\omega_{p}\right) \prod_{p \in P_{N} \backslash P} \phi_{\beta}\left(\omega_{p}\right)}{\sum_{\omega \in \mathcal{E}_{P}^{\nu}} \prod_{p \in P} \phi_{\beta}\left((\omega-\nu)_{p}\right) \prod_{p \in P_{N} \backslash P} \phi_{\beta}\left((\omega-\nu)_{p}\right)} \\
& \text { [4.6], (4.7] } \frac{\sum_{\omega \in \mathcal{E}_{P}^{\nu}} \prod_{p \in P} \phi_{\beta}\left(\omega_{p}\right) \prod_{p \in P_{N} \backslash P} \phi_{\beta}\left(\omega_{p}\right)}{\sum_{\omega \in \mathcal{E}_{P}^{\nu}} \prod_{p \in P} \phi_{\beta}(0) \prod_{p \in P_{N} \backslash P} \phi_{\beta}\left(\omega_{p}\right)} \\
& =\frac{\prod_{p \in P} \phi_{\beta}\left(\nu_{p}\right) \sum_{\omega \in \mathcal{E}_{P}^{\nu}} \prod_{p \in P_{N} \backslash P} \phi_{\beta}\left(\omega_{p}\right)}{\prod_{p \in P} \phi_{\beta}(0) \sum_{\omega \in \mathcal{E}_{P}^{\nu}} \prod_{p \in P_{N} \backslash P} \phi_{\beta}\left(\omega_{p}\right)} \\
& =\frac{\prod_{p \in P} \phi_{\beta}\left(\nu_{p}\right)}{\prod_{p \in P} \phi_{\beta}(0)}=\prod_{p \in \operatorname{supp} \nu} \frac{\phi_{\beta}\left(\nu_{p}\right)}{\phi_{\beta}(0)} \\
& \stackrel{[2.11}{\leq} \prod_{p \in \operatorname{supp} \nu} \lambda(\beta)=\lambda(\beta)^{|\operatorname{supp} \nu|} \text {. }
\end{aligned}
$$

Thus, $\mu_{\beta, N}\left(\left\{\sigma \in \Sigma_{E_{N}}:\left.(d \sigma)\right|_{\text {supp } \nu}=\nu\right\}\right) \leq \lambda(\beta)^{|\operatorname{supp} \nu|}$.
The following lemma is about how many irreducible plaquette configurations there can exist.

Lemma 4.11. For each $p_{0} \in P_{N}$ and each $m \geq 6$, there are at most $5^{m-1}(|G|-$ 1) ${ }^{m}$ irreducible $\nu \in \Sigma_{P_{N}}$ with $p_{0} \in \operatorname{supp} \nu$ and $|\operatorname{supp} \nu| \geq 2 m$.

Proof. We construct an injective map from the set of irreducible 2-forms in $\Sigma_{P_{N}}$ with $p_{0} \in \operatorname{supp} \nu$ and $|\operatorname{supp} \nu| \geq 2 m$ to the set of sequences $\nu^{(1)}, \nu^{(2)}, \ldots \nu^{(m)}$ of $G$-valued 2-forms on $P_{N}$.

Let $\nu \in \Sigma_{P_{N}}$ be irreducible, $p_{0} \in \operatorname{supp} \nu$ and $|\operatorname{supp} \nu|=2 m$. Define $\nu^{(0)}:=$ $0 \in \Sigma_{P_{N}}, p_{1}:=p_{0}$ and for $p \in P_{N}$ :

$$
\nu_{p}^{(1)}:= \begin{cases}\nu_{p_{1}} & \text { if } p=p_{1} \\ -\nu_{p_{1}} & \text { if } p=-p_{1} \\ 0 & \text { else. }\end{cases}
$$

Let $k \in\{1,2, \ldots, m\}$ and assume that the 2-forms $\nu^{(1)}, \nu^{(2)}, \ldots, \nu^{(k)}$ are given such that for each $j \in\{1,2, \ldots, k\}$, the two following statements hold:
(a) supp $\nu^{(j)} \backslash \operatorname{supp} \nu^{(j-1)}=\left\{p_{j},-p_{j}\right\}$ for some $p_{j} \in P_{N}$. This implies that $\left|\operatorname{supp} \nu^{(j)}\right|=2 j$.
(b) $\left.\nu\right|_{\text {supp } \nu^{(j)}}=\nu^{(j)}$.

We check that (a) and (b) hold if $k=1$. Since the sequence contains only one element, we only have to check the conditions for $\nu^{(1)}$. (a) supp $\nu^{(1)} \backslash \operatorname{supp} \nu^{(0)}=$ $\left\{p_{1},-p_{1}\right\} \backslash\{\varnothing\}=\left\{p_{1},-p_{1}\right\}$ and (b) $\left.\nu\right|_{\text {supp } \nu^{(1)}}=\nu^{(1)}$.

We prove that $\nu^{(m)}=\nu$ by considering two different cases: $d \nu^{(k)}=0$ and $d \nu^{(k)} \not \equiv 0$. First, we assume that there exists an sequence satisfying (a) and (b) for which $d \nu^{(k)}=0$. From (b) follows that $\left.\nu\right|_{\text {supp } \nu^{(k)}}=\nu^{(k)} \in \Sigma_{P_{N}}$. Since $\nu$ is irreducible, there exists no non-empty set $P \subsetneq \operatorname{supp} \nu$ such that $\left.\nu\right|_{P} \in \Sigma_{P_{N}}$. Since it followed from (a) that supp $\nu^{(k)}$ is non-empty, we have $\nu^{(k)}=\nu$. Since $\left|\operatorname{supp} \nu^{(k)}\right|=2 k$, we must have

$$
2 k=\left|\operatorname{supp} \nu^{(k)}\right|=|\operatorname{supp} \nu|=2 m .
$$

Thus, $k=m$ and $\nu^{(m)}=\nu$.
Consider the second case, where we assume that the sequence is such that $d \nu^{(k)} \not \equiv 0$. Since we had $k=m$ in the first case, we now have that $k<m$. Therefore, we expand the sequence by defining $\nu^{(k+1)}$ such that (a) and (b) hold. Since $d \nu^{(k)} \not \equiv 0$, there exists an oriented 3-cell $c \in B_{N}$ such that $\left(d \nu^{(k)}\right)_{c} \neq 0$. We assume that both the 3 -cells in the set $B_{N}$ and the arbitrary total ordering of the plaquettes are given. Denote the first 3-cell for which $\left(d \nu^{(k)}\right)_{c} \neq 0$ by $c_{k+1}$. Furthermore, $(d \nu)_{c_{k+1}}=0$. Hence, there exists at least one plaquette $p \in \partial_{c_{k+1}} \backslash \operatorname{supp} \nu^{(k)}$. We denote the first plaquette for which $p \in \partial_{c_{k+1}} \backslash \operatorname{supp} \nu^{(k)}$ by $p_{k+1}$ and let

$$
\nu_{p}^{(k+1)}:= \begin{cases}\nu_{p_{(k+1)}} & \text { if } p=p_{k+1} \\ -\nu_{p_{(k+1)}} & \text { if } p=-p_{k+1}, \\ \nu_{p}^{(k)} & \text { otherwise } .\end{cases}
$$

We saw earlier that (a) and (b) holds for $k=1$. We assume that $\nu^{(1)}, \ldots, \nu^{(k)}$ satisfies both (a) and (b), then they also hold for $\nu^{(k+1)}$. Therefore, (a) and (b) hold for all $\nu^{(k)}$, where $k \in\{1,2, \ldots, m\}$. Since it follows from a) that $\left|\operatorname{supp} \nu^{(m)}\right|=$ $2 m$, from (b) that $\left.\nu\right|_{\operatorname{supp} \nu^{(m)}}=\nu^{(m)}$ and we assumed that $|\operatorname{supp} \nu|=2 m$, we have that $\nu^{(m)}=\nu$. Therefore, the injective map is constructed.

Since this mapping is injective, we can derive the upper bound for the set of sequences instead of the irreducible vortices with $p_{0} \in \operatorname{supp} \nu$ and $|\operatorname{supp} \nu| \geq$
$2 m$. Therefore, we calculate how many such sequences $\left(\nu^{(1)}, \nu^{(2)}, \ldots, \nu^{(m)}\right)$ there are. Since $\nu_{p_{k}} \in G \backslash\{0\}$, there is for each $k \in\{1,2, \ldots, m\}|G|-1$ choices for $\nu_{p_{k}}$. For each $k \in\{1,2, \ldots, m-1\}$, there is a given 3 -cell $c_{k+1}$ for which we have at most 5 choices for $p_{k+1}$ for each $k \in\{1,2, \ldots, m-1\}$. Combining this, we have at most $5^{m-1}(|G|-1)^{m}$ sequences $\left(\nu^{(1)}, \nu^{(2)}, \ldots, \nu^{(m)}\right)$. Hence, there exists at most $5^{m-1}(|G|-1)^{m}$ irreducible vortices $p_{0} \in \operatorname{supp} \nu$ and $|\operatorname{supp} \nu| \geq 2 m$.

Recall the definition of $\lambda(\beta)$. Since $\rho(0)=1$ and the largest value $\Re \rho(g)=$ $\cos \left(\frac{g \cdot 2 \pi m}{n}\right)$ can take for $g \neq 0$ is $\cos (2 \pi / n)$, it follows that $\lambda(\beta)$ can be written as

$$
\begin{aligned}
\lambda(\beta) & :=\max _{g \in G \backslash\{0\}} \frac{e^{\beta \Re(\rho(g))}}{e^{\beta \Re((\rho(0)))}}=\max _{g \in G \backslash 0\}} e^{\beta \Re(\rho(g)-\rho(0))} \\
& =\max _{g \in G \backslash\{0\}} e^{\beta \Re(\rho(g)-1)}=e^{\beta(\cos (2 \pi / n)-1)} .
\end{aligned}
$$

Hence,

$$
\lim _{\beta \rightarrow \infty} \lambda(\beta)=\lim _{\beta \rightarrow \infty} e^{\beta(\cos (2 \pi / n)-1)}=0
$$

Thus, there always exists a $\beta_{0}$ for which $5(|G|-1) \lambda(\beta)^{2}<1$ for $\beta>\beta_{0}$. Finally, we have everything necessary to state and prove the proposition.

Proposition 4.12. Fix any $\beta_{0}>0$ such that $5(|G|-1) \lambda(\beta)^{2}<1$ for all $\beta>\beta_{0}$. Fix $p_{0} \in P_{N}$ and $M \geq 6$. Then

$$
\begin{aligned}
& \mu_{\beta, N}\left(\left\{\sigma \in \Sigma_{E_{N}}: \exists \text { a vortex } \nu \text { in } \sigma \text { with } p_{0} \in \operatorname{supp} \nu \text { and }|\operatorname{supp} \nu| \geq 2 M\right\}\right) \\
& \quad \leq C_{0}^{(M)} \lambda(\beta)^{2 M}
\end{aligned}
$$

for all $\beta>\beta_{0}$, where

$$
\begin{equation*}
C_{0}^{(M)}:=\frac{5^{M}(|G|-1)^{M}}{1-5(|G|-1) \lambda(\beta)^{2}} . \tag{4.8}
\end{equation*}
$$

Proof. Let $m \in \mathbb{Z}_{+}$and $p_{0} \in P_{N}$. For $\nu \in \Sigma_{P_{N}}$, we have

$$
\mu_{\beta, N}\left(\left\{\sigma \in \Sigma_{E_{N}}:\left.(d \sigma)\right|_{\operatorname{supp} \nu}=\nu\right\}\right) \stackrel{\sqrt{4.10}}{\leq}(\lambda(\beta))^{|\operatorname{supp} \nu|}=\lambda(\beta)^{2 m} .
$$

From Lemma 4.11 follows that there are at most $5^{m-1}(|G|-1)^{m}$ irreducible plaquette configurations $\nu \in \Sigma_{P_{N}}$, such that $p_{0} \in \operatorname{supp} \nu$ and $|\operatorname{supp} \nu|=2 m$. Therefore,

$$
\begin{aligned}
& \mu_{\beta, N}\left(\left\{\sigma \in \Sigma_{E_{N}}: \exists \text { a vortex } \nu \text { in } \sigma \text { with } p_{0} \in \operatorname{supp} \nu \text { and }|\operatorname{supp} \nu|=2 m\right\}\right) \\
& \leq 5^{m-1}(|G|-1)^{m} \lambda(\beta)^{2 m} \leq(5(|G|-1))^{m} \lambda(\beta)^{2 m}
\end{aligned}
$$

We sum over every $m \geq M$ :

$$
\begin{aligned}
& \mu_{\beta, N}\left(\left\{\sigma \in \Sigma_{E_{N}}: \exists \text { a vortex } \nu \text { in } \sigma \text { with } p_{0} \in \operatorname{supp} \nu \text { and } \mid \text { supp } \nu \mid \geq 2 M\right)\right. \\
& =\sum_{m=M}^{\infty} \mu_{\beta, N}\left(\left\{\sigma \in \Sigma_{E_{N}}: \exists \text { a vortex } \nu \text { in } \sigma \text { with } p_{0} \in \operatorname{supp} \nu\right.\right. \text { and } \\
& \quad|\operatorname{supp} \nu|=2 m\}) \\
& \leq \sum_{m=M}^{\infty}(5(|G|-1))^{m} \lambda(\beta)^{2 m}=\frac{(5(|G|-1))^{M} \lambda(\beta)^{2 M}}{1-5(|G|-1) \lambda(\beta)^{2}}=C_{0}^{(M)} \lambda(\beta)^{2 M} .
\end{aligned}
$$

The second to last step follows from the sum being a geometric sum, which converges when $(5(|G|-1)) \lambda(\beta)^{2}<1$.

### 4.2 Oriented surfaces

This section focuses on oriented surfaces. Their relation to simple loops is given by the second lemma. First, the definitions for an oriented surface and its boundary are given.

Definition 4.13. (Oriented surface) $A \mathbb{Z}$-valued 2 -form $q$ on $P_{N}$ is an oriented surface if we have for every $e \in E_{N}$ that

$$
\begin{equation*}
(\delta q)_{e}=\sum_{p \in \hat{\partial} e} q_{p} \in\{-1,0,1\} . \tag{4.9}
\end{equation*}
$$

If a 2-form $q$ is such that $\left|q_{p}\right|=1$ for all $p \in P_{N}$, then the oriented surface is a collection of oriented plaquettes, where every plaquette shares at least one edge with another plaquette in the collection. An example of an oriented surface and its boundary is illustrated in Figure 4.1.


Figure 4.1: Example of an oriented surface and its boundary when $\left|q_{p}\right|=1$.

Definition 4.14. (The boundary of an oriented surface) Let $q$ be an oriented surface. Then

$$
B_{q}:=\left\{e \in E_{N}: \sum_{p \in \hat{\partial} e} q_{p}=1\right\}
$$

is the boundary of $q$.
Before stating the first lemma, we define for an oriented surface $q$ that

$$
q_{p}^{+}:= \begin{cases}\max \left(q_{p}, 0\right) & \text { for } p \in P_{N}  \tag{4.10}\\ 0 & \text { for } p \notin P_{N}\end{cases}
$$

Lemma 4.15. Let $q$ be an oriented surface with boundary $B_{q}$. Then,

$$
\sum_{p \in P_{N}} q_{p}^{+}(d \sigma)_{p}=\sum_{e \in B_{q}} \sigma_{e}, \quad \text { for all } \sigma \in \Sigma_{E_{N}}
$$

Note that this lemma is a discrete version of Stokes theorem.
Proof. Let $\sigma \in \Sigma_{E_{N}}$ and $p \in P_{N}$. Recall that the set $E_{N}^{+}$only contains the positively oriented edges in the set $E_{N}$ and that the set $E_{N}^{-}$only contains the negatively oriented edges. Furthermore,

$$
\begin{equation*}
\sigma_{e}=-\sigma_{-e} . \tag{4.11}
\end{equation*}
$$

From equation 3.7) follows that $(d \sigma)_{p}=\sum_{e \in \partial p} \sigma_{e}$. Therefore,

$$
\begin{aligned}
\sum_{p \in P_{N}} q_{p}^{+}(d \sigma)_{p} & =\sum_{p \in P_{N}} q_{p}^{+} \sum_{e \in \partial \partial_{p}} \sigma_{e} \\
& =\sum_{e \in E_{N}} \sigma_{e} \sum_{p \in(\hat{\partial} e) \cap P_{N}} q_{p}^{+} \\
& \stackrel{\boxed{4.10}}{=} \sum_{e \in E_{N}} \sigma_{e} \sum_{p \in \hat{\partial} e} q_{p}^{+} \\
& =\sum_{e \in E_{N}^{+}} \sigma_{e} \sum_{p \in \hat{\partial} e} q_{p}^{+}+\sum_{e \in E_{N}^{-}} \sigma_{e} \sum_{p \in \hat{\partial} e} q_{p}^{+} \\
& \stackrel{\boxed{4.11]}}{=} \sum_{e \in E_{N}^{+}} \sigma_{e} \sum_{p \in \hat{\partial} e} q_{p}^{+}-\sum_{e \in E_{N}^{+}} \sigma_{e} \sum_{p \in \hat{\partial}(-e)} q_{p}^{+} \\
& =\sum_{e \in E_{N}^{+}} \sigma_{e}\left(\sum_{p \in \hat{\partial} e} q_{p}^{+}-\sum_{p \in \hat{\partial}(-e)} q_{p}^{+}\right)
\end{aligned}
$$

Recall from Example 3.16 that $\hat{\partial} e \cap \hat{\partial}(-e)=\varnothing$. Moreover, if $p \in \hat{\partial}(-e)$, then $-p \in \hat{\partial} e$. Since $q$ is a 2 -form, we have for $p \in \hat{\partial} e$ that $q_{-p}=-q_{p}$. Therefore,

$$
\begin{equation*}
q_{p}^{+}-q_{-p}^{+}=\max \left(q_{p}, 0\right)-\max \left(q_{-p}, 0\right)=\max \left(q_{p}, 0\right)-\max \left(-q_{p}, 0\right)=q_{p} \tag{4.12}
\end{equation*}
$$

and

$$
\begin{aligned}
\sum_{e \in E_{N}^{+}} \sigma_{e}\left(\sum_{p \in \hat{\partial} e} q_{p}^{+}-\sum_{p \in \hat{\partial}(-e)} q_{p}^{+}\right) & =\sum_{e \in E_{N}^{+}} \sigma_{e}\left(\sum_{p \in \hat{\partial} e} q_{p}^{+}-\sum_{p \in \hat{\partial} e} q_{-p}^{+}\right) \\
& \stackrel{4.12]}{=} \sum_{e \in E_{N}^{+}} \sigma_{e} \sum_{p \in \hat{\partial} e} q_{p} .
\end{aligned}
$$

Furthermore,

$$
\begin{equation*}
\sigma_{e} \sum_{p \in \hat{\partial} e} q_{p}=-\sigma_{-e} \sum_{p \in \hat{\partial} e}\left(-q_{-p}\right)=\sigma_{-e} \sum_{p \in \hat{\partial} e} q_{-p}=\sigma_{-e} \sum_{-p \in \hat{\partial} e} q_{p}=\sigma_{-e} \sum_{p \in \hat{\partial}(-e)} q_{p} . \tag{4.13}
\end{equation*}
$$

Therefore,

$$
\begin{aligned}
& \sum_{e \in E_{N}^{+}} \sigma_{e} \sum_{p \in \hat{\partial} e} q_{p}^{\underline{\sqrt[{[4.13}]]{=}}} \sum_{e \in E_{N}^{+}}\left(\frac{1}{2}\left(\sigma_{e} \sum_{p \in \hat{\partial} e} q_{p}+\sigma_{-e} \sum_{p \in \hat{\partial}(-e)} q_{p}\right)\right)=\frac{1}{2} \sum_{e \in E_{N}} \sigma_{e} \sum_{p \in \hat{\partial} e} q_{p} \\
& \text { [4.9) } \frac{1}{2} \sum_{\substack{e \in E_{N} \\
\sum_{p \in \hat{\partial}_{e} e} q_{p}=-1}} \sigma_{e} \sum_{p \in \hat{\partial} e} q_{p}+\frac{1}{2} \sum_{\substack{e \in E_{N}: \\
\sum_{p \in \hat{\partial}_{e} q_{p}=0}}} \sigma_{e} \sum_{p \in \hat{\partial} e} q_{p} \\
& +\frac{1}{2} \sum_{\substack{e \in E_{N}: \\
\sum_{p \in \hat{\partial}_{e} q_{P}=1}}} \sigma_{e} \sum_{p \in \hat{\partial} e} q_{p} \\
& =\frac{1}{2} \sum_{\substack{e \in E_{N}: \\
\sum_{p \in \hat{e} e}^{q_{p}=-1}}} \sigma_{e} \sum_{p \in \hat{\partial} e} q_{p}+\frac{1}{2} \sum_{\substack{e \in E_{N}: \\
\sum_{p \in \hat{\partial}_{e}} q_{p}=1}} \sigma_{e} \sum_{p \in \hat{\partial} e} q_{p} \\
& =\frac{1}{2} \sum_{\substack{e \in E_{N}: \\
\sum_{p \in \hat{\partial} e} q_{p}=1}} \sigma_{e} \sum_{p \in \hat{\partial}_{e}} q_{p}+\frac{1}{2} \sum_{\substack{e \in E_{N}: \\
\sum_{p \in \hat{e}_{e} q_{p}=1}}} \sigma_{e} \sum_{p \in \hat{\partial}_{e}} q_{p} \\
& =\sum_{e \in B_{q}} \sigma_{e} \sum_{p \in \hat{\partial} e} q_{p}=\sum_{e \in B_{q}} \sigma_{e} .
\end{aligned}
$$

Thus,

$$
\sum_{p \in P_{N}} q_{p}^{+}(d \sigma)_{p}=\sum_{e \in B_{q}} \sigma_{e}
$$

This lemma is applied in the final proof of this section. The next lemma says that a simple loop is the boundary of an oriented surface. This is illustrated in Figure 4.2.


Figure 4.2: A simple loop $\gamma$ and an oriented surface, whose boundary is $\gamma$.

Lemma 4.16. Let $\gamma$ in $B_{N}$ be a simple loop contained in the cube $B \subseteq B_{N}$. Then there exists an oriented surface $q$, whose support is in $B$, such that $\gamma$ is the boundary of $q$.

Proof. Let $\gamma$ be a given simple loop contained in the cube $B$. For each $e \in E_{N}$, define the 1-form $f^{\gamma}$ as

$$
f_{e}^{\gamma}:= \begin{cases}1 & \text { if } e \in \gamma \\ -1 & \text { if }-e \in \gamma \\ 0 & \text { otherwise }\end{cases}
$$

By definition, the cube $B$ contains the support of $f^{\gamma}$. For the co-derivative of $f^{\gamma}$, we have that $\delta f_{e}^{\gamma}=0$. Hence, the assumptions of Lemma 3.24 are fulfilled for $f^{\gamma}$. Therefore, there exists a 2 -form $q^{\gamma}$ on $P_{N}$, with support only in $B$, for which

$$
\begin{equation*}
f^{\gamma}=\delta q^{\gamma} \tag{4.14}
\end{equation*}
$$

From equation (3.12) follows that

$$
\begin{equation*}
\left(\delta q^{\gamma}\right)_{e}=\sum_{p \in \hat{\partial} e} q_{p}^{\gamma} \quad \text { for } e \in E_{N} \tag{4.15}
\end{equation*}
$$

By combining these conclusions, we obtain

$$
\sum_{p \in \hat{\partial} e} q_{p}^{\gamma} \stackrel{\sqrt[4.15)]{=}}{=}\left(\delta q^{\gamma}\right)_{e} \stackrel{\sqrt{4.14]}}{=} f_{e}^{\gamma} \in\{-1,0,1\}
$$

Thus, $q^{\gamma}$ is an oriented surface with support contained in $B$ and $\gamma$ is the boundary of $q^{\gamma}$.

This lemma is applied in the proof of Proposition 6.1. Next, we define internal plaquettes and edges.

Definition 4.17. (Internal plaquette) Let $q$ be an oriented surface and let $\hat{p} \in$ supp $q$. A plaquette $\hat{p}$ is an internal plaquette of $q$ if

$$
\sum_{p \in \hat{\partial} e} q_{p}=0 \quad \text { for each } e \in \partial \hat{p}
$$

We derive another definition of internal plaquettes. If $p$ is not an internal plaquette of the oriented surface $q$, then $\sum_{p \in \hat{\partial} e} q_{p} \in\{-1,1\}$. For edges in the boundary of $q$, we have $\sum_{p \in \hat{\partial} e} q_{p}=1$. Since $q_{-p}=-q_{p}$, we obtain

$$
\sum_{p \in \hat{\partial} e} q_{p}=-\sum_{p \in \hat{\partial} e} q_{-p}=-\sum_{-p \in \hat{\partial} e} q_{p}
$$

and

$$
\sum_{p \in \hat{\partial} e} q_{p}=-1 \Leftrightarrow \sum_{-p \in \hat{\partial} e} q_{p}=1
$$

Recall that $p \in \hat{\partial} e$ is equivalent to $e \in \partial p$. Therefore, a plaquette $p \in \operatorname{supp} q$ is internal if

$$
\begin{equation*}
(\partial p \cup \partial(-p)) \cap B_{q}=\varnothing \tag{4.16}
\end{equation*}
$$

Definition 4.18. (Internal edge) An edge $e \in E_{N}$ is an internal edge of an oriented surface $q$ if there exists a plaquette $p \in \operatorname{supp} q$, such that $e \in \partial p$ and neither $e$ nor $-e$ is in the boundary of $q$.

Last, we have the final lemma of this chapter.
Lemma 4.19. Let $\sigma \in \Sigma_{E_{N}}$ and let $\nu \in \Sigma_{P_{N}}$ be a vortex in $\sigma$. Let $q$ be an oriented surface. If there exists a box $B$, which contains the support of $\nu$, for which the intersection $(* * B) \cap \operatorname{supp} q$ only contains internal plaquettes of $q$, then

$$
\begin{equation*}
\sum_{p \in P_{N}} q_{p}^{+} \nu_{p}=0 \tag{4.17}
\end{equation*}
$$

Proof. Since $\nu$ is a 2 -form and the lattice is $\mathbb{Z}^{4}$, it follows from equation (3.16) that

$$
\begin{equation*}
*(* \nu)=(-1)^{2(4-2)} \nu=\nu . \tag{4.18}
\end{equation*}
$$

Since $\nu \in \Sigma_{P_{N}}$, we have $d \nu=0$. By Lemma 3.35, the co-derivative of $* \nu$ is

$$
\delta(* \nu)=(-1)^{4(2+1)+1} *(d(*(* \nu))) \stackrel{4.18}{=}-* d \nu=-1 \cdot 0=0 .
$$

Assume that the box $B$ contains supp $\nu$. By Lemma 3.31, we have that the 2 -form $* \nu$ has no support outside $* B$. Thus, the assumptions of Lemma 3.24 are fulfilled for $* \nu$. Therefore, there exists a 3 -form $g$ without support outside $* B$, such that $* \nu=-\delta g$. By Lemma 3.35, we have

$$
\begin{equation*}
\delta g=(-1)^{4(3+1)+1} *(d(* g))=(-1) * d(* g) . \tag{4.19}
\end{equation*}
$$

Since $* g$ is a 1 -form, $d(* g)$ is a 2 -form and

$$
\begin{equation*}
*(*(d(* g))) \stackrel{\sqrt{3.16}}{=}-1^{2(4-2)}(d(* g))=d(* g) . \tag{4.20}
\end{equation*}
$$

Thus,

$$
\begin{equation*}
\nu \stackrel{\sqrt[4.18]]{=}}{=} *(* \nu)=*(-\delta g) \stackrel{\sqrt[4.19]{=}}{=} *(* d(* g)) \stackrel{\sqrt{4.20}}{=} d(* g) . \tag{4.21}
\end{equation*}
$$

Since the support of $q$ is finite, we have that $\sum_{p \in P_{N}} q_{p}^{+} \nu_{p}$ is well defined. Let $B_{q}$ be the boundary of $q$. From Lemma 4.15 follows that

$$
\sum_{p \in P_{N}} q_{p}^{+} \nu_{p} \stackrel{4.21]}{=} \sum_{p \in P_{N}} q_{p}^{+}(d(* g))_{p}=\sum_{e \in B_{q}}(* g)_{e} .
$$

A $k$-form $f_{0}$ is elementary if there exists a $k$-cell $c$ for which supp $f_{0}=\{c,-c\}$. Since $g$ only has support inside $* B$, the support is finite. Therefore, we can write the 3 -form $g$ as a finite sum of elementary 3 -forms $g_{0}$, whose support is contained in $* B$. Since

$$
\sum_{e \in B_{q}}(* g)_{e}=\sum_{e \in B_{q}} \sum_{g_{0}}\left(* g_{0}\right)_{e},
$$

we only have to prove that

$$
\sum_{e \in B_{q}}\left(* g_{0}\right)_{e}=0
$$

for each $g_{0}$.
Take any elementary 3 -form $g_{0}$ whose support is contained in $* B$. Then, the Hodge dual of $g_{0}$ is an elementary 1 -form $* g_{0}$. Since $g_{0}$ is elementary, there exists
a 3-cell $c_{0}$ in $* B$, such that supp $g_{0}=\left\{c_{0},-c_{0}\right\}$. Likewise, we have supp $\left(* g_{0}\right)=$ $\left\{e_{0},-e_{0}\right\}$, where $e_{0}=* c_{0}$. Since $e_{0}=* c_{0}$, we have $\hat{\partial} e_{0}=\hat{\partial}\left(* c_{0}\right)=*\left(\partial c_{0}\right)$. Since $\partial c_{0}$ is in $* B$, it follows that $*\left(\partial c_{0}\right)=\hat{\partial} e_{0}$ is in $* * B$. Therefore, $\left(\hat{\partial} e_{0} \cap\right.$ $\operatorname{supp} q) \subseteq(* * B \cap \operatorname{supp} q)$. Since we assumed that $(* * B) \cap \operatorname{supp} q$ only contains internal plaquettes of $q$, the plaquettes $p \in \hat{\partial} e_{0} \cap \operatorname{supp} q$ must be internal plaquettes of $q$. Recall that a plaquette $p \in \hat{\partial} e_{0} \cap \operatorname{supp} q$ is an internal plaquette if $(\partial p \cup \partial(-p)) \cap B_{q}=\varnothing$. Since $\pm e_{0} \in \partial p \cup \partial(-p)$, it follows that $\pm e_{0} \notin B_{q}$. Therefore, $B_{q}$ only contains elements for which $* g_{0}$ is zero. Thus,

$$
\sum_{e \in B_{q}}\left(* g_{0}\right)_{e}=0
$$

and

$$
0=\sum_{e \in B_{q}}(* g)_{e}=\sum_{p \in P_{N}} q_{p}^{+} \nu_{p} .
$$

## Chapter 5

## The existence and translation invariance of the limit

This chapter focuses on the limit of the expectation of the Wilson loop observable $W_{\gamma}$. This limit is also called the infinite volume limit. Both the existence and the translation invariance of this limit will be proved. For these proofs, Ginibre's inequality is needed. Before stating the inequality, we define both cones and convex cones.

Definition 5.1. (Cone and convex cone) $A$ set $A \subset V$, where $V$ is a vector space, is a cone if for every $x \in A$ and positive $\alpha \in \mathbb{R}$ we have $\alpha x \in A$. The cone is convex if for positive $\alpha, \beta \in \mathbb{R}$ and $x, y \in A$ we have $\alpha x+\beta y \in A$ [6].

We will use the notation $\operatorname{Cone}(A)$ for a convex cone generated by the set $A$. This means that the cone is the intersection of all convex cones containing $A$.

Lemma 5.2. (Ginibre's inequality). Let $K$ be a compact metric space and let $\mu$ be a probability measure on $K$. Let $C(K)$ be the algebra of complex-valued continuous functions on $K$. Let $S$ be a subset of $C(K)$ which is both invariant under complex conjugation and such that for any $f_{1}, \ldots f_{m} \in S$ and any choice of signs $s_{1}, \ldots, s_{m} \in\{-1,1\}$, the inequality

$$
\begin{equation*}
\iint \mu(x) \mu(y) \prod_{i=1}^{m}\left(f_{i}(x)+s_{i} f_{i}(y)\right) \geq 0 \tag{5.1}
\end{equation*}
$$

holds. Let $h$ be a real-valued function and let $f \in C(K)$. Define

$$
\langle f\rangle_{h}=\int f(x) e^{-h(x)} \mu(x) / \int e^{-h(x)} \mu(x)
$$

Let Cone $(S)$ be the convex cone generated by the set $S$. Then, for any $f, g,-h$ in Cone $(S),\langle f\rangle_{h},\langle g\rangle_{h}$ and $\langle f g\rangle_{h}$ are real and

$$
\begin{equation*}
\langle f g\rangle_{h} \geq\langle f\rangle_{h}\langle g\rangle_{h} \tag{5.2}
\end{equation*}
$$

For the proof, see [7].
In the following theorem, there is a function from the set of $G$-valued 1-forms on the set of all oriented edges in $\mathbb{Z}^{4}$. Therefore, we introduce the two following notations: $E_{\infty}$ is the set of oriented edges in $\mathbb{Z}^{4}$ and $\Sigma_{E_{\infty}}$ is the set of $G$-valued 1 -forms on $E_{\infty}$. In the theorem, $f$ only depends on the spins of the edges in $E_{M}$ and we have that $f(\sigma)=f\left(\left.\sigma\right|_{E_{M}}\right)$ for $\sigma \in \Sigma_{E_{\infty}}$. To simplify, when $N \geq M$ we write $f(\sigma)$ instead of $f\left(\left.\sigma\right|_{E_{M}}\right)$ and $f$ for the natural restriction of $f$ to $\Sigma_{E_{N}}$. A translation $\tau$ of the lattice changes where the origin of the lattice is. That the limit is translation invariant means that it does not depend on where in the lattice the Wilson loop observable is.

Theorem 5.3. Let $G=\mathbb{Z}_{n}$ and $f: \Sigma_{E_{\infty}} \rightarrow \mathbb{R}$ be a real-valued function, which only depends on the spins of edges in $E_{M}$ for some integer $M \geq 1$. Let $\beta \geq 0$. Then,
(i) the limit $\lim _{N \rightarrow \infty} \mathbb{E}_{\beta, N}[f(\sigma)]$ exists.
(ii) for any translation $\tau$ of $\mathbb{Z}^{4}$, we have

$$
\lim _{N \rightarrow \infty} \mathbb{E}_{\beta, N}[f \circ \tau(\sigma)]=\lim _{N \rightarrow \infty} \mathbb{E}_{\beta, N}[f(\sigma)]
$$

Proof. (i) To prove the existence of the limit, we write $f$ as a finite sum of functions $g_{j}$, for which the limit of the expectation value exists. We will show that $\left\langle g_{j}\right\rangle_{h}$ ( $h$ is defined later in the proof) is both increasing and bounded, which implies that the limit for $f$ exists.

First, we prove that $f$ can be written as a finite sum and that we can use Ginibre's inequality on $\left\langle g_{j}\right\rangle_{h}$. For this, we need some statements from Example 4 in [7]. These statements hold when $N \geq 1$ for the group $\Gamma^{(N)}=\left(\sum_{E_{N}},+\right)$. First, the group homomorphisms $\Gamma^{(N)} \mapsto \mathbb{C} \backslash\{0\}$ are given by the functions $\sigma \mapsto \prod_{e \in E_{N}} e^{2 \pi i \sigma_{e} \sigma_{e}^{\prime} / n}$, where $\sigma^{\prime}$ is fixed and $\sigma, \sigma^{\prime} \in \sum_{E_{N}}$. We notice that $\sigma \mapsto$ $\rho\left((d \sigma)_{p}\right), \sigma \in E_{N}$ is one of these group homomorphisms. Second, let $S^{(N)}$ be the set of real parts of the group homomorphisms $\Gamma^{(N)} \mapsto \mathbb{C} \backslash\{0\}$. Denote the
convex cone of this set by Cone $\left(S^{(N)}\right)$, then

$$
\Re\left(\rho\left((d \sigma)_{p}\right)\right) \in S^{(N)} \subseteq \operatorname{Cone}\left(S^{(N)}\right)
$$

Third, inequality 5.1 holds for the functions $f_{1}, \ldots, f_{n} \in S^{(N)}$ and the signs $s_{1}, \ldots, s_{m} \in\{-1,1\}$ when $\mu$ is the uniform measure on $\sum_{E_{N}}$.

We show that $f$ can be written as a finite sum of functions $g \in S^{(M)}$. Since $M \leq N$, the set $\Sigma_{E_{M}}$ is finite. Thus, the group homomorphisms $\Gamma^{(M)} \mapsto \mathbb{C} \backslash$ $\{0\}$ span the set of all real-valued functions on the set $\Sigma_{E_{M}}$. Therefore, there exists scalars $a_{1}, a_{2}, \ldots, a_{m} \in \mathbb{R}$ and functions $g_{1}, g_{2}, \ldots, g_{m} \in S^{(M)}$ such that the function $f$ can be written as $f=a_{1} g_{1}+\ldots+a_{m} g_{m}$. Since the functions $g_{1}, \ldots, g_{m} \in S^{(M)}$, they are also in Cone( $\left.S^{(M)}\right)$. For $M^{\prime} \geq M$ and $j \in\{1, \ldots, m\}$, we simplify and write $g_{j}$ for the natural extension of $g_{j}$ from $\Sigma_{E_{M}}$ to the larger set $\Sigma_{E_{M^{\prime}}}$.

We define the function $-h$ and show that $-h \in \operatorname{Cone}\left(S^{\left(N^{\prime}\right)}\right)$. Fix $N^{\prime} \geq N \geq$ $M$ and let $\beta, \beta^{\prime} \geq 0$. Then, for $\sigma \in \Sigma_{E_{N^{\prime}}}$ define

$$
-h(\sigma):=-h(\sigma)_{N, N^{\prime}, \beta, \beta^{\prime}}:=\beta \sum_{p \in P_{N}} \Re\left(\rho\left((d \sigma)_{p}\right)\right)+\beta^{\prime} \sum_{p \in P_{N^{\prime}} \backslash P_{N}} \Re\left(\rho\left((d \sigma)_{p}\right)\right) .
$$

By definition, $S^{\left(N^{\prime}\right)}$ contains the function $\sigma \mapsto \Re\left(\rho\left((d \sigma)_{p}\right)\right), \sigma \in \Sigma_{E_{N^{\prime}}}$. Since both $\beta$ and $\beta^{\prime}$ are non-negative, it follows from the definition of a convex cone that $-h(\sigma) \in \operatorname{Cone}\left(S^{\left(N^{\prime}\right)}\right)$.

We check the assumptions for Ginibre's inequality. Since the metric space $K=\sum_{E_{N}}$ is finite and equipped with discrete topology, it is compact. The set $S=S^{\left(N^{\prime}\right)}$ contains only continuous functions and since it consists of the real parts of the group homomorphisms, it is invariant under complex conjugation. Let the probability measure be the uniform measure $\mu$, thus, the third statement from the example fulfils inequality (5.1). Since Ginibre's inequality will be applied to $\left\langle g_{j} \sum_{p \in P_{N^{\prime}} \backslash P_{N}} \Re\left(\rho\left(\left(d \sigma_{p}\right)\right)\right)\right\rangle_{h}$, the functions $g_{j}(\sigma), \sum_{p \in P_{N^{\prime}} \backslash P_{N}} \Re\left(\rho\left(\left(d \sigma_{p}\right)\right)\right)$ and $-h$ must be in Cone $\left(S^{\left(N^{\prime}\right)}\right)$. We proved earlier that both $g_{j}$ and $-h$ are in Cone $\left(S^{\left(N^{\prime}\right)}\right)$. Since $\sum_{p \in P_{N^{\prime}} \backslash P_{N}} \Re\left(\rho\left(\left(d \sigma_{p}\right)\right)\right)$ is a sum of elements in $S^{\left(N^{\prime}\right)}$, it follows that it is in Cone $\left(S^{\left(N^{\prime}\right)}\right)$. Therefore, all assumptions for Ginibre's inequality are fulfilled.

We show that the derivative of $\left\langle g_{j}\right\rangle_{h}$ is non-negative in $\beta^{\prime}$ for $j \in\{1,2, \ldots, m\}$. Ginibre's inequality will be applied in the calculations:

$$
\frac{d}{d \beta^{\prime}}\left\langle g_{j}\right\rangle_{h}=\frac{d}{d \beta^{\prime}}\left(\frac{\int g_{j}(\sigma) e^{-h(\sigma)} \mu(\sigma)}{\int e^{-h(\sigma)} \mu(\sigma)}\right)
$$

$$
\begin{aligned}
&= \frac{\frac{d}{d \beta^{\prime}}\left(\int g_{j}(\sigma) e^{-h(\sigma)} \mu(\sigma)\right) \cdot \int e^{-h(\sigma)} \mu(\sigma)-\int g_{j}(\sigma) e^{-h(\sigma)} \mu(\sigma) \cdot \frac{d}{d \beta^{\prime}} \int e^{-h(\sigma)} \mu(\sigma)}{\left(\int e^{-h(\sigma)} \mu(\sigma)\right)^{2}} \\
&= \frac{\sum_{p \in P_{N^{\prime}} \backslash P_{N}} \Re\left(\rho\left((d \sigma)_{p}\right)\right) \int g_{j}(\sigma) e^{-h(\sigma)} \mu(\sigma)}{\int e^{-h(\sigma)} \mu(\sigma)} \\
&-\left\langle g_{j}\right\rangle_{h} \frac{\sum_{p \in P_{N^{\prime}} \backslash P_{N}} \Re\left(\rho\left((d \sigma)_{p}\right)\right) \int e^{-h(\sigma)} \mu(\sigma)}{\int e^{-h(\sigma)} \mu(\sigma)} \\
&=\left\langle g_{j}(\sigma) \sum_{p \in P_{N^{\prime}} P_{N}} \Re\left(\rho\left((d \sigma)_{p}\right)\right)\right\rangle_{h}-\left\langle g_{j}(\sigma)\right\rangle_{h}\left\langle\sum_{p \in P_{N^{\prime}} \backslash P_{N}} \Re(\rho(d \sigma(p))\rangle_{h}\right. \\
& \stackrel{5.52}{\geq}\left\langle g_{j}(\sigma)\right\rangle_{h}\left\langle\sum_{p \in P_{N^{\prime}} \backslash P_{N}} \Re\left(\rho\left((d \sigma)_{p}\right)\right)\right\rangle_{h}-\left\langle g_{j}(\sigma)\right\rangle_{h}\left\langle\sum_{p \in P_{N^{\prime}} \backslash P_{N}} \Re\left(\rho\left((d \sigma)_{p}\right)\right\rangle_{h}\right. \\
&=0 .
\end{aligned}
$$

Thus, we have proved that $\left\langle g_{j}\right\rangle_{h_{N, N^{\prime}, \beta, \beta^{\prime}}}$ is increasing in $\beta^{\prime}$. Moreover, $g_{j}$ depends on a finite set of edges. Hence, $\left\|g_{j}\right\|_{\infty}<\infty$ and $\left\langle g_{j}\right\rangle_{h_{N, N^{\prime}, \beta, 0}}$ is bounded. The expectation value of $g_{j}$ is

$$
\mathbb{E}_{\beta, N}\left[g_{j}(\sigma)\right]=\left\langle g_{j}\right\rangle_{h_{N, N^{\prime}, \beta, 0}},
$$

which was proved to be both bounded and monotone in $\beta$. Therefore, the limit of $\mathbb{E}_{\beta, N}\left[g_{j}(\sigma)\right]$ exists.

We have left to prove that $\lim _{N \rightarrow \infty} \mathbb{E}_{\beta, N}[f(\sigma)]$ exists:

$$
\begin{aligned}
\sum_{j=1}^{m} a_{j} \lim _{N \rightarrow \infty} \mathbb{E}_{\beta, N}\left[g_{j}(\sigma)\right] & =\lim _{N \rightarrow \infty} \sum_{j=1}^{m} a_{j} \mu_{\beta, N}\left[g_{j}(\sigma)\right]=\lim _{N \rightarrow \infty} \mathbb{E}_{\beta, N}\left[\sum_{j=1}^{m} a_{j} g_{j}(\sigma)\right] \\
& =\lim _{N \rightarrow \infty} \mathbb{E}_{\beta, N}[f(\sigma)] .
\end{aligned}
$$

Thus, the limit exists.
ii) We prove that the limit is translation invariant. Let $\tau$ be a translation of $\mathbb{Z}^{4}$ and choose $N^{\prime} \geq N \geq M$ large enough for $B_{N^{\prime}}$ to contain $B_{N}, \tau B_{N}$ and $\tau^{-1} B_{N}$. An illustration of this translation is given in Figure 5.1.


Figure 5.1: The sets $B_{N}, B_{N^{\prime}}$ and $\tau B_{N^{\prime}}$.

Define the functions $h_{N}:=h_{N, N^{\prime}, \beta, 0}$ and $h_{N^{\prime}}:=h_{N, N^{\prime}, \beta, \beta}$. For a translation $\tau$ of $h_{N}$, we have

$$
h_{N} \circ \tau=\beta \sum_{p \in P_{N}} \Re\left(\left(\rho\left((d \sigma)_{\tau(p)}\right)\right)=\beta \sum_{p \in \tau\left(P_{N}\right)} \Re\left(\rho\left((d \sigma)_{p}\right)\right) .\right.
$$

Since both $\tau B_{N}$ and $\tau^{-1} B_{N}$ are in $B_{N^{\prime}}$ and $\left\langle g_{j}\right\rangle_{h}$ is increasing in $N$ (This can be proved with a similar argument to that in part i).), it follows that

$$
\begin{equation*}
\left\langle g_{j}\right\rangle_{h_{N} \circ \tau} \leq\left\langle g_{j}\right\rangle_{h_{N^{\prime}}} \tag{5.3}
\end{equation*}
$$

and

$$
\begin{equation*}
\left\langle g_{j}\right\rangle_{h_{N} \circ \tau^{-1}} \leq\left\langle g_{j}\right\rangle_{h_{N^{\prime}}} \tag{5.4}
\end{equation*}
$$

We derive two inequalities for the expectation value of $g_{j}$ and $g_{j} \circ \tau$ :

$$
\begin{aligned}
\mathbb{E}_{\beta, N}\left[g_{j}(\sigma)\right] & =\left\langle g_{j}\right\rangle_{h_{N}} \\
& =\left\langle g_{j} \circ \tau \circ \tau^{-1}\right\rangle_{h_{N}} \\
& =\left\langle g_{j} \circ \tau\right\rangle_{h_{N} \circ \tau} \\
& \stackrel{5.3]}{\leq}\left\langle g_{j} \circ \tau\right\rangle_{h_{N^{\prime}}}=\mathbb{E}_{\beta, N^{\prime}}\left[g_{j} \circ \tau(\sigma)\right]
\end{aligned}
$$

and

$$
\mathbb{E}_{\beta, N}\left[g_{j} \circ \tau(\sigma)\right]=\left\langle g_{j} \circ \tau\right\rangle_{h_{N}}=\left\langle g_{j}\right\rangle_{h_{N} \circ \tau^{-1}} \stackrel{\sqrt{5.4}}{\leq}\left\langle g_{j}\right\rangle_{h_{N^{\prime}}}=\mathbb{E}_{\beta, N^{\prime}}\left[g_{j}(\sigma)\right] .
$$

Take the limits, first for $N^{\prime}$ and then for $N$ on both inequalities. We obtain

$$
\lim _{N \rightarrow \infty} \mathbb{E}_{\beta, N}\left[g_{j}(\sigma)\right] \leq \lim _{N^{\prime} \rightarrow \infty} \mathbb{E}_{\beta, N^{\prime}}\left[g_{j} \circ \tau(\sigma)\right],
$$

and

$$
\lim _{N \rightarrow \infty} \mathbb{E}_{\beta, N}\left[g_{j} \circ \tau(\sigma)\right] \leq \lim _{N^{\prime} \rightarrow \infty} \mathbb{E}_{\beta, N^{\prime}}\left[g_{j}(\sigma)\right]
$$

Combining these inequalities, we obtain

$$
\lim _{N \rightarrow \infty} \mathbb{E}_{\beta, N}\left[g_{j} \circ \tau(\sigma)\right] \leq \lim _{N^{\prime} \rightarrow \infty} \mathbb{E}_{\beta, N^{\prime}}\left[g_{j}(\sigma)\right] \leq \lim _{N^{\prime} \rightarrow \infty} \mathbb{E}_{\beta, N^{\prime}}\left[g_{j} \circ \tau(\sigma)\right]
$$

and as a result,

$$
\lim _{N \rightarrow \infty} \mathbb{E}_{\beta, N}\left[g_{j} \circ \tau(\sigma)\right]=\lim _{N^{\prime} \rightarrow \infty} \mathbb{E}_{\beta, N^{\prime}}\left[g_{j}(\sigma)\right] .
$$

Thus, the limit of the expectation of $g_{j}$ is translation invariant, but we wish to prove this property for the expectation of $f$. Similarly to the end of the proof in part (i), we obtain

$$
\begin{aligned}
\lim _{N \rightarrow \infty} \mathbb{E}_{\beta, N}[f \circ \tau(\sigma)] & =\lim _{N \rightarrow \infty} \mathbb{E}_{\beta, N}\left[\sum_{j=1}^{m} a_{j} g_{j} \circ \tau(\sigma)\right]=\sum_{j=1}^{m} a_{j} \lim _{N \rightarrow \infty} \mathbb{E}_{\beta, N}\left[g_{j} \circ \tau(\sigma)\right] \\
& =\sum_{j=1}^{m} a_{j} \lim _{N \rightarrow \infty} \mathbb{E}_{\beta, N}\left[g_{j}\right]=\lim _{N \rightarrow \infty} \mathbb{E}_{\beta, N}\left[\sum_{j=1}^{m} a_{j} g_{j}\right] \\
& =\lim _{N \rightarrow \infty} \mathbb{E}_{\beta, N}[f(\sigma)] .
\end{aligned}
$$

Therefore, the limit of the expectation value of $f$ is translation invariant.
We check the assumptions for the Wilson loop observable

$$
W_{\gamma}=\rho\left(\sum_{e \in \gamma} \sigma_{e}\right), \sigma \in \Sigma_{E_{N}} .
$$

Since the Wilson loop observable maps $\Sigma_{E_{\infty}}$ to $\mathbb{R}$, has real values and only depends on the spins of edges in $E_{M}$, the function $f$ in Theorem 5.3 can be chosen as $W_{\gamma}$. Thus, the limit

$$
\lim _{N \rightarrow \infty} \mathbb{E}_{\beta, N}\left[W_{\gamma}\right]=\left\langle W_{\gamma}\right\rangle_{\beta}
$$

both exists and is translation invariant.

## Chapter 6

## Proof of the main theorem

The proof is divided into two cases; one when $\ell \lambda(\beta)^{12}$ is small and one when $\ell \lambda(\beta)^{12}$ is large. In the final part of the proof, these two parts are combined. We begin with the case when $\ell \lambda(\beta)^{12}$ is small.

### 6.1 The case when $\ell \lambda(\beta)^{12}$ is small

First, the proposition for the case when $\ell \lambda(\beta)^{12}$ is small is stated, then some lemmas necessary for the proof of the proposition are given and proved. Last, the proposition is proved.

Proposition 6.1. Consider lattice gauge theory with the structure group $G=\mathbb{Z}_{n}$ and a one-dimensional faithful representation $\rho$ of $G$. Let $\gamma$ be a simple oriented loop in $\mathbb{Z}^{4}, \ell=|\gamma|$ the length of it and $\ell_{c}$ the number of corner edges in $\gamma$. Let $N$ be large enough so that $\gamma \subseteq E_{N}$ and such that there exists a cube $B$ of width $|\gamma|$ containing $\gamma$ inside $B_{N}$. Let $\beta_{0}>0$ be such that $5(|G|-1) \lambda(\beta)^{2}<1$ for all $\beta>\beta_{0}$. Then for all $\beta>\beta_{0}$, we have

$$
\left|\mathbb{E}_{\beta, N}\left[W_{\gamma}\right]-e^{-\ell(1-\theta(\beta))}\right| \leq C_{A} e^{2 C^{*} \ell \lambda(\beta)^{12}}\left(\sqrt{\frac{\ell_{c}}{\ell}}+\lambda(\beta)^{2}\right)
$$

Here $\theta(\beta)$ is defined by equation (2.7) and $C_{A}$ by equation (6.41). Furthermore, $C_{A}$ only depend on $\beta_{0}, G$ and $\rho$.

To prove this proposition, we show that it is very likely that only minimal vortices, which are centred at the edges in $\gamma$, have an impact on a Wilson loop (A minimal vortex is centred at $d x_{j}$ if it can be written as $d\left(g d x_{j}\right)$ for a $g \in G \backslash\{0\}$
and a $d x_{j} \in E_{N}$.). A number of lemmas will be needed. We begin with defining the constant $C^{*}$ and recalling the definitions for $\theta(\beta)$ and $\lambda(\beta)$. Let $\beta_{0}>0$ be given and define

$$
\begin{equation*}
C^{*}:=\sup _{\beta>\beta_{0}}\left[(1-\theta(\beta)) \lambda(\beta)^{-12}\right] . \tag{6.1}
\end{equation*}
$$

Recall that

$$
\theta(\beta)=\frac{\sum_{g \in G} \rho(g) \phi_{\beta}(g)^{12}}{\sum_{g \in G} \phi_{\beta}(g)^{12}}
$$

and

$$
\lambda(\beta)=\max _{g \in G \backslash\{0\}} \frac{\phi_{\beta}(g)}{\phi_{\beta}(0)},
$$

where

$$
\phi_{\beta}(g):=e^{\beta \Re \rho(g)}, \quad g \in G .
$$

### 6.1.1 Technical lemmas

Lemma 6.2. The function $\theta(\beta)$ is a real-valued function and can be written as

$$
\begin{equation*}
\theta(\beta)=\frac{\sum_{g \in G} \Re(\rho(g)) \phi_{\beta}(g)^{12}}{\sum_{g \in G} \phi_{\beta}(g)^{12}} \tag{6.2}
\end{equation*}
$$

Proof. Since $\rho(g)^{-1}=\rho(g)^{*}$, we have

$$
\Re(\rho(-g))=\Re\left(\rho(g)^{-1}\right)=\Re\left(\rho(g)^{*}\right)=\Re(\rho(g)) .
$$

Therefore,

$$
\phi_{\beta}(-g)=e^{\beta \Re(\rho(-g))}=e^{\beta \Re(\rho(g))}=\phi_{\beta}(g) \quad \text { for } g \in G .
$$

Furthermore, since $g \rightarrow-g$ is a bijection, we have that

$$
\begin{aligned}
\theta(\beta) & =\frac{1}{2} \frac{\sum_{g \in G} \rho(g) \phi_{\beta}(g)^{12}}{\sum_{g \in G} \phi_{\beta}(g)^{12}}+\frac{1}{2} \frac{\sum_{g \in G} \rho(-g) \phi_{\beta}(-g)^{12}}{\sum_{g \in G} \phi_{\beta}(-g)^{12}} \\
& =\frac{1}{2} \frac{\sum_{g \in G} \rho(g) \phi_{\beta}(g)^{12}}{\sum_{g \in G} \phi_{\beta}(g)^{12}}+\frac{1}{2} \frac{\sum_{g \in G} \rho(-g) \phi_{\beta}(g)^{12}}{\sum_{g \in G} \phi_{\beta}(g)^{12}} \\
& =\frac{\sum_{g \in G}(\rho(g)+\rho(-g)) \phi_{\beta}(g)^{12}}{\sum_{g \in G} \phi_{\beta}(g)^{12}} \\
& =\frac{\sum_{g \in G} \Re(\rho(g)) \phi_{\beta}(g)^{12}}{\sum_{g \in G} \phi_{\beta}(g)^{12}} \in \mathbb{R} .
\end{aligned}
$$

Lemma 6.3. Let $\beta \geq 0$ be such that $5(|G|-1) \lambda(\beta)^{2}<1$. Then

$$
\begin{equation*}
0<1-\frac{2}{5^{6}(|G|-1)^{5}}<1-C^{*} \lambda(\beta)^{12} \leq \theta(\beta) \leq 1 \tag{6.3}
\end{equation*}
$$

Proof. We calculate the upper bound for $1-C^{*} \lambda(\beta)^{12}$ :

$$
\begin{align*}
1-C^{*} \lambda(\beta)^{12} & =1-\sup _{\beta>\beta_{0}}\left[(1-\theta(\beta)) \lambda(\beta)^{-12}\right] \lambda(\beta)^{12} \\
& \leq 1-\left((1-\theta(\beta)) \lambda(\beta)^{-12}\right) \lambda(\beta)^{12}  \tag{6.4}\\
& =\theta(\beta)
\end{align*}
$$

Since Lemma 6.2 gives that $\theta(\beta)$ is real, we have

$$
\begin{equation*}
\theta(\beta) \leq|\theta(\beta)|=\left|\frac{\sum_{g \in G} \rho(g) \phi_{\beta}(g)^{12}}{\sum_{g \in G} \phi_{\beta}(g)^{12}}\right| \leq \frac{\sum_{g \in G}|\rho(g)| \phi_{\beta}(g)^{12}}{\sum_{g \in G} \phi_{\beta}(g)^{12}} \stackrel{(2.3)}{=} \frac{\sum_{g \in G} \phi_{\beta}(g)^{12}}{\sum_{g \in G} \phi_{\beta}(g)^{12}}=1 . \tag{6.5}
\end{equation*}
$$

We calculate an upper bound for $C^{*}$. Notice that $\Re \rho(g) \in[-1,1[$ for $g \in G \backslash\{0\}$ and $\phi_{\beta}>0$. Hence,

$$
\begin{aligned}
& \stackrel{(2.2)}{=} \sup _{\beta>\beta_{0}}\left[\frac{\sum_{g \in G \backslash\{0\}}(1-\Re \rho(g)) \phi_{\beta}(g)^{12}}{\sum_{g \in G} \phi_{\beta}(g)^{12}} \lambda(\beta)^{-12}\right] \\
& \leq \sup _{\beta>\beta_{0}}\left[\frac{\sum_{g \in G \backslash\{0\}} 2 \phi_{\beta}(g)^{12}}{\sum_{g \in G} \phi_{\beta}(g)^{12}} \lambda(\beta)^{-12}\right] \\
& =\sup _{\beta>\beta_{0}}\left[\frac{\sum_{g \in G \backslash\{0\}} 2 \phi_{\beta}(g)^{12}}{\phi_{\beta}(0)^{12}+\sum_{g \in G \backslash\{0\}} \phi_{\beta}(g)^{12}} \lambda(\beta)^{-12}\right] \\
& \leq \sup _{\beta>\beta_{0}}\left[2 \sum_{g \in G \backslash\{0\}} \frac{\phi_{\beta}(g)^{12}}{\phi_{\beta}(0)^{12}} \lambda(\beta)^{-12}\right] \\
& =\sup _{\beta>\beta_{0}}\left[2 \sum_{g \in G \backslash\{0\}} \lambda(\beta)^{12} \lambda(\beta)^{-12}\right] \\
& =2 \sum_{g \in G \backslash\{0\}} 1=2(|G|-1) \text {. }
\end{aligned}
$$

Therefore,

$$
\begin{align*}
C^{*} \lambda(\beta)^{12} & \leq 2(|G|-1) \lambda(\beta)^{12}=\frac{2}{5^{6}(|G|-1)^{5}}\left(5(|G|-1) \lambda(\beta)^{2}\right)^{6} \\
& <\frac{2}{5^{6}(|G|-1)^{5}} . \tag{6.6}
\end{align*}
$$

Thus,

$$
0<1-\frac{2}{5^{6}(|G|-1)^{5}} \stackrel{\sqrt{6.6 \mid}}{<} 1-C^{*} \lambda(\beta)^{12} \stackrel{\sqrt{6.4]}}{\leq} \theta(\beta) \stackrel{\sqrt{6.5]}}{\leq} 1 .
$$

Notice that $C^{*}>0$ by equations (6.4) and (6.5). The next lemma gives two inequalities for the function $\theta$.

Lemma 6.4. Let $\beta>\beta_{0}$. For $0 \leq j \leq \ell$, we have

$$
\begin{equation*}
\theta^{-j} \leq 2 e^{C^{*} \ell \lambda(\beta)^{12}} \tag{6.7}
\end{equation*}
$$

For $1 \leq j \leq \ell$, we have

$$
\begin{equation*}
\theta^{-j}-1 \leq 2 j C^{*} \lambda(\beta)^{12} e^{C^{*} \ell \lambda(\beta)^{12}} \tag{6.8}
\end{equation*}
$$

For the proof of this lemma, the two following inequalities are needed.
First, for $x \in\left[0, \frac{1}{2}\right]$, we have

$$
\begin{equation*}
(1-x)^{-1} \leq 2 e^{x}, \tag{6.9}
\end{equation*}
$$

and second, Bernoulli's inequality: For $x \geq-1$ and $n \geq 1$, we have

$$
\begin{equation*}
(1+x)^{n} \geq 1+n x . \tag{6.10}
\end{equation*}
$$

Proof of Lemma 6.4. First, notice that

$$
0<C^{*} \lambda(\beta)^{12} \stackrel{\sqrt{6.6}}{<} \frac{2}{5^{6}(|G|-1)^{5}} \leq \frac{2}{5^{6}}<\frac{1}{2}
$$

Therefore, inequality (6.9) can be applied on $\left(1-C^{*} \lambda(\beta)^{12}\right)^{-1}$. For $0 \leq j \leq \ell$, we have

$$
\begin{equation*}
\theta^{-j} \leq \theta^{-\ell} \stackrel{\sqrt{6.3]}}{\leq}\left(\left(1-C^{*} \lambda(\beta)^{12}\right)^{-1}\right)^{\ell} \stackrel{\sqrt{6.9})}{\leq} 2 e^{C^{*} \ell \lambda(\beta)^{12}} \tag{6.11}
\end{equation*}
$$

For $1 \leq j \leq \ell$, we have by inequality (6.10) that

$$
\theta^{j}=(1+(\theta-1))^{j} \geq 1+j(\theta-1) .
$$

Thus,

$$
1-\theta^{j} \leq j(1-\theta)
$$

and

$$
\begin{equation*}
\theta^{-j}-1=\theta^{-j}\left(1-\theta^{j}\right) \leq j(1-\theta) \theta^{-j} \stackrel{\sqrt{6.3},, \sqrt{6.11]}}{\leq} 2 j C^{*} \lambda(\beta)^{12} e^{C^{*} \ell \lambda(\beta)^{12}} . \tag{6.12}
\end{equation*}
$$

Lemma 6.5. For $\ell \geq 1$ and $\theta=\theta(\beta)$, we have

$$
\begin{equation*}
\left|\theta^{\ell}-e^{-\ell(1-\theta)}\right| \leq \frac{e^{2 C^{*} \ell \lambda(\beta)^{12}}}{2 \ell} \tag{6.13}
\end{equation*}
$$

We state two inequalities necessary for the proof of the lemma. First, for $\ell \geq 1,|x| \leq 1$ and $|y| \leq 1$, we have

$$
\begin{equation*}
\left|x^{\ell}-y^{\ell}\right| \leq \ell|x-y| \tag{6.14}
\end{equation*}
$$

Indeed, by the intermediate value theorem, there exists a $\xi \in[-1,1]$ such that

$$
\left|x^{\ell}-y^{\ell}\right|=\left|\ell \xi^{\ell-1}\right||x-y| \leq \ell|x-y| .
$$

Second, for $x>0$, we have

$$
\begin{equation*}
1+x \leq e^{x} \tag{6.15}
\end{equation*}
$$

This follows from the Taylor series for the exponential function:

$$
e^{x}=\sum_{k=0}^{\infty} \frac{x^{n}}{n!},
$$

which is also otherwise useful in the proof.
Proof of Lemma 6.5. Since $0 \leq 1-\theta<1$, it follows from the Taylor series that

$$
\begin{equation*}
e^{-(1-\theta)} \leq 1+(-(1-\theta))+\frac{(-(1-\theta))^{2}}{2}=\theta+\frac{(1-\theta)^{2}}{2} \tag{6.16}
\end{equation*}
$$

From equation 6.15) follows that

$$
\begin{equation*}
1+(\theta-1) \leq e^{(\theta-1)} \Leftrightarrow \theta \leq e^{-(1-\theta)} \tag{6.17}
\end{equation*}
$$

Thus,

$$
\begin{align*}
& \mid \theta^{\ell}-e^{-\ell(1-\theta)} \stackrel{\sqrt{6.144}}{\leq}\left|\theta-e^{-(1-\theta)}\right| \stackrel{\sqrt[6.17]{6}}{=} \ell\left(\theta-e^{-(1-\theta)}\right) \\
& \stackrel{\sqrt{6.16 \operatorname{lin}} \ell\left(2 \theta-2 \theta-(1-\theta)^{2}\right)}{2}=\frac{\ell(1-\theta)^{2}}{2} \\
& \stackrel{\sqrt{6.3)}\left(\ell C^{*} \lambda(\beta)^{12}\right)^{2}}{2 \ell}  \tag{6.18}\\
& \stackrel{(6.155}{\leq} e^{2 C^{*} \ell \lambda(\beta)^{12}} \\
& 2 \ell
\end{align*}
$$

### 6.1.2 Probabilistic bounds on vortex sizes

We want to calculate the probabilities for the events $A_{1}, A_{2}$ and $A_{3}$, which are soon defined. Before that, some necessary definitions and notations are given. Let $\gamma$ be a simple loop, then the cube $B$ is defined as a fixed cube that contains $\gamma$, is inside $B_{N}$ and is of width $\ell=|\gamma|$. Let $q$ be an fixed oriented surface such that its support is contained in $B$ and such that $\gamma$ is the boundary of $q$ (we will later prove that $q$ exists). Denote the support of this oriented surface by $Q:=\operatorname{supp} q$. Then, define the set $Q^{\prime}$ as the set of plaquettes $p \in Q$ for which the oriented loop $\gamma$ does not intersect any cube of width $b+2$ containing $p$, where $b$ is chosen as the smallest number for which the support of any irreducible $\omega \in \sum_{P_{N}}$ with $|\operatorname{supp} \nu| \leq 48$ is inside the cube. Last, the set

$$
\gamma_{c}:=\{e \in \gamma:|\partial(\hat{\partial} e) \cap(\gamma \cup-\gamma)| \geq 2\}
$$

is the set of corner edges in $\gamma$. Recall that the number of corner edges in $\gamma$ is given by $\ell_{c}$.

We define the events:
$A_{1}:=\{$ There exists no vortex $\nu$ in $\sigma$ with $|\operatorname{supp} \nu| \geq 50$ whose support intersects $Q$ \}
$A_{2}:=\{$ There exists no vortex $\nu$ in $\sigma$ with $|\operatorname{supp} \nu| \geq 14$ and

$$
\left.\operatorname{supp} \nu \cap\left(Q \backslash Q^{\prime}\right) \neq \varnothing\right\}
$$

$A_{3}:=\left\{\right.$ There exists no vortex $\nu$ in $\sigma: \operatorname{supp} \nu=\hat{\partial} e \cup \hat{\partial}(-e)$ for some $\left.e \in \gamma_{c}\right\}$.
Notice that the vortex in the event $A_{3}$ is minimal by Lemma 4.7. The probabilities of these events are calculated in the following three lemmas.

Lemma 6.6. Let $\beta_{0}>0$ be such that $5(|G|-1) \lambda(\beta)^{2}<1$ for all $\beta>\beta_{0}$. Then

$$
\begin{equation*}
\mu_{\beta, N}\left(A_{1}\right) \geq 1-C_{0}^{(25)} \ell^{4} \lambda(\beta)^{50} . \tag{6.19}
\end{equation*}
$$

Proof. Fix $p_{0} \in P_{N}$ and define the event

$$
\begin{aligned}
E_{1}\left(p_{0}\right):= & \left\{\text { There exists at least one vortex } \nu \text { in } \sigma \text { with } p_{0} \in \operatorname{supp} \nu\right. \\
& \text { and }|\operatorname{supp} \nu| \geq 50\} .
\end{aligned}
$$

The assumptions for Proposition 4.12 are fulfilled for $E_{1}$. Therefore,

$$
\begin{equation*}
\mu_{\beta, N}\left(E_{1}\left(p_{0}\right)\right) \leq C_{0}^{(25)} \lambda(\beta)^{50}, \quad \text { for any } p_{0} \in P_{N} \tag{6.20}
\end{equation*}
$$

Since $B$ is of width $\ell$ and $\operatorname{supp} q$ is contained in $B$, there is at most $\ell^{4}$ choices for $p_{0} \in \operatorname{supp} \nu$ for which supp $\nu$ intersects $Q$. Hence, by a union bound and inequality (6.20), we have

$$
\mu_{\beta, N}\left(A_{1}^{c}\right) \leq \ell^{4} C_{0}^{(25)} \lambda(\beta)^{50}
$$

Thus,

$$
\mu_{\beta, N}\left(A_{1}\right)=1-\mu_{\beta, N}\left(A_{1}^{c}\right) \geq 1-C_{0}^{(25)} \ell^{4} \lambda(\beta)^{50}
$$

Lemma 6.7. Let $\beta_{0}>0$ be such that $5(|G|-1) \lambda(\beta)^{2}<1$ for all $\beta>\beta_{0}$. Then

$$
\mu_{\beta, N}\left(A_{2}\right) \geq 1-C_{1} C_{0}^{(7)} \ell \lambda(\beta)^{14}
$$

Proof. Fix $p_{0} \in P_{N}$. The event $E_{2}$ is defined by

$$
E_{2}\left(p_{0}\right):=\left\{\sigma \in \Sigma_{E_{N}}: \exists \text { a vortex } \nu \text { in } \sigma \text { with } p_{0} \in \operatorname{supp} \nu \text { and }|\operatorname{supp} \nu| \geq 14\right\}
$$

By Proposition 4.12, we have

$$
\mu_{\beta, N}\left(E_{2}\left(p_{0}\right)\right) \leq C_{0}^{(7)} \lambda(\beta)^{14}
$$

Consider the plaquettes $p \in Q \backslash Q^{\prime}$. By definition, any cube of width $b+2$ that contains one of these plaquettes intersects $\gamma$. Hence

$$
\begin{equation*}
\left|Q \backslash Q^{\prime}\right| \leq C_{1}|\gamma|=C_{1} \ell \tag{6.21}
\end{equation*}
$$

where the constant $C_{1}$ depends on the width of the box.
By a union bound,

$$
\mu_{\beta, N}\left(A_{2}^{c}\right) \leq C_{1} \ell \mu_{\beta, N}\left(E_{2}\left(p_{0}\right)\right) .
$$

Thus,

$$
\mu_{\beta, N}\left(A_{2}\right)=1-\mu_{\beta, N}\left(A_{2}^{c}\right) \geq 1-C_{1} \ell \mu_{\beta, N}\left(E_{2}\left(p_{0}\right)\right) \geq 1-C_{1} C_{0}^{(7)} \ell \lambda(\beta)^{14}
$$

Lemma 6.8. Let $\beta_{0}>0$ be such that $5(|G|-1) \lambda(\beta)^{2}<1$ for all $\beta>\beta_{0}$. Then

$$
\begin{equation*}
\mu_{\beta, N}\left(A_{3}\right) \geq 1-C^{(6)} \ell_{c} \lambda(\beta)^{12} \tag{6.22}
\end{equation*}
$$

Proof. By Proposition 4.12 for a minimal vortex and the fact that there are at most $\ell_{c}$ corner edges for which supp $\nu=\hat{\partial} e \cup \hat{\partial}(-e)$, we have that

$$
\mu_{\beta, N}\left(A_{3}^{c}\right) \leq \ell_{c} C^{(6)} \lambda(\beta)^{12}
$$

and

$$
\mu_{\beta, N}\left(A_{3}\right) \geq 1-C^{(6)} \ell_{c} \lambda(\beta)^{12} .
$$

### 6.1.3 The main argument

We continue with some more definitions and notations for sets and conditional probability. Define the set of edges

$$
\gamma_{1}:=\gamma \backslash \gamma_{c}
$$

and the random set of edges

$$
\gamma^{\prime}:=\left\{e \in \gamma_{1}: \exists p, p^{\prime} \in \hat{\partial} e \text { with }(d \sigma)_{p} \neq(d \sigma)_{p^{\prime}}\right\} .
$$

Notice that $\gamma_{1}$ is the set of the non-corner edges of $\gamma$ (Recall that both corner and non-corner edges are illustrated in Figure 2.2.).

We investigate how $d \sigma$ can be written on two different plaquettes $p, p^{\prime} \in \hat{\partial} e$ :

$$
(d \sigma)_{p}=\sum_{e^{\prime} \in \partial p} \sigma_{e^{\prime}}=\sigma_{e}+\sum_{e^{\prime} \in \partial p \backslash\{e\}} \sigma_{e^{\prime}}
$$

and

$$
(d \sigma)_{p^{\prime}}=\sum_{e^{\prime} \in \partial p^{\prime}} \sigma_{e^{\prime}}=\sigma_{e}+\sum_{e^{\prime} \in \partial p^{\prime} \backslash\{e\}} \sigma_{e^{\prime}} .
$$

Therefore, the event $(d \sigma)_{p} \neq(d \sigma)_{p^{\prime}}$ does not depend on the value of $\sigma_{e}$. Hence, if we know the spins of the edges which are not in $\pm \gamma_{1}$, then the edges in $\gamma^{\prime}$ are determined. Furthermore, a plaquette cannot contain two non-corner edges. Therefore, when we condition on $\left(\sigma_{e}\right)_{e \notin \pm \gamma_{1}}$, the spins $\left(\sigma_{e}\right)_{e \in \gamma_{1}}$ are independent. This is concluded in the following lemma.

Lemma 6.9. The random set $\gamma^{\prime}$ is determined by the spins $\left(\sigma_{e}\right)_{e \notin \pm \gamma_{1}}$ and conditioning on the latter, $\left(\sigma_{e}\right)_{e \in \gamma_{1}}$ are independent.

Lastly, we have one more inequality for an expectation value, the expected value of the number of edges in the set $\gamma^{\prime}$.

Lemma 6.10. Let $\beta_{0}>0$ be such that $5(|G|-1) \lambda(\beta)^{2}<1$ for all $\beta>\beta_{0}$. Then

$$
\begin{equation*}
\mathbb{E}_{\beta, N}\left[\left|\gamma^{\prime}\right|\right] \leq C_{0}^{(6)} \ell \lambda(\beta)^{12} . \tag{6.23}
\end{equation*}
$$

Proof. If $(d \sigma)_{p} \neq 0$, then the plaquette is in the support of a vortex $\nu$. Since the size of the vortex is 12 or larger, we have that $|\operatorname{supp} \nu| \geq 12$ if $\operatorname{supp} \nu$ contains $(d \sigma)_{p} \neq 0$. Then, from Proposition 4.12 follows for a fixed $p_{0}$ that

$$
\mu_{\beta, N}\left\{\text { A plaquette } p \text { is such that }(d \sigma)_{p} \neq 0\right\} \leq C_{0}^{(6)} \lambda(\beta)^{12} .
$$

Since $|\gamma|=\ell$ and $\gamma^{\prime} \subseteq \gamma$, there are at most $\ell$ plaquettes in $\gamma^{\prime}$ for which $(d \sigma)_{p} \neq 0$. Thus,

$$
\mathbb{E}_{\beta, N}\left[\left|\gamma^{\prime}\right|\right] \leq C_{0}^{(6)} \ell \lambda(\beta)^{12}
$$

We prove the proposition. Six inequalities and equations for expectations will be calculated and then combined to obtain the desired inequality.

Proof of Proposition 6.1. Since $\gamma$ is a simple loop and $B \subseteq B_{N}$, we have by Lemma 4.16 there exists an oriented surface $q$ such that its support is contained in $B$ and $\gamma$ is the boundary of $q$. Recall that the support of this oriented surface is $Q$. Let $\sigma \sim \mu_{\beta, N}$.

Since $\sigma \in E_{N}$, it follows from Lemma 4.5 that we can write $d \sigma$ as a sum of vortices $\nu_{1}, \nu_{2}, \ldots$ with disjoint supports. Let the decomposition be fixed and define the set $V$ as the set of vortices in this decomposition for which the support intersects $Q$. For the oriented surface $q$ with boundary $\gamma$, we have by Lemma 4.15 that

$$
\sum_{e \in \gamma} \sigma_{e}=\sum_{p \in P_{N}} q_{p}^{+}(d \sigma)_{p} \quad \text { for } \sigma \in \Sigma_{E_{N}} .
$$

Therefore, the Wilson loop observable $W_{\gamma}$ can be written as

$$
W_{\gamma}=\rho\left(\sum_{e \in \gamma} \sigma_{e}\right)=\rho\left(\sum_{p \in P_{N}} q_{p}^{+}(d \sigma)_{p}\right)=\rho\left(\sum_{v \in V} \sum_{p \in \operatorname{supp} \nu} q_{p}^{+}(d \sigma)_{p}\right)
$$

Note that since $|\rho(g)|=1$ for $g \in G$, we have $\left|W_{\gamma}\right|=1$.
The first expectation $\left(\mathbb{E}_{\beta, N}\left[\left|W_{\gamma}-W_{\gamma}^{0}\right|\right]\right)$
We define the subset $V_{0}$ of $V$ as

$$
V_{0}:=\{\nu \in V:|\operatorname{supp} \nu| \leq 48\}
$$

and $W_{\gamma}^{0}$ as

$$
W_{\gamma}^{0}:=\rho\left(\sum_{\nu \in V_{0}} \sum_{p \in \operatorname{supp} \nu} q_{p}^{+}(d \sigma)_{p}\right) .
$$

We estimate the expectation value of $\left|W_{\gamma}-W_{\gamma}^{0}\right|$. First, consider the case when the event $A_{1}$ occurs. Then there is no vortex with $|\operatorname{supp} \nu| \geq 50$, thus, $V_{0}=V$, $W_{\gamma}=W_{\gamma}^{0}$ and $\left|W_{\gamma}-W_{\gamma}^{0}\right|=0$. Second, consider the case when the event $A_{1}$ does not occur. Then,

$$
\left|W_{\gamma}-W_{\gamma}^{0}\right| \leq\left|W_{\gamma}\right|+\left|W_{\gamma}^{0}\right|=1+1=2 .
$$

Thus, by Lemma 6.6, the expectation value is

$$
\begin{equation*}
\mid \mathbb{E}_{\beta, N}\left[\left|W_{\gamma}-W_{\gamma}^{0}\right|\right] \leq 0 \cdot \mu_{\beta, N}\left(A_{1}\right)+2 \cdot \mu_{\beta, N}\left(A_{1}^{c}\right) \leq 2 C_{0}^{(25)} \ell^{4} \lambda(\beta)^{50} \tag{6.24}
\end{equation*}
$$

The second expectation $\left(\mathbb{E}_{\beta, N}\left[\left|W_{\gamma}^{0}-W_{\gamma}^{3}\right|\right]\right)$
Consider a cube $B$ of width $b$, where $b$ is chosen as the smallest number for which the support of any irreducible $\omega \in \sum_{P_{N}}$ with $|\operatorname{supp} \nu| \leq 48$ is inside the cube. Let

$$
V_{1}:=\left\{\nu \in V_{0}: \operatorname{supp} \nu \cap Q^{\prime} \neq \varnothing\right\} .
$$

Fix $\nu \in V_{1}$. Since $B$ is a cube of width $b$, we have from Example 3.30 that the cube $* * B$ is of width $b+2$. By definition, for plaquettes $p \in Q^{\prime}$ the intersection between a cube of width $b+2$ containing $p$ and $\gamma$ is empty. Hence, the intersection $(* * B) \cap Q$ only consists of internal plaquettes of $q$. Recall that $\gamma$ is the boundary of $q$. Thus, from Lemma 4.19 follows that

$$
\sum_{p \in \operatorname{supp} \nu} q_{p}^{+}(d \sigma)_{p}=0 \quad \text { for } v \in V_{1} .
$$

We define the set

$$
V_{2}:=V_{0} \backslash V_{1}
$$

and

$$
W_{\gamma}^{0}=\rho\left(\sum_{\nu \in V_{0}} \sum_{p \in \operatorname{supp} \nu} q_{p}^{+}(d \sigma)_{p}\right)=\rho\left(\sum_{\nu \in V_{2}} \sum_{p \in \operatorname{supp} \nu} q_{p}^{+}(d \sigma)_{p}\right) .
$$

We define a subset $V_{3}$ of $V_{2}$, which only contains minimal vortices. Let $V_{3}:=\left\{\nu \in V_{2}:|\operatorname{supp} \nu|=12\right\}=\left\{\nu \in V:|\operatorname{supp} \nu|=12\right.$ and $\left.\operatorname{supp} \nu \cap Q^{\prime}=\varnothing\right\}$ and

$$
W_{\gamma}^{3}:=\rho\left(\sum_{\nu \in V_{3}} \sum_{p \in \operatorname{supp} \nu} q_{p}^{+}(d \sigma)_{p}\right) .
$$

We calculate the expectation value of $\left|W_{\gamma}^{0}-W_{\gamma}^{3}\right|$. Consider the event $A_{2}$. If $A_{2}$ occurs, then there is no vortex $\nu$ with $|\operatorname{supp} \nu| \geq 14$. Thus, $V_{3}=V_{2}$ and $W_{\gamma}^{3}=W_{\gamma}^{0}$. Similarly as earlier, if $A_{2}$ does not occur, then $\left|W_{\gamma}^{0}-W_{\gamma}^{3}\right| \leq 2$. By Lemma 6.7, the expectation value is

$$
\begin{equation*}
\mathbb{E}_{\beta, N}\left[\left|W_{\gamma}^{0}-W_{\gamma}^{3}\right|\right] \leq 0 \cdot \mu_{\beta, N}\left(A_{2}\right)+2 \cdot \mu_{\beta, N}\left(A_{2}^{c}\right) \leq 2 C_{1} C_{0}^{(7)} \ell \lambda(\beta)^{14} \tag{6.25}
\end{equation*}
$$

## The third expectation $\left(\mathbb{E}_{\beta, N}\left[\left|W_{\gamma}^{3}-W_{\gamma}^{4}\right|\right]\right)$

Let

$$
V_{4}:=\left\{\nu \in V_{3}: \exists e \in \gamma \text { such that supp } \nu=\hat{\partial} e \cup \hat{\partial}(-e)\right\}
$$

and

$$
W_{\gamma}^{4}:=\rho\left(\sum_{\nu \in V_{4}} \sum_{p \in \operatorname{supp} \nu} q_{p}^{+}(d \sigma)_{p}\right) .
$$

Consider $\nu \in V_{3} \backslash V_{4}$. Since $\nu$ is minimal, by Lemma 4.7 there exists an edge $e \in E_{N}$ such that supp $\nu=\hat{\partial} e \cup \hat{\partial}(-e)$. Furthermore, since $\nu \notin V_{4}$, the edge $e$ cannot be in the boundary of $q$. Thus, $e$ is an internal edge of $q$. By Lemma 4.16

$$
\sum_{p \in \operatorname{supp} \nu} q_{p}^{+}(d \sigma)_{p}=0 \quad \text { for } v \in V_{3} \backslash V_{4} .
$$

Therefore,

$$
W_{\gamma}^{3}=\rho\left(\sum_{\nu \in V_{3}} \sum_{p \in \operatorname{supp} \nu} q_{p}^{+}(d \sigma)_{p}\right)=\rho\left(\sum_{\nu \in V_{4}} \sum_{p \in \operatorname{supp} \nu} q_{p}^{+}(d \sigma)_{p}\right)=W_{\gamma}^{4} .
$$

Since $W_{\gamma}^{3}=W_{\gamma}^{4}$, the expectation of $\left|W_{\gamma}^{3}-W_{\gamma}^{4}\right|$ is

$$
\begin{equation*}
\mathbb{E}_{\beta, N}\left[\left|W_{\gamma}^{3}-W_{\gamma}^{4}\right|\right]=\mathbb{E}_{\beta, N}[0]=0 \tag{6.26}
\end{equation*}
$$

The fourth expectation $\left(\mathbb{E}_{\beta, N}\left[\left|W_{\gamma}^{4}-W_{\gamma}^{5}\right|\right]\right)$
Let

$$
V_{5}:=\left\{\nu \in V_{4}: \exists e \in \gamma_{1} \text { such that supp } \nu=\hat{\partial} e \cup \hat{\partial}(-e)\right\}
$$

and

$$
W_{\gamma}^{5}:=\rho\left(\sum_{\nu \in V_{5}} \sum_{p \in \operatorname{supp} \nu} q_{p}^{+}(d \sigma)_{p}\right) .
$$

If the event $A_{3}$ occurs, then $V_{4}=V_{5}$ and $W_{\gamma}^{4}=W_{\gamma}^{5}$. If $A_{3}$ does not occur, then $\left|W_{\gamma}^{4}-W_{\gamma}^{5}\right| \leq 2$. Thus, by Lemma 6.8, the expectation value of $\left|W_{\gamma}^{4}-W_{\gamma}^{5}\right|$ is

$$
\begin{equation*}
\mathbb{E}_{\beta, N}\left[\left|W_{\gamma}^{4}-W_{\gamma}^{5}\right|\right] \leq 0 \cdot \mu_{\beta, N}\left(A_{3}\right)+2 \mu_{\beta, N}\left(A_{3}^{c}\right) \leq 2 C^{(6)} \ell_{c} \lambda(\beta)^{12} \tag{6.27}
\end{equation*}
$$

The fifth expectation $\left(\mathbb{E}_{\beta, N}\left[\left|W_{\gamma}^{5}-W_{\gamma}^{6}\right|\right]\right)$
We first define a set of edges:

$$
E_{5}:=\left\{e \in \gamma: \exists \nu \in V_{5} \text { such that supp } \nu=\hat{\partial} e \cup \hat{\partial}(-e)\right\} .
$$

Then $W_{\gamma}^{5}$ can be written as

$$
\begin{aligned}
W_{\gamma}^{5} & =\rho\left(\sum_{\nu \in V_{5}} \sum_{p \in \operatorname{supp} \nu} q_{p}^{+}(d \sigma)_{p}\right)=\rho\left(\sum_{e \in E_{5}} \sum_{p \in \pm \hat{\partial}_{e}} q_{p}^{+}(d \sigma)_{p}\right) \\
& =\rho\left(\sum_{e \in E_{5}} \sum_{p \in \hat{\partial} e} q_{p}(d \sigma)_{p}\right) .
\end{aligned}
$$

Define

$$
E_{6}:=\left\{e \in \gamma_{1}:(d \sigma)_{p}=(d \sigma)_{p^{\prime}} \text { for all } p, p^{\prime} \in \hat{\partial}(-e)\right\}
$$

and

$$
W_{\gamma}^{6}:=\rho\left(\sum_{e \in E_{6}} \sum_{p \in \hat{\partial} e} q_{p}^{+}(d \sigma)_{p}\right) .
$$

Notice that the set $E_{6}$ can be written as

$$
\begin{equation*}
E_{6}=\gamma_{1} \backslash \gamma^{\prime} \tag{6.28}
\end{equation*}
$$

Consider two disjoint edges $e, e^{\prime} \in \gamma_{1}$, then $\hat{\partial} e \cup \hat{\partial}(-e)$ and $\hat{\partial} e^{\prime} \cup \hat{\partial}\left(-e^{\prime}\right)$ are disjoint. Since any vortex $\nu \in V_{5}$ is minimal, we have by Lemma 4.9 for a vortex $\nu \in V_{5}$ that $(d \sigma)_{p}=d\left(g d x_{j}\right)_{p}=\{-g, g\}$ for all $p \in P_{N}$. Thus, $(d \sigma)_{p}$ does not depend on $p \in \hat{\partial} e$. For $e \in E_{6}$, fix $p_{e} \in \hat{\partial} e$. Since $q$ is an oriented surface and $E_{6}$ only contains edges in the boundary of $q$, we obtain

$$
\sum_{p \in \hat{\partial} e} q_{p}(d \sigma)_{p}=\sum_{p \in \hat{\partial} e} q_{p}(d \sigma)_{p_{e}}=\left(\sum_{p \in \hat{\partial} e} q_{p}\right)(d \sigma)_{p_{e}}=(d \sigma)_{p_{e}}, \quad \text { for } e \in E_{6}
$$

Therefore,

$$
W_{\gamma}^{6}=\rho\left(\sum_{e \in E_{6}}(d \sigma)_{p_{e}}\right)
$$

If the event $A_{2}$ occurs, then $E_{5}=E_{6}$ and $W_{\gamma}^{5}=W_{\gamma}^{6}$. If $A_{2}$ does not occur, then $\left|W_{\gamma}^{5}-W_{\gamma}^{6}\right| \leq 2$. Hence, by Lemma 6.7.

$$
\begin{equation*}
\mathbb{E}_{\beta, N}\left[\left|W_{\gamma}^{5}-W_{\gamma}^{6}\right|\right] \leq 0 \cdot \mu_{\beta, N}\left(A_{2}\right)+2 \mu_{\beta, N}\left(A_{2}^{c}\right)=2 C_{1} C_{0}^{(7)} \ell \lambda(\beta)^{14} \tag{6.29}
\end{equation*}
$$

The sixth expectation $\left(\left|\mathbb{E}_{\beta, N}\left[W_{\gamma}^{6}\right]-\theta^{\ell}\right|\right)$
The strategy will be to first compute the conditional expectation of $W_{\gamma}^{6}$ given $\left(\sigma_{e}\right)_{e \neq \pm \gamma_{1}}$ and then apply the law of total expectation. Given $\left(\sigma_{e}\right)_{e \notin \pm \gamma_{1}}$, denote the conditional probability by $\mu_{\beta, N}^{\prime}$ and the conditional expectation by $\mathbb{E}_{\beta, N}^{\prime}$.

We calculate the conditional expectation value of $W_{\gamma}^{6}$ given $\left(\sigma_{e}\right)_{e \neq \pm \gamma_{1}}$. By Lemma 6.9, we have that the spins $\left(\sigma_{e}\right)_{e \in \gamma_{1}}$ are independent. For $e \in \gamma_{1} \backslash \gamma^{\prime}$, define $\sigma^{e}:=\sum_{\hat{e} \in p_{e} \backslash\{e\}} \sigma_{\hat{e}}$ such that

$$
\begin{equation*}
(d \sigma)_{p}=\sigma^{e}+\sigma_{e} \quad \text { for each } p \in \hat{\partial} e \tag{6.30}
\end{equation*}
$$

Then,

$$
\begin{aligned}
& \mathbb{E}_{\beta, N}^{\prime}\left[W_{\gamma}^{6}\right]=\mathbb{E}_{\beta, N}^{\prime}\left[\rho\left(\sum_{e \in E_{6}}(d \sigma)_{p_{e}}\right)\right] \\
& \stackrel{\sqrt{6.28}}{=} \mathbb{E}_{\beta, N}^{\prime}\left[\rho\left(\sum_{e \in \gamma_{1} \backslash \gamma^{\prime}}(d \sigma)_{p_{e}}\right)\right] \\
&=\mathbb{E}_{\beta, N}^{\prime}\left[\prod_{e \in \gamma_{1} \backslash \gamma^{\prime}} \rho\left((d \sigma)_{p_{e}}\right)\right] \\
& \Perp \prod_{e \in \gamma_{1} \backslash \gamma^{\prime}} \mathbb{E}_{\beta, N}^{\prime}\left[\rho\left((d \sigma)_{p_{e}}\right)\right] \\
& \stackrel{\boxed{66.30}}{=} \prod_{e \in \gamma_{1} \backslash \gamma^{\prime}} \mathbb{E}_{\beta, N}^{\prime}\left[\rho\left(\sigma^{e}+\sigma_{e}\right)\right] \\
&=\prod_{e \in \gamma_{1} \backslash \gamma^{\prime}}\left(\sum_{g \in G} \rho\left(\sigma^{e}+g\right) \frac{\phi_{\beta}\left(\sigma^{e}+g\right)^{12}}{\sum_{g \in G} \phi_{\beta}\left(\sigma^{e}+g\right)^{12}}\right) \\
&=\prod_{e \in \gamma_{1} \backslash \gamma^{\prime}}\left(\frac{\sum_{g \in G} \rho(g) \phi_{\beta}(g)^{12}}{\sum_{g \in G} \phi_{\beta}(g)^{12}}\right) \\
&=\left(\frac{\sum_{g \in G} \rho(g) \phi_{\beta}(g)^{12}}{\sum_{g \in G} \phi_{\beta}(g)^{12}}\right)
\end{aligned}
$$

$$
=\theta(\beta)^{\left|\gamma_{1} \wedge \gamma^{\prime}\right|} .
$$

By the law of total expectation, we have

$$
\begin{align*}
\mathbb{E}_{\beta, N}\left[W_{\gamma}^{6}\right] & =\mathbb{E}_{\beta, N}\left[\mathbb{E}_{\beta, N}^{\prime}\left[W_{\gamma}^{6}\right]\right]=\mathbb{E}_{\beta, N}\left[\theta(\beta)^{\left|\gamma_{1} \backslash \gamma^{\prime}\right|}\right]=\mathbb{E}_{\beta, N}\left[\theta(\beta)^{\left|\gamma_{1}\right|} \theta(\beta)^{-\left|\gamma^{\prime}\right|}\right]  \tag{6.31}\\
& =\theta(\beta)^{\left|\gamma_{1}\right|} \mathbb{E}_{\beta, N}\left[\theta(\beta)^{-\left|\gamma^{\prime}\right|}\right]
\end{align*}
$$

Note that since $\theta \leq 1$, it follows that

$$
\begin{equation*}
\theta(\beta)^{\left|\gamma_{1}\right|} \leq 1 \tag{6.32}
\end{equation*}
$$

Since $\left|\gamma_{1}\right|$ is the number of non-corner edges in $\gamma$ and $\ell_{c}$ is the number of corner edges in $\gamma$, we have

$$
\begin{equation*}
\left|\gamma_{1}\right|+\ell_{c}=\ell \tag{6.33}
\end{equation*}
$$

Let $j$ be such that $0 \leq j \leq \ell$. Then,

$$
\begin{align*}
& \left|\mathbb{E}_{\beta, N}\left[W_{\gamma}^{6}\right]-\theta^{\ell} \stackrel{[6.31}{=} \theta^{\left|\gamma_{1}\right|} \mathbb{E}_{\beta, N}\left[\theta^{-\left|\gamma^{\prime}\right|}\right]-\theta^{\ell}\right| \\
& \stackrel{\sqrt{6.33}}{=}\left|\theta^{\left|\gamma_{1}\right|}\right| \mathbb{E}_{\beta, N}\left[\theta^{-\left|\gamma^{\prime}\right|}\right]-\theta^{\left|\gamma_{1}\right|}+\theta^{\left|\gamma_{1}\right|}-\theta^{\left|\gamma_{1}\right|+\ell_{c}} \mid \\
& =\left|\theta^{\left|\gamma_{1}\right|}\left(\mathbb{E}_{\beta, N}\left[\theta^{-\left|\gamma^{\prime}\right|}\right]-1\right)-\theta^{\left|\gamma_{1}\right|}\left(\theta^{\ell_{c}}-1\right)\right| \\
& \leq\left|\theta^{\left|\gamma_{1}\right|}\left(\mathbb{E}_{\beta, N}\left[\theta^{-\left|\gamma^{\prime}\right|}\right]-1\right)\right|+\left|\theta^{\left|\gamma_{1}\right|}\left(\theta^{\ell_{c}}-1\right)\right| \\
& \stackrel{\sqrt{6.32}}{\leq}\left|\mathbb{E}_{\beta, N}\left[\theta^{-\left|\gamma^{\prime}\right|}\right]-1\right|+\left|\theta^{\ell_{c}}-1\right| \\
& \stackrel{\sqrt{6.12} \mid}{\leq}\left|\mathbb{E}_{\beta, N}\left[\theta^{-\left|\gamma^{\prime}\right|}\right]-1\right|+2 C^{*} \ell_{c} \lambda(\beta)^{12} e^{C^{*} \ell \lambda(\beta)^{12}} \\
& =\left|\mathbb{E}_{\beta, N}\left[\theta^{-\left|\gamma^{\prime}\right|}\right]-\mathbb{E}_{\beta, N}\left[\theta^{-\left|\gamma^{\prime}\right|} \mathbb{I}_{\left\{\left|\gamma^{\prime}\right| \leq j\right\}}\right]+\mathbb{E}_{\beta, N}\left[\theta^{-\left|\gamma^{\prime}\right|} \mathbb{I}_{\left\{\left|\gamma^{\prime}\right| \leq j\right\}}\right]-1\right| \\
& +2 C^{*} \ell_{c} \lambda(\beta)^{12} e^{C^{*} \ell \lambda(\beta)^{12}} \\
& \leq\left|\mathbb{E}_{\beta, N}\left[\theta^{-\left|\gamma^{\prime}\right|}\right]-\mathbb{E}_{\beta, N}\left[\theta^{-\left|\gamma^{\prime}\right|} \mathbb{I}_{\left\{\left|\gamma^{\prime}\right| \leq j\right\}}\right]\right|+\left|\mathbb{E}_{\beta, N}\left[\theta^{-\left|\gamma^{\prime}\right|} \mathbb{I}_{\left\{\left|\gamma^{\prime}\right| \leq j\right\}}\right]-1\right| \\
& +2 C^{*} \ell_{c} \lambda(\beta)^{12} e^{C^{*} \ell \lambda(\beta)^{12}} \tag{6.34}
\end{align*}
$$

We focus on the first term on the right-hand side above:

$$
\begin{align*}
\left|\mathbb{E}_{\beta, N}\left[\theta^{-\left|\gamma^{\prime}\right|}\right]-\mathbb{E}_{\beta, N}\left[\theta^{-\left|\gamma^{\prime}\right|} \mathbb{I}_{\left\{\left|\gamma^{\prime}\right| \leq j\right\}}\right]\right| & =\mathbb{E}_{\beta, N}\left[\theta^{-\left|\gamma^{\prime}\right|} \mathbb{I}_{\left\{\left|\gamma^{\prime}\right|>j\right\}}\right] \\
& \leq \theta^{-\ell} \mu_{\beta, N}\left(\left\{\sigma \in \Sigma_{E_{N}}:\left|\gamma^{\prime}\right|>j\right\}\right) \\
& \leq \frac{\theta^{-\ell} \mathbb{E}_{\beta, N}\left[\left|\gamma^{\prime}\right|\right]}{j}  \tag{6.35}\\
& \stackrel{\sqrt{6.23}}{\leq} \frac{\theta^{-\ell} C_{0}^{(6)} \ell \lambda(\beta)^{12}}{j} \\
& \stackrel{(6.11]}{\leq} \frac{2 C_{0}^{(6)} \ell \lambda(\beta)^{12} e^{C^{*} \ell \lambda(\beta)^{12}}}{j}
\end{align*}
$$

Then we focus on the second term:

$$
\left.\begin{array}{rl}
\left|\mathbb{E}_{\beta, N}\left[\theta^{-\left|\gamma^{\prime}\right|} \mathbb{I}_{\left\{\left|\gamma^{\prime}\right| \leq j\right\}}\right]-1\right| & \leq\left|\mathbb{E}_{\beta, N}\left[\left(\theta^{-\left|\gamma^{\prime}\right|}-1\right) \mathbb{I}_{\left\{\left|\gamma^{\prime}\right| \leq j\right\}}\right]\right|+\mathbb{E}_{\beta, N}\left[\mathbb{I}_{\left\{\left|\gamma^{\prime}\right| \leq j\right\}}\right] \\
& \leq\left(\theta^{-j}-1\right)+\frac{\mathbb{E}_{\beta, N}\left[\left|\gamma^{\prime}\right|\right]}{j}  \tag{6.36}\\
\frac{\sqrt{6.12 \mid}, \sqrt[(6.23 \mid]{\leq}}{}{ }^{*} *
\end{array}\right](\beta)^{12} e^{C^{*} \ell \lambda(\beta)^{12}}+\frac{C_{0}^{(6)} \ell \lambda(\beta)^{12}}{j} .
$$

We combine these two inequalities and choose $j=\sqrt{\ell}$ :

$$
\begin{align*}
& \left|\mathbb{E}_{\beta, N}\left[\theta^{-\left|\gamma^{\prime}\right|}\right]-\mathbb{E}_{\beta, N}\left[\theta^{-\left|\gamma^{\prime}\right|} \mathbb{I}_{\left\{\left|\gamma^{\prime}\right| \leq j\right\}}\right]\right|+\left|\mathbb{E}_{\beta, N}\left[\theta^{-\left|\gamma^{\prime}\right|} \mathbb{I}_{\left\{\left|\gamma^{\prime}\right| \leq j\right\}}\right]-1\right| \\
& \stackrel{(6.35),(6.36}{=} \frac{2 C_{0}^{(6)} \ell \lambda(\beta)^{12} e^{C^{*} \ell \lambda(\beta)^{12}}}{j}+2 C^{*} j \lambda(\beta)^{12} e^{C^{*} \ell \lambda(\beta)^{12}}+\frac{C_{0}^{(6)} \ell \lambda(\beta)^{12}}{j}  \tag{6.37}\\
& \stackrel{j=\sqrt{\ell}}{=} 2 C_{0}^{(6)} \sqrt{\ell} \lambda(\beta)^{12} e^{C^{*} \ell \lambda(\beta)^{12}}+2 C^{*} \sqrt{\ell} \lambda(\beta)^{12} e^{C^{*} \ell \lambda(\beta)^{12}}+C_{0}^{(6)} \sqrt{\ell} \lambda(\beta)^{12} \\
& \quad=2\left(C_{0}^{(6)}+C^{*}\right) \sqrt{\ell} \lambda(\beta)^{12} e^{C^{*} \ell \lambda(\beta)^{12}}+C_{0}^{(6)} \sqrt{\ell} \lambda(\beta)^{12} .
\end{align*}
$$

We return to inequality (6.34) and obtain

$$
\begin{align*}
\mid \mathbb{E}_{\beta, N}\left[W_{\gamma}^{6}\right]-\theta^{\ell} \stackrel{\mid 6.37] ~}{\mid 6.37}_{\leq} & 2\left(C_{0}^{(6)}+C^{*}\right) \sqrt{\ell} \lambda(\beta)^{12} e^{C^{*} \ell \lambda(\beta)^{12}}+C_{0}^{(6)} \sqrt{\ell} \lambda(\beta)^{12} \\
& +2 \ell_{c} C^{*} \lambda(\beta)^{12} e^{c^{*} \ell \lambda(\beta)^{12}} \\
= & 2\left(\frac{C_{0}^{(6)}+C^{*}}{\sqrt{\ell}}+\frac{C^{*} \ell_{c}}{\ell}\right) \ell \lambda(\beta)^{12} e^{C^{*} \ell \lambda(\beta)^{12}}+C_{0}^{(6)} \sqrt{\ell} \lambda(\beta)^{12} . \tag{6.38}
\end{align*}
$$

## Combining the expectation values

We combine equations (6.24), (6.25), (6.26), (6.27), (6.29) and (6.38):

$$
\begin{aligned}
\left|\mathbb{E}_{\beta, N}\left[W_{\gamma}\right]-\theta^{\ell}\right| \leq & \left|\mathbb{E}_{\beta, N}\left[W_{\gamma}^{6}\right]-\theta^{\ell}\right|+\mathbb{E}_{\beta, N}\left[\left|W_{\gamma}^{6}-W_{\gamma}^{5}\right|\right]+\mathbb{E}_{\beta, N}\left[\left|W_{\gamma}^{5}-W_{\gamma}^{4}\right|\right] \\
& +\mathbb{E}_{\beta, N}\left[\left|W_{\gamma}^{4}-W_{\gamma}^{3}\right|\right]+\mathbb{E}_{\beta, N}\left[\left|W_{\gamma}^{3}-W_{\gamma}^{0}\right|\right]+\mathbb{E}_{\beta, N}\left[\left|W_{\gamma}^{0}-W_{\gamma}\right|\right] \\
\leq & 2\left(\frac{C_{0}^{(6)}+C^{*}}{\sqrt{\ell}}+\frac{C^{*} \ell_{c}}{\ell}\right) \ell \lambda(\beta)^{12} e^{C^{*} \ell \lambda(\beta)^{12}}+C_{0}^{(6)} \sqrt{\ell} \lambda(\beta)^{12} \\
& +2 C_{1} C_{0}^{(7)} \ell \lambda(\beta)^{14}+2 C_{0}^{(6)} \ell_{c} \lambda(\beta)^{12}+0+2 C_{1} C_{0}^{(7)} \ell \lambda(\beta)^{14} \\
& +2 C_{0}^{(25)} \ell^{4} \lambda(\beta)^{50} \\
= & 2\left(\frac{C_{0}^{(6)}+C^{*}}{C^{*} \sqrt{\ell}}+\frac{\ell_{c}}{\ell}\right) C^{*} \ell \lambda(\beta)^{12} e^{C^{*} \ell \lambda(\beta)^{12}} \\
& +\left(\frac{C_{0}^{(6)}}{2 C^{*} \sqrt{\ell}}+\frac{2 C_{1} C_{0}^{(7)} \lambda(\beta)^{2}}{C^{*}}+\frac{C_{0}^{(6)} \ell}{C^{*} \ell}\right) 2 C^{*} \ell \lambda(\beta)^{12} \\
& +\frac{C_{0}^{(25)} \lambda(\beta)^{2}}{2^{3}\left(C^{*}\right)^{4}}\left(2 C^{*} \ell \lambda(\beta)^{12}\right)^{4} .
\end{aligned}
$$

Note that for $x>0$ the two following inequalities hold:

$$
x \leq e^{x} \text { and } x^{4} \leq 5 e^{x}
$$

Hence,

$$
\begin{gathered}
C^{*} \ell \lambda(\beta)^{12} e^{C^{*} \ell \lambda(\beta)^{12}} \leq e^{2 C^{*} \ell \lambda(\beta)^{12}} \\
2 C^{*} \ell \lambda(\beta)^{12} \leq e^{2 C^{*} \ell \lambda(\beta)^{12}}
\end{gathered}
$$

and

$$
\left(2 C^{*} \ell \lambda(\beta)^{12}\right)^{4} \leq 5 e^{2 C^{*} \ell \lambda(\beta)^{12}}
$$

Furthermore, since $1 \leq \ell_{c} \leq \ell$, we have

$$
\begin{equation*}
\frac{1}{\sqrt{\ell}} \leq \sqrt{\frac{\ell_{c}}{\ell}} \text { and } \frac{\ell_{c}}{\ell} \leq \sqrt{\frac{\ell_{c}}{\ell}} \tag{6.39}
\end{equation*}
$$

Therefore,

$$
\begin{align*}
&\left|\mathbb{E}_{\beta, N}\left[W_{\gamma}\right]-\theta^{\ell}\right| \\
& \leq\left(2 \frac{C_{0}^{(6)}+C^{*}}{C^{*} \sqrt{\ell}}+\frac{2 \ell_{c}}{\ell}+\frac{C_{0}^{(6)}}{2 C^{*} \sqrt{\ell}}+\frac{2 C_{1} C_{0}^{(7)} \lambda(\beta)^{2}}{C^{*}}+\frac{C_{0}^{(6)} \ell_{c}}{C^{*} \ell}+\frac{5 C_{0}^{(25)} \lambda(\beta)^{2}}{2^{3}\left(C^{*}\right)^{4}}\right) \\
& \cdot e^{2 C^{*} \ell \lambda(\beta)^{12}} \\
&=\left(\frac{5 C_{0}^{(6)}+4 C^{*}}{2 C^{*} \sqrt{\ell}}+\frac{2 \ell_{c}}{\ell}+\frac{2 C_{1} C_{0}^{(7)} \lambda(\beta)^{2}}{C^{*}}+\frac{C_{0}^{(6)} \ell_{c}}{C^{*} \ell}+\frac{5 C_{0}^{(25)} \lambda(\beta)^{2}}{2^{3}\left(C^{*}\right)^{4}}\right) \cdot e^{2 C^{*} \ell \lambda(\beta)^{12}} \\
& \frac{\sqrt[66.39]{\leq}}{\leq}\left(\frac{5 C_{0}^{(6)}+4 C^{*}}{2 C^{*}}+2+\frac{2 C_{1} C_{0}^{(7)}}{C^{*}}+\frac{C_{0}^{(6)}}{C^{*}}+\frac{5 C_{0}^{(25)}}{2^{3}\left(C^{*}\right)^{4}}\right)\left(\sqrt{\frac{\ell_{c}}{\ell}}+\lambda(\beta)^{2}\right) \\
&=\left(\frac{7 C_{0}^{(6)}}{2 C^{*}}+4+\frac{2 C_{1} C_{0}^{(7)}}{C^{*}}+\frac{5 C_{0}^{(25)}}{2^{3}\left(C^{*}\right)^{4}}\right)\left(\sqrt{\frac{\ell_{c}}{\ell}}+\lambda(\beta)^{2}\right) e^{2 C^{*} \ell \lambda(\beta)^{12}} \\
&=\left(C_{A}-1 / 2\right)\left(\sqrt{\frac{\ell_{c}}{\ell}}+\lambda(\beta)^{2}\right) e^{2 C^{*} \ell \lambda(\beta)^{12}},
\end{align*}
$$

where

$$
\begin{equation*}
C_{A}=\frac{7 C_{0}^{(6)}}{2 C^{*}}+\frac{2 C_{1} C_{0}^{(7)}}{C^{*}}+\frac{5 C_{0}^{(25)}}{2^{3}\left(C^{*}\right)^{4}}+\frac{9}{2} \tag{6.41}
\end{equation*}
$$

Finally, we can calculate the last inequality:

$$
\begin{aligned}
\left|\mathbb{E}_{\beta, N}\left[W_{\gamma}\right]-e^{-\ell(1-\theta(\beta))}\right| & =\left|\mathbb{E}_{\beta, N}\left[W_{\gamma}\right]-\theta^{\ell}+\theta^{\ell}-e^{-\ell(1-\theta(\beta))}\right| \\
& \leq\left|\mathbb{E}_{\beta, N}\left[W_{\gamma}\right]-\theta^{\ell}\right|+\left|\theta^{\ell}-e^{-\ell(1-\theta(\beta))}\right|
\end{aligned}
$$

$$
\begin{aligned}
&\left.\frac{\sqrt{6.13), \sqrt{6.40}}}{\left(C_{A}\right.}-1 / 2\right)\left(\sqrt{\frac{\ell_{c}}{\ell}}+\lambda(\beta)^{2}\right) e^{2 C^{*} \ell \lambda(\beta)^{12}}+\frac{e^{2 C^{*} \ell \lambda(\beta)^{12}}}{2 \ell} \\
&=\left(C_{A}-1 / 2+\frac{1}{2} \cdot \frac{1}{\ell\left(\sqrt{\frac{\ell_{c}}{\ell}}+\lambda(\beta)^{2}\right)}\right)\left(\sqrt{\frac{\ell_{c}}{\ell}}+\lambda(\beta)^{2}\right) \\
& \cdot e^{2 C^{*} \ell \lambda(\beta)^{12}} \\
& \quad \leq\left(C_{A}-1 / 2+\frac{1}{2}\right)\left(\sqrt{\frac{\ell_{c}}{\ell}}+\lambda(\beta)^{2}\right) e^{2 C^{*} \ell \lambda(\beta)^{12}} \\
&= C_{A}\left(\sqrt{\frac{\ell_{c}}{\ell}}+\lambda(\beta)^{2}\right) e^{2 C^{*} \ell \lambda(\beta)^{12}} .
\end{aligned}
$$

### 6.2 The case when $\ell \lambda(\beta)^{12}$ is large

We now consider the case when $\ell \lambda(\beta)^{12}$ is large. First, some notations and definitions of importance are stated.

Let the set $K$ be a finite and non-empty index set. The arguments of the maxima are the elements for which the value of the function is maximised and is denoted by arg max. Likewise, the arguments of the minima are the elements for which the value of the function is minimised and denoted by arg min. Given $g_{k} \in G$, define for each $K$, the set

$$
G_{0}\left[\left(g_{k}\right)_{k \in K}\right]:=\underset{g \in G}{\arg \max } \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right) .
$$

Next, an assumption is given and we prove that it is satisfied if $\beta_{0}$ is chosen large enough.
$(\star)$ For all $\beta \geq \beta_{0}$ and choices of $g_{k} \in G$ for $k \in K$ and any $g^{\prime} \in G_{0}\left[\left(g_{k}\right)_{k \in K}\right]$, the following inequality holds:

$$
\sum_{g \in G \backslash G_{0}\left[\left(g_{k}\right)_{k \in K}\right]} \prod_{k \in K} \frac{\phi_{\beta}\left(g+g_{k}\right)^{2}}{\phi_{\beta}\left(g^{\prime}+g_{k}\right)^{2}} \leq \frac{1-\cos (2 \pi / n)}{8} .
$$

Let $\beta \geq 0$ and $g \in G$. Since

$$
\prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)=\prod_{k \in K} e^{\beta \Re\left(\rho\left(g+g_{k}\right)\right)}=e^{\beta \sum_{k \in K} \Re\left(\rho\left(g+g_{k}\right)\right)},
$$

we obtain

$$
G_{0}\left[\left(g_{k}\right)_{k \in K}\right]=\underset{g \in G}{\arg \max } e^{\beta \sum_{k \in K} \Re\left(\rho\left(g+g_{k}\right)\right)}=\underset{g \in G}{\arg \max } \sum_{k \in K} \Re\left(\rho\left(g+g_{k}\right)\right) .
$$

We see that $G_{0}\left[\left(g_{k}\right)_{k \in K}\right]$ does not depend on $\beta$. Let $g^{\prime} \in G_{0}\left[\left(g_{k}\right)_{k \in K}\right]$. Then $g^{\prime}$ is one of the elements in $G$ for which $\sum_{k \in K} \Re\left(\rho\left(g+g_{k}\right)\right)$ takes its maximum value. For $g \in G \backslash G_{0}\left[\left(g_{k}\right)_{k \in K}\right]$, we have

$$
\sum_{k \in K} \Re\left(\rho\left(g^{\prime}+g_{k}\right)\right)>\sum_{k \in K} \Re\left(\rho\left(g+g_{k}\right)\right) .
$$

Thus, by taking the limit when $\beta \rightarrow \infty$ we obtain

$$
\lim _{\beta \rightarrow \infty} \prod_{k \in K} \frac{\phi_{\beta}\left(g+g_{k}\right)^{2}}{\phi_{\beta}\left(g^{\prime}+g_{k}\right)^{2}}=\lim _{\beta \rightarrow \infty} \exp \left(2 \beta \sum_{k \in K}\left(\Re \rho\left(g+g_{k}\right)-\Re \rho\left(g^{\prime}+g_{k}\right)\right)\right)=0 .
$$

Thus, there exists a $\beta_{0}$ for which ( $\star$ ) holds. This implies that there exists a $\beta_{0}$ for which $(\star)$ holds for all $K$ simultaneously, since $\beta_{0}$ can be chosen as the maximum of $\beta_{0}$ for all fixed $K$.

Proposition 6.11. Consider lattice gauge theory with structure group $G=\mathbb{Z}_{n}$ and a one-dimensional faithful representation $\rho$ of $G$. Let $\gamma$ be a simple oriented loop in $\mathbb{Z}^{4}, \ell=|\gamma|$ the length of it and $\ell_{c}$ the number of corner edges in $\gamma$. Let $N$ be large enough so that the edges of $\gamma$ are internal edges of $B_{N}$. Let $\beta_{0}>0$ satisfy the assumption $(\star)$ when applied with sets $K$ with $|K|=6$, and be such that $2 \lambda\left(\beta_{0}\right)^{2|K|} \leq 1$. Then for all $\beta>\beta_{0}$, we have

$$
\begin{equation*}
\left|\mathbb{E}_{\beta, N}\left[W_{\gamma}\right]\right| \leq e^{-C_{*}\left(\ell-\ell_{c}\right) \lambda(\beta)^{12}}, \quad \text { where } C_{*}:=\frac{1-\cos (2 \pi / n)}{4} \tag{6.42}
\end{equation*}
$$

To prove this proposition, three lemmas will be applied. We begin with stating the main lemma of this section, which will require two lemmas to be proved.

Lemma 6.12. Let $K$ be a finite and non-empty index set and assume that $\beta_{0}=$ $\beta_{0}(|K|)>0$ satisfies $(\star)$ applied with $K$ and

$$
\begin{equation*}
2 \lambda\left(\beta_{0}\right)^{2|K|} \leq 1 \tag{6.43}
\end{equation*}
$$

Then, for all choices of $g_{k} \in G$ for $k \in K$ and in the setting of Lemma 6.11, we have

$$
\begin{equation*}
\left|\frac{\sum_{g \in G} \rho(g) \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}{\sum_{g \in G} \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}\right| \leq 1-C_{*} \lambda(\beta)^{2|K|} \tag{6.44}
\end{equation*}
$$

The left-hand side of inequality (6.44) is the conditional expected value of a single spin (see the beginning of the proof of this proposition). To lighten notations, we define

$$
\begin{equation*}
S_{\beta}\left(\left(g_{k}\right)_{k \in K}\right):=\frac{\sum_{g \in G} \rho(g) \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}{\sum_{g \in G} \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}} \tag{6.45}
\end{equation*}
$$

and

$$
\begin{equation*}
\omega:=1-S_{\beta}\left(\left(g_{k}\right)_{k \in K}\right) . \tag{6.46}
\end{equation*}
$$

We state and prove two lemmas for the case when the set $G_{0}\left[\left(g_{k}\right)_{k \in K}\right]$ only consists of one element (We will later prove that $G_{0}\left[\left(g_{k}\right)_{k \in K}\right]=\{0\}$ in this case.). Thereafter, the proof of Lemma 6.12 is given.

Lemma 6.13. Let $K$ be a finite and non-empty index set and assume that $\beta_{0}=$ $\beta_{0}(|K|)>0$ satisfies $(\star)$ applied with $K$. If $G_{0}\left[\left(g_{k}\right)_{k \in K}\right]=\{0\}$, then

$$
\begin{align*}
& \text { (i) }|\omega|^{2} \leq \Re \omega,  \tag{6.47}\\
& \text { (ii) }\left|S_{\beta}\left(\left(g_{k}\right)_{k \in K}\right)\right| \leq 1-C_{*} \sum_{g \in G \backslash\{0\}} \frac{\prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}{\prod_{k \in K} \phi_{\beta}\left(0+g_{k}\right)^{2}} . \tag{6.48}
\end{align*}
$$

Proof. (i) From equations (6.45) and (6.46) follows that

$$
\begin{align*}
\omega & =1-\frac{\sum_{g \in G} \rho(g) \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}{\sum_{g \in G} \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}} \\
& =\frac{\sum_{g \in G}(1-\rho(g)) \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}{\sum_{g \in G} \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}} \\
& =\frac{\sum_{g \in G}(1-\rho(g)) \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}{\prod_{k \in K} \phi_{\beta}\left(0+g_{k}\right)^{2}} \cdot \frac{\prod_{k \in K} \phi_{\beta}\left(0+g_{k}\right)^{2}}{\sum_{g \in G} \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}  \tag{6.49}\\
& =\left(\sum_{g \in G}(1-\rho(g)) \prod_{k \in K} \frac{\phi_{\beta}\left(g+g_{k}\right)^{2}}{\phi_{\beta}\left(0+g_{k}\right)^{2}}\right) \cdot \frac{\prod_{k \in K} \phi_{\beta}\left(0+g_{k}\right)^{2}}{\sum_{g \in G} \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}} .
\end{align*}
$$

Since $G_{0}\left[\left(g_{k}\right)_{k \in K}\right]=\{0\}$, by $(\star)$ we have

$$
\begin{equation*}
\sum_{g \in G \backslash\{0\}} \prod_{k \in K} \frac{\phi_{\beta}\left(g+g_{k}\right)^{2}}{\phi_{\beta}\left(0+g_{k}\right)^{2}} \leq \frac{1-\cos (2 \pi / n)}{8} . \tag{6.50}
\end{equation*}
$$

We focus on the second factor in (6.49):

$$
\begin{align*}
& 1 \geq \frac{\prod_{k \in K} \phi_{\beta}\left(0+g_{k}\right)^{2}}{\sum_{g \in G} \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}} \\
& =1-\frac{\sum_{g \in G \backslash\{0\}} \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}{\sum_{g \in G} \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}} \\
& >1-\frac{\sum_{g \in G \backslash\{0\}} \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}{\prod_{k \in K} \phi_{\beta}\left(0+g_{k}\right)^{2}}  \tag{6.51}\\
& \stackrel{\sqrt[6.50)]{\geq}}{\geq 1}-\frac{1-\cos (2 \pi / n)}{8} \\
& =\frac{7+\cos (2 \pi / n)}{8} \\
& \geq \frac{6}{8}>\frac{1}{2} .
\end{align*}
$$

We calculate $\Re \omega$. Notice first that

$$
\begin{equation*}
\min _{g \in G \backslash\{0\}}(1-\Re \rho(g))=1-\cos (2 \pi / n)>0 . \tag{6.52}
\end{equation*}
$$

Thus,

$$
\begin{align*}
& \Re \omega=\left(\sum_{g \in G}(1-\Re(\rho(g))) \frac{\prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}{\prod_{k \in K} \phi_{\beta}\left(0+g_{k}\right)^{2}}\right) \cdot \frac{\prod_{k \in K} \phi_{\beta}\left(0+g_{k}\right)^{2}}{\sum_{g \in G} \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}} \\
& =\left(\sum_{g \in G \backslash\{0\}}(1-\Re(\rho(g))) \frac{\prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}{\prod_{k \in K} \phi_{\beta}\left(0+g_{k}\right)^{2}}\right) \cdot \frac{\prod_{k \in K} \phi_{\beta}\left(0+g_{k}\right)^{2}}{\sum_{g \in G} \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}} \\
& \stackrel{(6.51),(6.52)}{(1}-\cos (2 \pi / n)) \sum_{g \in G \backslash\{0\}} \frac{\prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}{\prod_{k \in K} \phi_{\beta}\left(0+g_{k}\right)^{2}} \cdot \frac{1}{2} \\
& \stackrel{\sqrt{6.50}}{\geq} \frac{1}{2}(1-\cos (2 \pi / n)) \frac{1-\cos (2 \pi / n)}{8}=\frac{(1-\cos (2 \pi / n))^{2}}{4^{2}} \text {. } \tag{6.53}
\end{align*}
$$

Since

$$
\begin{equation*}
\max _{g \in G}|1-\rho(g)| \leq 1+1=2 \tag{6.54}
\end{equation*}
$$

we have

$$
\begin{aligned}
|\omega| & =\left|\left(\sum_{g \in G \backslash\{0\}}(1-\rho(g)) \frac{\prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}{\prod_{k \in K} \phi_{\beta}\left(0+g_{k}\right)^{2}}\right) \cdot \frac{\prod_{k \in K} \phi_{\beta}\left(0+g_{k}\right)^{2}}{\sum_{g \in G} \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}\right| \\
& \left.\leq \max _{g \in G}|1-\rho(g)| \sum_{g \in G \backslash\{0\}} \prod_{k \in K} \frac{\phi_{\beta}\left(g+g_{k}\right)^{2}}{\phi_{\beta}\left(0+g_{k}\right)^{2}} \cdot \frac{\prod_{k \in K} \phi_{\beta}\left(0+g_{k}\right)^{2}}{\sum_{g \in G} \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}} \right\rvert\, \\
& \stackrel{(6.54\}}{\leq} 2 \sum_{g \in G \backslash\{0\}} \prod_{k \in K} \frac{\phi_{\beta}\left(g+g_{k}\right)^{2}}{\phi_{\beta}\left(0+g_{k}\right)^{2}} \cdot \frac{\prod_{k \in K} \phi_{\beta}\left(0+g_{k}\right)^{2}}{\sum_{g \in G} \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}} \\
& \stackrel{(6.51]}{\leq} 2 \sum_{g \in G \backslash\{0\}} \prod_{k \in K} \frac{\phi_{\beta}\left(g+g_{k}\right)^{2}}{\phi_{\beta}\left(0+g_{k}\right)^{2}} \cdot 1 \\
& \frac{[6.500}{\leq} \frac{1-\cos (2 \pi / n)}{4} .
\end{aligned}
$$

Therefore,

$$
|\omega|^{2} \leq\left(\frac{1-\cos (2 \pi / n)}{4}\right)^{2} \stackrel{\sqrt{6.53 \mid}}{\leq} \Re \omega .
$$

(ii) We calculate $\left|S_{\beta}\left(\left(g_{k}\right)_{k \in K}\right)\right|=|1-\omega|$ :

$$
\begin{aligned}
|1-\omega| & =\sqrt{(1-\Re \omega)^{2}+(\Im \omega)^{2}}=\sqrt{1-2 \Re \omega+|\omega|^{2}} \\
& \leq \sqrt{1-2 \Re \omega+|\omega|^{2}+\Re \omega\left(-|\omega|^{2}+\Re \omega\right)+\frac{|\omega|^{4}}{4}} \\
& =\sqrt{\left(1-\Re \omega+\frac{|\omega|^{2}}{2}\right)^{2}}=1-\Re \omega+\frac{|\omega|^{2}}{2} \\
& \stackrel{\sqrt{6.47 \mid}}{\leq} 1-\Re \omega+\frac{\Re \omega}{2}=1-\frac{\Re \omega}{2} \\
& \stackrel{\sqrt{6.53)}}{\leq} 1-\frac{1}{2}\left(\frac{(1-\cos (2 \pi / n))}{2} \sum_{g \in G \backslash\{0\}} \frac{\prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}{\prod_{k \in K} \phi_{\beta}\left(0+g_{k}\right)^{2}}\right) \\
& =1-C_{*} \sum_{g \in G \backslash\{0\}} \frac{\prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}{\prod_{k \in K} \phi_{\beta}\left(0+g_{k}\right)^{2}} .
\end{aligned}
$$

Before the following lemma, the notation $\angle\left(z_{1}, z_{2}\right)$ is introduced. Let $z_{1}, z_{2} \in$ $\mathbb{C} \backslash\{0\}$, then $\angle\left(z_{1}, z_{2}\right)$ is the absolute value of the smallest angle between these numbers.

Lemma 6.14. Let $K$ be a finite and non-empty index set and assume that $\beta_{0}=$ $\beta_{0}(|K|)>0$ satisfies $(\star)$ applied with $K$. For each $k \in K$, let $g_{k} \in G$. If
$G_{0}\left[\left(g_{k}\right)_{k \in K}\right]=\{0\}$, then

$$
\begin{equation*}
\max _{g \in G \backslash\{0\}} \prod_{k \in K} \frac{\phi_{\beta}\left(g+g_{k}\right)}{\phi_{\beta}\left(0+g_{k}\right)} \geq \max _{g \in G \backslash\{0\}} \prod_{k \in K} \frac{\phi_{\beta}(g+0)}{\phi_{\beta}(0+0)} . \tag{6.55}
\end{equation*}
$$

Proof. Since

$$
\prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)=\exp \left(\beta \Re\left(\rho(g) \sum_{k \in K} \rho\left(g_{k}\right)\right)\right),
$$

inequality (6.55) can be written as

$$
\max _{g \in G \backslash\{0\}} \frac{\exp \left(\beta \Re\left(\rho(g) \sum_{k \in K} \rho\left(g_{k}\right)\right)\right)}{\exp \left(\beta \Re\left(\rho(0) \sum_{k \in K} \rho\left(g_{k}\right)\right)\right)} \geq \max _{g \in G \backslash\{0\}} \frac{\exp \left(\beta \Re\left(\rho(g) \sum_{k \in K} \rho(0)\right)\right)}{\exp \left(\beta \Re\left(\rho(0) \sum_{k \in K} \rho(0)\right)\right)},
$$

which is equivalent to

$$
\begin{aligned}
& \max _{g \in G \backslash\{0\}} \exp \left(\beta\left(\Re\left(\rho(g) \sum_{k \in K} \rho\left(g_{k}\right)\right)-\Re\left(\rho(0) \sum_{k \in K} \rho\left(g_{k}\right)\right)\right)\right) \\
& \geq \max _{g \in G \backslash\{0\}} \exp \left(\beta\left(\Re\left(\rho(g) \sum_{k \in K} \rho(0)\right)-\Re\left(\rho(0) \sum_{k \in K} \rho(0)\right)\right)\right)
\end{aligned}
$$

and

$$
\begin{align*}
& \max _{g \in G \backslash\{0\}}\left(\Re\left(\rho(g) \sum_{k \in K} \rho\left(g_{k}\right)\right)-\Re\left(\rho(0) \sum_{k \in K} \rho\left(g_{k}\right)\right)\right) \\
& \geq \max _{g \in G \backslash\{0\}}\left(\Re\left(\rho(g) \sum_{k \in K} \rho(0)\right)-\Re\left(\rho(0) \sum_{k \in K} \rho(0)\right)\right) . \tag{6.56}
\end{align*}
$$

Since $\rho(g)$ and $\sum_{k \in K} \rho\left(-g_{k}\right)$ are non-zero complex numbers, we have

$$
\begin{align*}
\Re\left(\rho(g) \sum_{k \in K} \rho\left(g_{k}\right)\right) & =\Re\left(\sum_{k \in K} \overline{\rho\left(-g_{k}\right)} \rho(g)\right) \\
& =\left|\sum_{k \in K} \rho\left(g_{k}\right)\right||\rho(g)| \cos \left(\angle\left(\sum_{k \in K} \rho\left(-g_{k}\right), \rho(g)\right)\right)  \tag{6.57}\\
& =\left|\sum_{k \in K} \rho\left(g_{k}\right)\right| \cos \left(\angle\left(\sum_{k \in K} \rho\left(-g_{k}\right), \rho(g)\right)\right) .
\end{align*}
$$

Fix $\hat{g} \in G \backslash\{0\}$. Then $\hat{g} \notin G_{0}\left[\left(g_{k}\right)_{k \in K}\right]$ and

$$
\prod_{k \in K} \phi_{\beta}\left(\hat{g}+g_{k}\right)<\prod_{k \in K} \phi_{\beta}\left(0+g_{k}\right),
$$

which is equivalent to

$$
\begin{equation*}
\Re\left(\rho(g) \sum_{k \in K} \rho\left(g_{k}\right)\right)<\Re\left(\rho(0) \sum_{k \in K} \rho\left(g_{k}\right)\right)=\Re\left(1 \cdot \sum_{k \in K} \rho\left(g_{k}\right)\right) . \tag{6.58}
\end{equation*}
$$

From equation (6.57) follows that inequality (6.58) is equivalent to

$$
\begin{aligned}
& \left|\sum_{k \in K} \rho\left(g_{k}\right)\right| \cos \left(\angle\left(\sum_{k \in K} \rho\left(-g_{k}\right), \rho(g)\right)\right) \\
& <\left|\sum_{k \in K} \rho\left(g_{k}\right)\right| \cos \left(\angle\left(\sum_{k \in K} \rho\left(-g_{k}\right), 1\right)\right) .
\end{aligned}
$$

Thus,

$$
\begin{equation*}
\cos \left(\angle\left(\sum_{k \in K} \rho\left(-g_{k}\right), \rho(g)\right)\right)-\cos \left(\angle\left(\sum_{k \in K} \rho\left(-g_{k}\right), 1\right)\right)<0 \tag{6.59}
\end{equation*}
$$

We write $G_{0}\left[\left(g_{k}\right)_{k \in K}\right]$ with $\arg \max$ and $\arg \min$. Note that $\cos (x)$ is decreasing in $[0, \pi]$ and therefore takes it largest value when the angle is as small as possible. Therefore,

$$
\begin{aligned}
&\{0\}=G_{0}\left[\left(g_{k}\right)_{k \in K}\right] \\
&=\underset{g \in G}{\arg \max } \exp \left(\beta \Re\left(\rho(g) \sum_{k \in K}\left(g+g_{k}\right)\right)\right) \\
& \stackrel{\boxed{6.577}}{=} \underset{g \in G}{\arg \max } \exp \left(\beta\left|\sum_{k \in K} \rho\left(g_{k}\right)\right| \cos \left(\angle\left(\sum_{k \in K} \rho\left(-g_{k}\right), \rho(g)\right)\right)\right) \\
&=\underset{g \in G}{\arg \max } \cos \left(\angle\left(\sum_{k \in K} \rho\left(-g_{k}\right), \rho(g)\right)\right) \\
&=\underset{g \in G}{\arg \min } \angle\left(\sum_{k \in K} \rho\left(-g_{k}\right), \rho(g)\right) .
\end{aligned}
$$

Thus, $\rho(0)$ is the point in $\rho(G)$ that is closest to $\sum_{k \in K} \rho\left(-g_{k}\right)$. Similarly, we have for $\hat{g}$ that

$$
\hat{g} \in \underset{g \in G \backslash G_{0}\left[\left(g_{k}\right)_{k \in K]}\right]}{\arg \max } \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)=\underset{g \in G \backslash\{0\}}{\arg \min } \angle\left(\sum_{k \in K} \rho\left(-g_{k}\right), \rho(g)\right) .
$$

Hence, $\rho(\hat{g})$ is the point second closest to $\sum_{k \in K} \rho\left(-g_{k}\right)$. These results for 0 and $\hat{g}$ are illustrated in Figure 6.1. Note that there are two possibilities for where $\sum_{k \in K} \rho\left(-g_{k}\right)$ is.


Figure 6.1: The points in $\rho(G)$ and the two possibilities for $\sum_{k \in K} \rho\left(-g_{k}\right)$. The figure is from [1].

We see that the angle $\angle(1, \rho(\hat{g}))$ can be written as a sum of two angles:

$$
\angle\left(\sum_{k \in K} \rho\left(-g_{k}\right), 1\right)+\angle\left(\sum_{k \in K} \rho\left(-g_{k}\right), \rho(\hat{g})\right)=\angle(1, \rho(\hat{g})) .
$$

Since $\cos (x)$ is decreasing on $[0, \pi]$, we have

$$
\cos \left(\angle\left(\sum_{k \in K} \rho\left(-g_{k}\right), \rho(\hat{g})\right)\right) \geq \cos (\angle(1, \rho(\hat{g}))
$$

Furthermore,

$$
1=\cos (\angle(1,1)) \geq \cos \left(\angle\left(\sum_{k \in K} \rho\left(-g_{k}\right), 1\right)\right) .
$$

Therefore, we have

$$
\begin{aligned}
& \cos (\angle(1,1))+\cos \left(\angle\left(\sum_{k \in K} \rho\left(-g_{k}\right), \rho(\hat{g})\right)\right) \\
& \geq \cos \left(\angle\left(\sum_{k \in K} \rho\left(-g_{k}\right), 1\right)\right)+\cos (\angle(1, \rho(\hat{g}))
\end{aligned}
$$

and

$$
\begin{align*}
& \cos \left(\angle\left(\sum_{k \in K} \rho\left(-g_{k}\right), \rho(\hat{g})\right)\right)-\cos \left(\angle\left(\sum_{k \in K} \rho\left(-g_{k}\right), 1\right)\right)  \tag{6.60}\\
& \geq \cos (\angle(1, \rho(\hat{g})))-\cos (\angle(1,1))
\end{align*}
$$

Since

$$
\begin{equation*}
\left|\sum_{k \in K} \rho\left(g_{k}\right)\right| \leq \sum_{k \in K}\left|\rho\left(g_{k}\right)\right| \leq \sum_{k \in K} 1=|K| \tag{6.61}
\end{equation*}
$$

we have

$$
\begin{aligned}
& \left.\Re\left(\rho(\hat{g}) \sum_{k \in K} \rho\left(g_{k}\right)\right)-\Re(\rho(0)) \sum_{k \in K} \rho\left(g_{k}\right)\right) \\
& \stackrel{\sqrt{6.57}}{=}\left|\sum_{k \in K} \rho\left(g_{k}\right)\right| \cos \left(\angle\left(\sum_{k \in K} \rho\left(-g_{k}\right), \rho(\hat{g})\right)\right) \\
& -\left|\sum_{k \in K} \rho\left(g_{k}\right)\right| \cos \left(\angle\left(\sum_{k \in K} \rho\left(-g_{k}\right), \rho(1)\right)\right) \\
& \stackrel{\sqrt{6.59}, \sqrt[\text { (6.61 }]{ }|K|}{|K|}\left(\cos \left(\angle\left(\sum_{k \in K} \rho\left(-g_{k}\right), \rho(\hat{g})\right)\right)-\cos \left(\angle\left(\sum_{k \in K} \rho\left(-g_{k}\right), \rho(1)\right)\right)\right) \\
& \stackrel{\sqrt{66.60}}{\geq}|K| \cos (\angle(1, \rho(\hat{g}))-\cos (\angle(1,1))) \\
& =\left|\sum_{k \in K} \rho(0)\right|\left(\cos \left(\angle\left(\sum_{k \in K} \rho(0), \rho(\hat{g})\right)\right)-\cos \left(\angle\left(\sum_{k \in K} \rho(0), \rho(0)\right)\right)\right) \\
& \left.\stackrel{\sqrt{6.57}}{=} \Re\left(\rho(\hat{g}) \sum_{k \in K} \rho(0)\right)-\Re(\rho(0)) \sum_{k \in K} \rho(0)\right) \text {. }
\end{aligned}
$$

Since $\hat{g} \in G \backslash\{0\}$, it follows that

$$
\begin{aligned}
& \max _{g \in G \backslash\{0\}}\left(\Re\left(\rho(g) \sum_{k \in K} \rho\left(g_{k}\right)\right)-\Re(\rho(0)) \sum_{k \in K} \rho\left(g_{k}\right)\right) \\
& \geq \max _{g \in G \backslash\{0\}}\left(\Re\left(\rho(g) \sum_{k \in K} \rho(0)\right)-\Re(\rho(0)) \sum_{k \in K} \rho(0)\right) .
\end{aligned}
$$

This is the same as inequality (6.56), which was shown to be equivalent to inequality (6.55).

We continue to the proof of Lemma 6.12.
Proof of Lemma 6.12. We investigate how many elements $G_{0}\left[\left(g_{k}\right)_{k \in K}\right]$ can consist of:

$$
\begin{aligned}
\prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2} & =\exp \left(\beta \sum_{k \in K} \Re\left(\rho\left(g+g_{k}\right)\right)\right)=\exp \left(\beta \sum_{k \in K} \Re\left(\rho(g) \rho\left(g_{k}\right)\right)\right) \\
& =\exp \left(\beta \Re\left(\rho(g) \sum_{k \in K} \rho\left(g_{k}\right)\right)\right) .
\end{aligned}
$$

Recall that $G_{0}\left[\left(g_{k}\right)_{k \in K}\right]$ is the set of the elements $g \in G$ for which the above function is maximised.

First, if $\sum_{k \in K} \rho\left(g_{k}\right)=0$, then $\prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}=1$ for all $g \in G$. Thus, $G_{0}\left[\left(g_{k}\right)_{k \in K}\right]$ consists of $G$ elements.

Second, we consider the case when $\sum_{k \in K} \rho\left(g_{k}\right) \in \mathbb{R} \backslash\{0\}$. If $\sum_{k \in K} \rho\left(g_{k}\right)>0$, then the function is maximised when $\Re \rho(g)$ is maximised, i.e. $g=0$. Otherwise, it is maximised when $\Re \rho(g)$ is minimised. This happens when $\rho(g)=e^{g \cdot 2 \pi i m / n}$ is as close to $e^{\pi i}$ as possible. Thus, there is 1 or 2 elements in $G_{0}\left[\left(g_{k}\right)_{k \in K}\right]$.

Third, if $\sum_{k \in K} \rho\left(g_{k}\right) \notin \mathbb{R}$, then $\sum_{k \in K} \rho\left(g_{k}\right)=r e^{i \theta}$. Hence,

$$
\Re\left(\rho(g) \sum_{k \in K} \rho\left(g_{k}\right)\right)=r \Re\left(\rho(g) e^{i \theta}\right) .
$$

This is maximised when $\rho(g)=e^{g \cdot 2 \pi i m / n}$ is as close as possible to $e^{-i \theta}$. Thus, we have that there are 1 or 2 elements in $G_{0}\left[\left(g_{k}\right)_{k \in K}\right]$.

To conclude, there are three possibilities for the numbers of elements in $G_{0}\left[\left(g_{k}\right)_{k \in K}\right]: 1,2$ or $G$ elements. We continue by proving that inequality (6.44) holds in all three cases.

First case: G elements
If $\left|G_{0}\left[\left(g_{k}\right)_{k \in K}\right]\right|=G$, then $\prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}=1$ for all $g \in G$. Hence

$$
\begin{aligned}
\left|\frac{\sum_{g \in G} \rho(g) \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}{\sum_{g \in G} \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}\right| & =\frac{\left|\sum_{g \in G} \rho(g)\right|}{|G|}=\frac{\left|\sum_{k=0}^{n-1} e^{k \cdot 2 \pi i m / n}\right|}{|G|} \\
& =\frac{1}{|G|} \frac{e^{2 \pi i m}-1}{e^{2 \pi i m / n}-1}=0 \leq 1-C_{*} \lambda(\beta)^{2|K|} .
\end{aligned}
$$

## Second case: two elements

If $\left|G_{0}\left[\left(g_{k}\right)_{k \in K}\right]\right|=2$, then

$$
\begin{aligned}
\left|S_{\beta}\left(\left(g_{k}\right)_{k \in K}\right)\right|= & \left|\frac{\sum_{g \in G} \rho(g) \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}{\sum_{g \in G} \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}\right| \\
\leq & \left|\frac{\sum_{g \in G} \rho(g) \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}{\sum_{g \in G_{0}\left[\left(g_{k}\right)_{k \in K}\right]} \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}\right| \\
\leq & \left|\frac{\sum_{g \in G_{0}\left[\left(g_{k}\right)_{k \in K}\right]} \rho(g) \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}{\sum_{g \in G_{0}\left[\left(g_{k}\right)_{k \in K}\right]} \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}\right| \\
& +\left|\frac{\sum_{g \in G \backslash G_{0}\left[\left(g_{k}\right)_{k \in K}\right]} \rho(g) \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}{\sum_{g \in G_{0}\left[\left(g_{k}\right)_{k \in K}\right]} \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}\right| .
\end{aligned}
$$

Notice that for every $g \in G_{0}\left[\left(g_{k}\right)_{k \in K}\right]$, the function $\phi_{\beta}\left(g+g_{k}\right)$ takes the same
value. Let $g^{\prime} \in G_{0}\left[\left(g_{k}\right)_{k \in K}\right]$. Then,

$$
\begin{align*}
\left|S_{\beta}\left(\left(g_{k}\right)_{k \in K}\right)\right| \leq & \left|\frac{\sum_{g \in G_{0}\left[\left(g_{k}\right)_{k \in K}\right]} \rho(g)}{G_{0}\left[\left(g_{k}\right)_{k \in K}\right]}\right|+\left|\frac{\sum_{g \in G \backslash G_{0}\left[\left(g_{k}\right)_{k \in K}\right]} \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}{\sum_{g \in G_{0}\left[\left(g_{k}\right)_{k \in K}\right]} \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}\right| \\
& \cdot|\rho(g)| \\
& \stackrel{2.33}{=} \frac{\left|\sum_{g \in G_{0}\left[\left(g_{k}\right)_{k \in K}\right]} \rho(g)\right|}{2}+\left|\frac{\sum_{g \in G \backslash G_{0}\left[\left(g_{k}\right)_{k \in K}\right]} \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}{\sum_{g \in G_{0}\left[\left(g_{k}\right)_{k \in K}\right]} \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}\right| \\
\leq & \frac{\left|\sum_{g \in G_{0}\left[\left(g_{k}\right)_{k \in K}\right]} \rho(g)\right|}{2}+\left|\sum_{g \in G \backslash G_{0}\left[\left(g_{k}\right)_{k \in K}\right]} \prod_{k \in K} \frac{\phi_{\beta}\left(g+g_{k}\right)^{2}}{\phi_{\beta}\left(g^{\prime}+g_{k}\right)^{2}}\right| \\
& \stackrel{(\star)}{\leq} \frac{\left|\sum_{g \in G_{0}\left[\left(g_{k}\right)_{k \in K]} \rho\right.} \rho(g)\right|}{2}+\frac{1-\cos (2 \pi / n)}{8} . \tag{6.62}
\end{align*}
$$

Since $G_{0}\left[\left(g_{k}\right)_{k \in K}\right]=\{j, j+1\}$ for a $j \in \mathbb{Z}_{n}$, we have

$$
\begin{aligned}
\frac{\left|\sum_{g \in G_{0}\left[\left(g_{k}\right)_{k \in K}\right]}^{2} \rho(g)\right|}{2} & =\frac{|\rho(j)+\rho(j+1)|}{2} \\
& \leq \frac{\sqrt{(1+\cos (2 \pi / n))^{2}+\sin ^{2}(2 \pi / n)}}{2} \\
& =\frac{\sqrt{1+2 \cos (2 \pi / n)+\cos ^{2}(2 \pi / n)+\sin ^{2}(2 \pi / n)}}{2} \\
& =\sqrt{\frac{1+\cos (2 \pi / n)}{2}} \\
& =\sqrt{1-\frac{1-\cos (2 \pi / n)}{2}} \\
& \leq \sqrt{1-\frac{1-\cos (2 \pi / n)}{4}} \\
& \leq 1-\frac{1-\cos (2 \pi / n)}{4}
\end{aligned}
$$

Combining this with inequality (6.62), we have

$$
\begin{aligned}
\left|S_{\beta}\left(\left(g_{k}\right)_{k \in K}\right)\right| & \leq 1-\frac{1-\cos (2 \pi / n)}{4}+\frac{1-\cos (2 \pi / n)}{8} \\
& =1-\frac{1-\cos (2 \pi / n)}{8}=1-\frac{1}{2} C_{*} \\
& \stackrel{\sqrt{6.43}}{\leq} 1-C_{*} \lambda\left(\beta_{0}\right)^{2|K|} .
\end{aligned}
$$

## Third case: one element

Last, we focus on the case when $\left|G_{0}\left[\left(g_{k}\right)_{k \in K}\right]\right|=1$. First, we investigate which
element is in $G_{0}\left[\left(g_{k}\right)_{k \in K}\right]$. For any $g^{\prime} \in G$, we have

$$
\begin{aligned}
\left|S_{\beta}\left(\left(g^{\prime}+g\right)_{k \in K}\right)\right| & =\left|\frac{\sum_{g \in G} \rho(g) \prod_{k \in K} \phi_{\beta}\left(g+\left(g^{\prime}+g_{k}\right)\right)^{2}}{\sum_{g \in G} \prod_{k \in K} \phi_{\beta}\left(g+\left(g^{\prime}+g_{k}\right)\right)^{2}}\right| \\
& =\left|\rho\left(g^{\prime}\right)\right|\left|\frac{\sum_{g \in G} \rho(g) \prod_{k \in K} \phi_{\beta}\left(g+\left(g^{\prime}+g_{k}\right)\right)^{2}}{\sum_{g \in G} \prod_{k \in K} \phi_{\beta}\left(g+\left(g^{\prime}+g_{k}\right)\right)^{2}}\right| \\
& =\left|\frac{\sum_{g \in G} \rho(g) \rho\left(g^{\prime}\right) \prod_{k \in K} \phi_{\beta}\left(g+\left(g^{\prime}+g_{k}\right)\right)^{2}}{\sum_{g \in G} \prod_{k \in K} \phi_{\beta}\left(g+\left(g^{\prime}+g_{k}\right)\right)^{2}}\right| \\
& =\left|\frac{\sum_{g \in G} \rho\left(g+g^{\prime}\right) \prod_{k \in K} \phi_{\beta}\left(\left(g+g^{\prime}\right)+g_{k}\right)^{2}}{\sum_{g \in G} \prod_{k \in K} \phi_{\beta}\left(\left(g+g^{\prime}\right)+g_{k}\right)^{2}}\right| \\
& =\left|\frac{\sum_{g \in G} \rho(g) \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}{\sum_{g \in G} \prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}\right| \\
& =\left|S_{\beta}\left(\left(g_{k}\right)_{k \in K}\right)\right| .
\end{aligned}
$$

Thus, it implies that $0 \in G_{0}\left[\left(g_{k}\right)_{k \in K}\right]$, which only consists of one element. Therefore, $G_{0}\left[\left(g_{k}\right)_{k \in K}\right]=\{0\}$. Then, from Lemma 6.13 follows that

$$
\left|S_{\beta}\left(\left(g_{k}\right)_{k \in K}\right)\right| \leq 1-C_{*} \sum_{g \in G \backslash\{0\}} \frac{\prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}{\prod_{k \in K} \phi_{\beta}\left(0+g_{k}\right)^{2}}
$$

Therefore, inequality (6.44) holds if

$$
\begin{aligned}
\sum_{g \in G \backslash\{0\}} \frac{\prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}{\prod_{k \in K} \phi_{\beta}\left(0+g_{k}\right)^{2}} & \geq \lambda(\beta)^{2|K|}=\left(\max _{g \in G \backslash\{0\}} \frac{\phi_{\beta}(g)}{\phi_{\beta}(0)}\right)^{2|K|} \\
& =\max _{g \in G \backslash\{0\}} \prod_{k \in K} \frac{\phi_{\beta}(g+0)^{2}}{\phi_{\beta}(0+0)^{2}},
\end{aligned}
$$

which holds if

$$
\begin{equation*}
\max _{g \in G \backslash\{0\}} \prod_{k \in K} \frac{\phi_{\beta}\left(g+g_{k}\right)}{\phi_{\beta}\left(0+g_{k}\right)} \geq \max _{g \in G \backslash\{0\}} \prod_{k \in K} \frac{\phi_{\beta}(g+0)}{\phi_{\beta}(0+0)} . \tag{6.63}
\end{equation*}
$$

Inequality (6.63) was proved to be true in Lemma 6.14. Thus,

$$
\begin{aligned}
\left|S_{\beta}\left(\left(g_{k}\right)_{k \in K}\right)\right| & \leq 1-C_{*} \sum_{g \in G \backslash\{0\}} \frac{\prod_{k \in K} \phi_{\beta}\left(g+g_{k}\right)^{2}}{\prod_{k \in K} \phi_{\beta}\left(0+g_{k}\right)^{2}} \\
& \leq 1-C_{*} \sum_{g \in G \backslash\{0\}} \frac{\prod_{k \in K} \phi_{\beta}(g+0)^{2}}{\prod_{k \in K} \phi_{\beta}(0+0)^{2}} \\
& =1-C_{*} \lambda(\beta)^{2|K|} .
\end{aligned}
$$

We continue to the proof of Proposition 6.11.
Proof of Proposition 6.11. Recall, by Lemma 6.9 follows that the spins $\left(\sigma_{e}\right)_{e \in \gamma_{1}}$ are independent if the spins $\left(\sigma_{e}\right)_{e \notin \pm \gamma_{1}}$ are given. Take an edge $e \in \gamma_{1}$. For $p \in \hat{\partial} e$, let $\sigma_{p}^{e}:=\sum_{e^{\prime} \in \partial p \backslash\{e\}} \sigma_{e^{\prime}}$. For $g \in G$, the conditional probability when $\sigma_{e}=g$ is

$$
\mu_{\beta, N}^{\prime}\left(\sigma_{e}=g\right)=\frac{\prod_{p \in \hat{\partial} e} \phi_{\beta}\left(\sigma_{p}^{e}+g\right)^{2}}{\sum_{g^{\prime} \in G} \prod_{p \in \hat{\partial} e} \phi_{\beta}\left(\sigma_{p}^{e}+g^{\prime}\right)^{2}}
$$

and the conditional expectation of $\rho\left(\sigma_{e}\right)$ is

$$
\mathbb{E}_{\beta, N}^{\prime}\left[\rho\left(\sigma_{e}\right)\right]=\frac{\sum_{g \in G} \rho(g) \prod_{p \in \hat{\jmath} e} \phi_{\beta}\left(\sigma_{p}^{e}+g\right)^{2}}{\sum_{g \in G} \prod_{p \in \hat{\partial} e} \phi_{\beta}\left(\sigma_{p}^{e}+g\right)^{2}}
$$

Since the spins $\left(\sigma_{e}\right)_{e \in \gamma_{1}}$ are independent, we obtain

$$
\begin{equation*}
\mathbb{E}_{\beta, N}^{\prime}\left[\prod_{e \in \gamma_{1}} \rho\left(\sigma_{e}\right)\right]=\prod_{e \in \gamma_{1}} \mathbb{E}_{\beta, N}^{\prime}\left[\rho\left(\sigma_{e}\right)\right] . \tag{6.64}
\end{equation*}
$$

Before calculating the expectation of $W_{\gamma}$, we apply Lemma 6.12 for $K=\hat{\partial} e$ and $g_{k}=\sigma_{p}^{e}$. Since $|K|=6$, we obtain

\[

\]

Thus, the upper bound for the expectation is

$$
\begin{aligned}
\left|\mathbb{E}_{\beta, N}\left[W_{\gamma}\right]\right| & =\left|\mathbb{E}_{\beta, N}\left[\rho\left(\sum_{e \in \gamma} \sigma_{e}\right)\right]\right|=\left|\mathbb{E}_{\beta, N}\left[\rho\left(\sum_{e \in \gamma \backslash \gamma_{1}} \sigma_{e}\right) \rho\left(\sum_{e \in \gamma_{1}} \sigma_{e}\right)\right]\right| \\
& =\left|\mathbb{E}_{\beta, N}\left[\prod_{e \in \gamma \backslash \gamma_{1}} \rho\left(\sigma_{e}\right) \prod_{e \in \gamma_{1}} \rho\left(\sigma_{e}\right)\right]\right| \\
& =\left|\mathbb{E}_{\beta, N}\left[\mathbb{E}_{\beta, N}^{\prime}\left[\prod_{e \in \gamma \backslash \gamma_{1}} \rho\left(\sigma_{e}\right) \prod_{e \in \gamma_{1}} \rho\left(\sigma_{e}\right)\right]\right]\right| \\
& =\left|\mathbb{E}_{\beta, N}\left[\prod_{e \in \gamma \gamma_{1}} \rho\left(\sigma_{e}\right) \mathbb{E}_{\beta, N}^{\prime}\left[\prod_{e \in \gamma_{1}} \rho\left(\sigma_{e}\right)\right]\right]\right| \\
& \leq \mathbb{E}_{\beta, N}\left[\left|\prod_{e \in \gamma \backslash \gamma_{1}} \rho\left(\sigma_{e}\right) \mathbb{E}_{\beta, N}^{\prime}\left[\prod_{e \in \gamma_{1}} \rho\left(\sigma_{e}\right)\right]\right|\right]
\end{aligned}
$$

$$
\begin{aligned}
& =\mathbb{E}_{\beta, N}\left[\left|\prod_{e \in \gamma \backslash \gamma_{1}} \rho\left(\sigma_{e}\right)\right| \cdot\left|\mathbb{E}_{\beta, N}^{\prime}\left[\prod_{e \in \gamma_{1}} \rho\left(\sigma_{e}\right)\right]\right|\right] \\
& =\mathbb{E}_{\beta, N}\left[1 \cdot\left|\mathbb{E}_{\beta, N}^{\prime}\left[\prod_{e \in \gamma_{1}} \rho\left(\sigma_{e}\right)\right]\right|\right] \\
& \stackrel{(6.64}{=} \mathbb{E}_{\beta, N}\left[\prod_{e \in \gamma_{1}}\left|\mathbb{E}_{\beta, N}^{\prime}\left[\rho\left(\sigma_{e}\right)\right]\right|\right] \\
& \stackrel{(6.65)}{\leq} \mathbb{E}_{\beta, N}\left[\prod_{e \in \gamma_{1}} e^{-C_{*} \lambda(\beta)^{12}}\right] \\
& =\mathbb{E}_{\beta, N}\left[e^{-C_{*} \lambda(\beta)^{12}\left|\gamma_{1}\right|}\right]=e^{-C_{*} \lambda(\beta)^{12}\left|\gamma_{1}\right|} \\
& =e^{-C_{*}\left(\ell-\ell_{c}\right) \lambda(\beta)^{12}} .
\end{aligned}
$$

### 6.3 The last part of the proof

Recall that the inequality we want to prove is inequality (2.12):

$$
\left|\left\langle W_{\gamma}\right\rangle_{\beta}-e^{-\ell(1-\theta(\beta))}\right| \leq C^{\prime}\left[\frac{\sqrt{\ell_{c}}}{\ell}+\lambda(\beta)^{2}\right]^{C^{\prime \prime}}
$$

We begin with defining the constants $C^{\prime}$ and $C^{\prime \prime}$. Let

$$
\begin{equation*}
C^{\prime}:=\sqrt{2}\left(C_{A} 2^{4 C^{*} / C_{*}}\right)^{1 /\left(1+4 C^{*} / C_{*}\right)} \text { and } C^{\prime \prime}:=1 /\left(1+4 C^{*} / C_{*}\right) . \tag{6.66}
\end{equation*}
$$

The proof is divided in two cases, one when $\ell_{c} \geq \frac{\ell}{2}$ and one when $\ell_{c}<\frac{\ell}{2}$. The inequality

$$
\begin{equation*}
\left|\mathbb{E}_{\beta, N}\left[W_{\gamma}\right]-e^{-\ell(1-\theta(\beta))}\right| \leq C^{\prime}\left[\frac{\sqrt{\ell_{c}}}{\ell}+\lambda(\beta)^{2}\right]^{C^{\prime \prime}} \tag{6.67}
\end{equation*}
$$

will be proved, which then implies inequality 2.12. Since we earlier proved that the limit $\left\langle W_{\gamma}\right\rangle$ exists and both the right-hand side of the inequality and $e^{-\ell(1-\theta(\beta))}$ are independent of $N$, taking the limits of inequality (6.67) when $N$ goes to infinity proves inequality (2.12).
6.3.1 First case: $\ell_{c} \geq \frac{\ell}{2}$

Assume that $\ell_{c} \geq \frac{\ell}{2}$. First, a lower bound for the right-hand side of 667) is calculated. From the definitions of $C^{*}$ and $C_{*}$ (equations 6.42 and (6.1))
follows that they are positive. Therefore, $4 C^{*} / C_{*}>0$ and $\left.C^{\prime \prime} \in\right] 0,1[$. From the definition of $C_{A}$ (equation 6.41) follows that $C_{A} \geq \frac{9}{2}$. Thus,

$$
\begin{aligned}
C^{\prime} & =\sqrt{2}\left(C_{A} 2^{4 C^{*} / C_{*}}\right)^{1 /\left(1+4 C^{*} / C_{*}\right)} \\
& =\sqrt{2}\left(C_{A}\right)^{1 /\left(1+4 C^{*} / C_{*}\right)} \cdot\left(2^{4 C^{*} / C_{*}}\right)^{1 /\left(1+4 C^{*} / C_{*}\right)} \\
& \geq \sqrt{2} \cdot 1 \cdot 2=2 \sqrt{2} .
\end{aligned}
$$

From the assumption that $\ell_{c} \geq \frac{\ell}{2}$ follows that

$$
\sqrt{\frac{l_{c}}{l}} \geq \sqrt{\frac{1}{2}}
$$

Therefore, the lower bound of the right-hand side of inequality (6.67) is

$$
C^{\prime}\left[\sqrt{\frac{l_{c}}{l}}+\lambda(\beta)^{2}\right]^{C^{\prime \prime}} \geq 2 \sqrt{2}\left(\sqrt{\frac{1}{2}}+0\right)^{C^{\prime \prime}}=2^{\frac{3}{2}-\frac{C^{\prime \prime}}{2}} \geq 2
$$

For the left-hand side of inequality (6.67), we obtain

$$
\begin{aligned}
\left|\mathbb{E}_{\beta, N}\left[W_{\gamma}\right]-e^{-\ell(1-\theta(\beta))}\right| & \leq\left|\mathbb{E}_{\beta, N}\left[W_{\gamma}\right]\right|+\left|e^{-\ell(1-\theta(\beta))}\right| \\
& =\left|\mathbb{E}_{\beta, N}\left[\rho\left(\sum_{e \in \gamma} \sigma_{e}\right)\right]\right|+\left|e^{-\ell(1-\theta(\beta))}\right| \\
& \leq \mathbb{E}_{\beta, N}\left[\left|\rho\left(\sum_{e \in \gamma} \sigma_{e}\right)\right|\right]+1 \\
& \leq \mathbb{E}_{\beta, N}[1]+1=1+1=2 .
\end{aligned}
$$

Hence, we obtain

$$
\left|\mathbb{E}_{\beta, N}\left[W_{\gamma}\right]-e^{-\ell(1-\theta(\beta))}\right| \leq 2 \leq C^{\prime}\left(\frac{\sqrt{\ell_{c}}}{\ell}+\lambda(\beta)^{2}\right)^{C^{\prime \prime}}
$$

6.3.2 Second case: $\ell_{c}<\frac{\ell}{2}$

Assume that $\ell_{c}<\frac{\ell}{2}$. Then $\ell-\ell_{c} \geq \frac{1}{2} \ell$. From this assumption and Proposition 6.11 follows that

$$
\left|\mathbb{E}_{\beta, N}\left[W_{\gamma}\right]\right| \leq e^{-C_{*}\left(\ell-\ell_{c}\right) \lambda(\beta)^{12}} \leq e^{-\frac{1}{2} C_{*} \ell \lambda(\beta)^{12}}
$$

Therefore,

$$
\begin{align*}
\left|\mathbb{E}_{\beta, N}\left[W_{\gamma}\right]-e^{-\ell(1-\theta(\beta))}\right| & \leq\left|\mathbb{E}_{\beta, N}\left[W_{\gamma}\right]\right|+\left|e^{-\ell(1-\theta(\beta))}\right| \\
& \leq e^{-\frac{1}{2} C_{*} \ell \lambda(\beta)^{12}}+e^{-\ell(1-\theta(\beta))} . \tag{6.68}
\end{align*}
$$

Since $5(|G|-1) \lambda(\beta)^{2}<1$ and $0<\lambda(\beta)<1$, we have

$$
\begin{align*}
& 1-5(|G|-1) \lambda(\beta)^{2}>0 \\
\Leftrightarrow & \frac{\lambda(\beta)^{10}}{5}-(|G|-1) \lambda(\beta)^{12} \geq 0 \\
\Rightarrow & \frac{1}{5}-(|G|-1) \lambda(\beta)^{12} \geq 0 \\
\Leftrightarrow & 1-(|G|-1) \lambda(\beta)^{12} \geq \frac{4}{5} \geq \frac{1}{2} \tag{6.69}
\end{align*}
$$

The lower bound for $1-\theta(\beta)$ is then

$$
\begin{aligned}
1-\theta(\beta) & \stackrel{\sqrt{6.2)}}{=} 1-\frac{\sum_{g \in G} \Re(\rho(g)) e^{12 \beta \Re \rho(g)}}{\sum_{g \in G} e^{12 \beta \Re \rho(g)}} \\
& =\frac{\sum_{g \in G}(1-\Re(\rho(g))) e^{12 \beta \Re \rho(g)}}{\sum_{g \in G} e^{12 \beta \Re \rho(g)}} \\
& =\frac{\sum_{g \in G}(1-\Re(\rho(g))) e^{12 \beta(\Re \rho(g)-1)}}{\sum_{g \in G} e^{12 \beta(\Re \rho(g)-1)}} \\
& =\frac{\sum_{g \in G \backslash\{0\}}(1-\Re(\rho(g))) e^{12 \beta(\Re \rho(g)-1)}}{1+\sum_{g \in G \backslash\{0\}} e^{12 \beta(\Re \rho(g)-1)}} \\
& \stackrel{\sqrt{6.52]}}{\geq} \sum_{g \in G \backslash\{0\}}(1-\cos (2 \pi / n)) e^{12 \beta(\Re \rho(g)-1)} \\
& =\frac{(1-\cos (2 \pi / n)) \sum_{g \in G \backslash\{0\}} e^{12 \beta(\Re \rho(g)-1)}}{1+(|G|-1) \max _{g \in G \backslash\{0\}} e^{12 \beta(\Re \rho(g)-1)}} \\
& \geq \frac{(1-\cos (2 \pi / n)) \max _{g \in G \backslash\{0\}} e^{12 \beta(\Re \rho(g)-1)}}{1+(|G|-1) \max _{g \in G \backslash\{0\}} e^{12 \beta(\Re \rho(g)-1)}} \\
& \stackrel{\sqrt{2.88}}{=} \frac{(1-\cos (2 \pi / n)) \lambda(\beta)^{12}}{1+(|G|-1) \lambda(\beta)^{12}} \\
& =\frac{(1-\cos (2 \pi / n)) \lambda(\beta)^{12}\left(1-(|G|-1) \lambda(\beta)^{12}\right)}{\left(1+(|G|-1) \lambda(\beta)^{12}\right)\left(1-(|G|-1) \lambda(\beta)^{12}\right)} \\
& =\frac{(1-\cos (2 \pi / n)) \lambda(\beta)^{12}\left(1-(|G|-1) \lambda(\beta)^{12}\right)}{\left.1-\left((|G|-1) \lambda(\beta)^{12}\right)^{2}\right)} \\
& \geq(1-\cos (2 \pi / n)) \lambda(\beta)^{12}\left(1-(|G|-1) \lambda(\beta)^{12}\right) \\
& =4 C_{*} \lambda(\beta)^{12}\left(1-(|G|-1) \lambda(\beta)^{12}\right) \\
& \left.\frac{6.69}{\geq}\right) \\
& 2 C_{*} \lambda(\beta)^{12} .
\end{aligned}
$$

Thus,

$$
\begin{equation*}
e^{-\ell(1-\theta)} \leq e^{-2 C_{*} \ell \lambda(\beta)^{12}} \tag{6.70}
\end{equation*}
$$

and

$$
\begin{align*}
&\left|\mathbb{E}_{\beta, N}\left[W_{\gamma}\right]-e^{-\ell(1-\theta(\beta)}\right| \overline{(6.68), ~(6.70]} e^{\left.\frac{1}{2} C_{*} \ell \lambda(\beta)\right)^{12}}+e^{-2 \ell C_{*} \lambda(\beta)^{12}} \\
& \leq e^{-\frac{1}{2} \ell C_{*} \lambda(\beta)^{12}}+e^{-\frac{1}{2} \ell C_{*} \lambda(\beta)^{12}}  \tag{6.71}\\
&=2 e^{-\frac{1}{2} \ell C_{*} \lambda(\beta)^{12}}
\end{align*}
$$

From Proposition 6.1 follows that

$$
\begin{equation*}
\left|\mathbb{E}_{\beta, N}\left[W_{\gamma}\right]-e^{-\ell(1-\theta(\beta))}\right| \leq C_{A} e^{2 C^{*} \ell \lambda(\beta)^{12}}\left(\sqrt{\frac{\ell_{c}}{\ell}}+\lambda(\beta)^{2}\right) \tag{6.72}
\end{equation*}
$$

By combining this with inequality (6.71), we obtain

$$
\begin{aligned}
& \left|\mathbb{E}_{\beta, N}\left[W_{\gamma}\right]-e^{-\ell(1-\theta(\beta))}\right|^{1+4 C^{*} / C_{*}} \\
& \stackrel{\sqrt{6.722}, \sqrt[6.71 \mid]{ }}{\leq}\left(C_{A} e^{2 C^{*} \ell \lambda(\beta)^{12}}\left(\sqrt{\frac{\ell_{c}}{\ell}}+\lambda(\beta)^{2}\right)\right)^{1} \cdot\left(2 e^{-\frac{1}{2} \ell C_{*} \lambda(\beta)^{12}}\right)^{4 C^{*} / C_{*}} \\
& \quad \leq C_{A} 2^{4 C^{*} / C_{*}}\left(\sqrt{\frac{\ell_{c}}{\ell}}+\lambda(\beta)^{2}\right)\left(e^{2 C^{*} \ell \lambda(\beta)^{12}} \cdot e^{-2 C^{*} \ell \lambda(\beta)^{12}}\right) \\
& \quad=C_{A} 2^{4 C^{*} / C_{*}}\left(\sqrt{\frac{\ell_{c}}{\ell}}+\lambda(\beta)^{2}\right)
\end{aligned}
$$

Thus,

$$
\begin{aligned}
\left|\mathbb{E}_{\beta, N}\left[W_{\gamma}\right]-e^{-\ell(1-\theta(\beta))}\right| \leq & \left(C_{A} 2^{4 C^{*} / C_{*}}\left(\sqrt{\frac{\ell_{c}}{\ell}}+\lambda(\beta)^{2}\right)\right)^{1 /\left(1+4 C^{*} / C_{*}\right)} \\
= & \left(C_{A} 2^{4 C^{*} / C_{*}}\right)^{1 /\left(1+4 C^{*} / C *\right)}\left(\sqrt{\frac{\ell_{c}}{\ell}}+\lambda(\beta)^{2}\right)^{1 /\left(1+4 C^{*} / C_{*}\right)} \\
\leq & \sqrt{2}\left(C_{A} 2^{4 C^{*} / C_{*}}\right)^{1 /\left(1+4 C^{*} / C *\right)} \\
& \cdot\left(\sqrt{\frac{\ell_{c}}{\ell}}+\lambda(\beta)^{2}\right)^{1 /\left(1+4 C^{*} / C_{*}\right)} \\
= & C^{\prime}\left(\sqrt{\frac{\ell_{c}}{\ell}}+\lambda(\beta)^{2}\right)^{C^{\prime \prime}}
\end{aligned}
$$

We have proved the main theorem.

## Chapter 7

## Svensk sammanfattning

Denna avhandling är baserad på artikeln "Wilson loops in finite Abelian lattice gauge theories" av M. Forsström, J. Lenells och F. Viklund [1], där väntevärdet för en Wilsonloopobservabel uppskattas. Avhandlingens mål är att förklara artikeln så att matematikstuderande på magisternivå kan förstå den. I denna avhandling behandlas därför den matematiska teori som behövs för att förstå artikeln. Därtill ges bevisen i artikeln i en mera detaljerad och förklarande form.

Det första kapitlet i avhandlingen är en introduktion som ger en bakgrund till problemet som behandlas i denna avhandling. Matematiska modeller inom gittergaugeteorier har länge studerats och målet med dem är att försöka förklara kvantfältteorier inom standardmodellen. Även om dessa modeller inte ännu kan användas inom kvantfysiken är de värdefulla också som matematiska modeller.

Det andra kapitlet introducerar artikelns huvudsats. Den teori som behövs för att satsen ska kunna förstås diskuteras och satsen ges. Grundläggande teori om grupper och representationer behandlas och både gruppen och representationen som används i artikeln definieras. Gruppen som används är $G=\left(\mathbb{Z}_{n},+\right)$ som är en ändlig abelsk grupp. Representationen som används är endimensionell, unitär och injektiv. Det visas att om en representation uppfyller dessa krav för den valda gruppen så ges den av $\rho(g)=e^{g \cdot 2 \pi i m / n}$, där $g$ tillhör gruppen $\mathbb{Z}_{n}$ och $m \in\{0,1, \ldots, n-1\}$ är relativt primt till $n$. Därtill definieras det fyrdimensionella gittret $\mathbb{Z}^{4}$ vars noder är i varje heltalspunkt. Delmängder av detta gitter och begrepp som hör till det ges, bland annat plaketter ${ }^{11}$ som bildas av två riktade bågar. Därefter tas begrepp för loopar och bågar upp innan ett mått

[^0]med Wilsonverkan ges och Wilsonloopobservabeln $W_{\gamma}$ definieras. Gränsvärdet för denna observabels väntevärde betecknas med $\left\langle W_{\gamma}\right\rangle_{\beta}$. Kapitlet avslutas med att satsen ges och kommenteras. Som följande ges satsen.

Sats. Lait $n$ vara ett heltal med $n \geq 2$. Låt gittergaugeteorin ha strukturgruppen $G=\mathbb{Z}_{n}$. Låt representationen $\rho$ vara en endimensionell och injektiv representation av $G$. Låt $\gamma \in \mathbb{Z}^{4}$ vara en riktad cykel, $\ell$ dess längd och $\ell_{c}$ antalet hörnbågar $i \gamma$. Då existerar gränsvärdet av Wilsonloopobservabelns väntevärde och $\beta_{0}$ kan väljas så att det existerar konstanter $C^{\prime}\left(\beta_{0}\right)$ och $C^{\prime \prime}\left(\beta_{0}\right)$ för vilka följande olikhet gäller för alla $\beta \geq \beta_{0}$ :

$$
\left|\left\langle W_{\gamma}\right\rangle_{\beta}-e^{-\ell(1-\theta(\beta))}\right| \leq C^{\prime}\left(\frac{\sqrt{\ell_{c}}}{\ell}+\lambda(\beta)^{2}\right)^{C^{\prime \prime}}
$$

Funktionen $\theta(\beta)$ är definierad $i$ ekvation (2.7) och $\lambda(\beta)$ i ekvation (2.8).
Från denna sats kan dras som slutsats att väntevärdet på Wilsonloopobservabeln tar ett värde väldigt nära 0 om $\ell(1-\theta(\beta))$ är stort och ett värde väldigt nära $1 \mathrm{om} \ell(1-\theta(\beta))$ är litet.

Det tredje kapitlet behandlar teori om diskret yttre algebra då $\mathbb{Z}^{r}$ har valts som gitter. Denna teori diskuteras alltså mera generellt än vad som skulle behövas för resten av avhandlingen. Riktade differentialceller och -former både definieras och ges exempel på (se Figur 3.1 för exempel på differentialceller). Dessa kallas också för $k$-celler och $k$-former. En riktad 1-cell är detsamma som en båge i gittret $\mathbb{Z}^{r}$ och en riktad 2-cell detsamma som en plakett. Därefter ges två operatorer för $k$-former. Dessa är den yttre derivatan, som avbildar en $k$-form på en $k+1$-form, och koderivatan som avbildar en $k$-form på en $k-1$-form. Exempel illustrerar hur dessa används. Ränder och koränder för $k$-former och speciellt en plakett behandlas. Dessutom ges två olika versioner av Poincarés lemma som både appliceras på en viss mängd av 2 -former och används för att visa att måttet med Wilsonverkan kan skrivas om så att det blir ett mått för plakettkonfigurationer. Till sist diskuteras den diskreta Hodge dual-operatorn och bijektionen mellan det ursprungliga gittret och det som skapas av Hodge dual-operatorn. Lemman med Hodge dual-operatorn som är nödvändiga för fjärde kapitlets bevis ges.

Det fjärde kapitlet behandlar teori för virvlar ${ }^{2}$ och riktade ytor. Dessa begrepp definieras och illustreras, observera att dessa definitioner kan variera mellan

[^1]olika källor. Först behandlas uppdelningar av virvlar, sedan definieras minimala virvlar och lemman med dessa som behövs till huvudbeviset ges och bevisas. Dessutom ges en proposition för en viss fördelning för virvlar som har en stor roll i huvudbeviset. Till sist ges några lemman för riktade ytor och en riktad ytas koppling till loopar förklaras. Dessutom definieras inre bågar och inre plaketter och ett lemma för riktade ytor ges.

I det femte kapitlet behandlas gränsvärdet för väntevärdet av en Wilsonloopobservabel. Det viktigaste resultatet i detta kapitel är en sats som säger att måtten $\mathbb{E}_{\beta, N}$ konvergerar svagt i topologin av lokal konvergens när $N$ närmar sig oändligheten och det begränsade måttet är translationsinvariant i $\mathbb{Z}^{r}$. Denna sats bevisas med hjälp av Ginibres olikhet. Till sist tillämpas satsen på Wilsonloopobservabeln för att visa att dess väntevärdes gränsvärde både existerar och är invariant under translation.

I det sjätte kapitlet bevisas slutligen huvudsatsen med hjälp av den teori och de resultat som getts i de tidigare kapitlen. Beviset delas upp i två fall, först i Proposition 6.1 där $\ell \lambda(\beta)^{12}$ är stort och därefter i Proposition 6.11 där $\ell \lambda(\beta)^{12}$ är litet. Beviset för Proposition 6.1 är uppdelat i flera delar där olika väntevärden beräknas innan de kombineras för att bevisa den önskade olikheten. Till beviset behövs flera lemman som ges innan. Dessa lemman är främst olikheter för funktionen $\theta$ samt sannolikheter för olika händelser, vilka beräknas med hjälp av propositionen från det fjärde kapitlet. Därefter bevisas Proposition 6.11. Till detta bevis behövs några lemman som först ges och bevisas och även vissa slutsatser från beviset för Proposition 6.1 används. Därefter kombineras dessa två fall för att bevisa olikheten

$$
\left|\mathbb{E}_{\beta, N}\left[W_{\gamma}\right]-e^{-\ell(1-\theta(\beta))}\right| \leq C^{\prime}\left(\frac{\sqrt{\ell_{c}}}{\ell}+\lambda(\beta)^{2}\right)^{C^{\prime \prime}}
$$

Med hjälp av resultatet för gränsvärdets existens för Wilsonloopobservabeln kan satsen därefter bevisas genom att ta gränsvärdet då $N$ närmar sig oändligheten.

För de som är intresserade av liknande problem finns bland annat artiklarna "Wilson loops in Ising lattice gauge theory" av S. Chatterjee [2] och "Wilson loop expectations in lattice gauge theories with finite gauge groups" av S. Cao [3] som båda publicerades år 2020.

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[^0]:    ${ }^{1}$ En svensk term är inte etablerad (eng. plaquette).

[^1]:    ${ }^{2}$ En svensk term är inte etablerad (eng. vortex).

